Effects of Superradiance in Active Galactic Nuclei

with Himanshu Verma, Kingman Cheung, Joseph Silk [arXiv: 2404.09955]

Priyanka Sarmah

Postdoctoral fellow

National Tsing Hua University, Taiwan

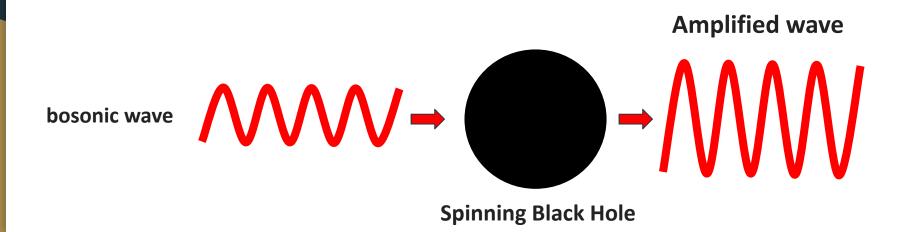


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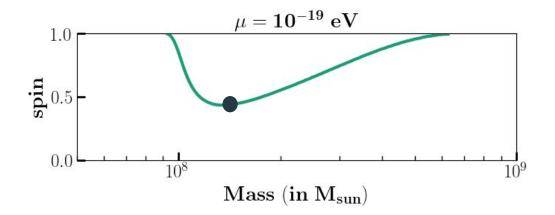


Black Hole Superradiance

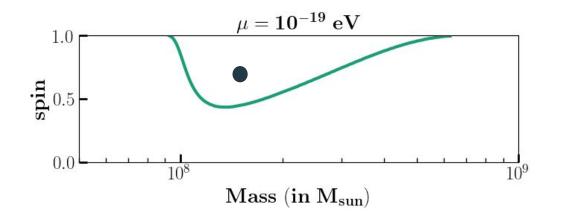
- Spinning supermassive BH opens a room for ultralight scalar particles to get produced through a phenomenon- *Superradiance (SR)*
- A bosonic cloud grow near the BH, *draining* the angular momentum of the BH



Observational signatures of Superradiance

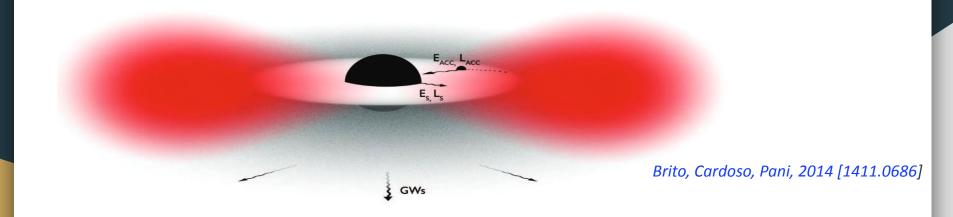


Observational signatures of Superradiance



Observation of a BH inside the depletion region in the Regge plane exclude the scalar

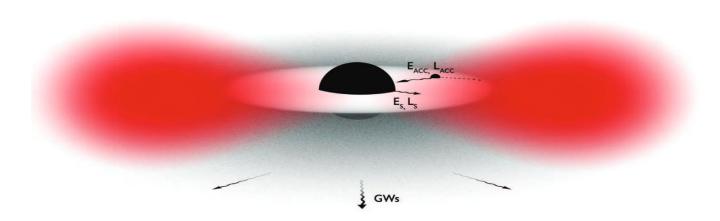
Realistic environment for BH Superradiance: The Active Galactic Nucleus (AGN)



- Key points: Role of accretion in adding mass and angular momentum to the BH
- 2 competing process: Spin up- accretion, Spin down- Superradiance



How do the characteristics of AGN alter due to Superradiance history of the BH?



Key Findings

As the accreting SMBH spins down due to superradiance:

- **Sudden drops** in the time-variation of the luminosities of AGNs in various wavelength bands.
- Observation of depletion regions in various planes of band-luminosities and f_{Edd} and accumulation of AGN along the boundaries of the depletion region.

Superradiance in a nutshell

• Condition of Superradiance(SR):

 $\omega_{\rm R} < m\Omega$,

 $\omega_{\rm R}, \Omega$ = angular velocity of the particle and BH

• Consequence of Superradiance: Growth of scalar cloud, BH loses mass and angular momentum.

• Angular momentum lost till : $\tilde{a} \sim \tilde{a}_{critical} = 4\alpha m/(m^2 + \alpha^2)$,

gravitational fine structure constant - $\alpha \sim GM\mu$

Luminosity, Eddington Ratio of AGN

• Total Luminosity :

$$L = \epsilon(\tilde{a})\dot{M}_{\rm disk}c^2$$

Radiative efficiency

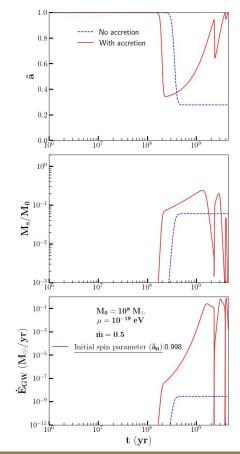
Fanidakis et al, 2011, MNRAS, 410, 53

 $\dot{m} \equiv \dot{M}_{\rm disk} c^2 / L_{\rm Edd}$ $L = \epsilon(\tilde{a}) \dot{m} L_{\rm Edd}.$

$$L_{\rm Edd} = \frac{4\pi G M m_p c}{\sigma_T} \approx 1.26 \times 10^{38} \rm erg/s \frac{M}{M_{\odot}}$$

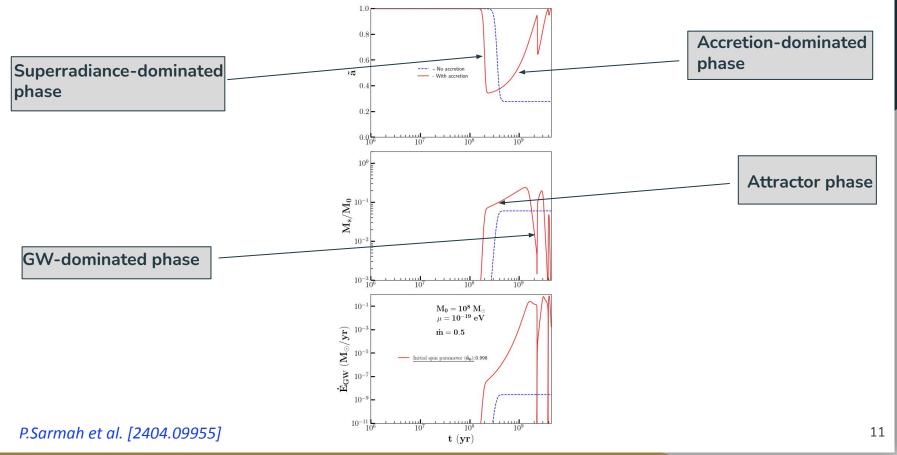
• Eddington Ratio: $f_{\text{Edd}} \equiv L/L_{\text{Edd}}$, $f_{\text{Edd}} = \epsilon(\tilde{a})\dot{m}$.

Time evolution of accreting BH + scalar cloud system



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Time evolution of accreting BH + scalar cloud system

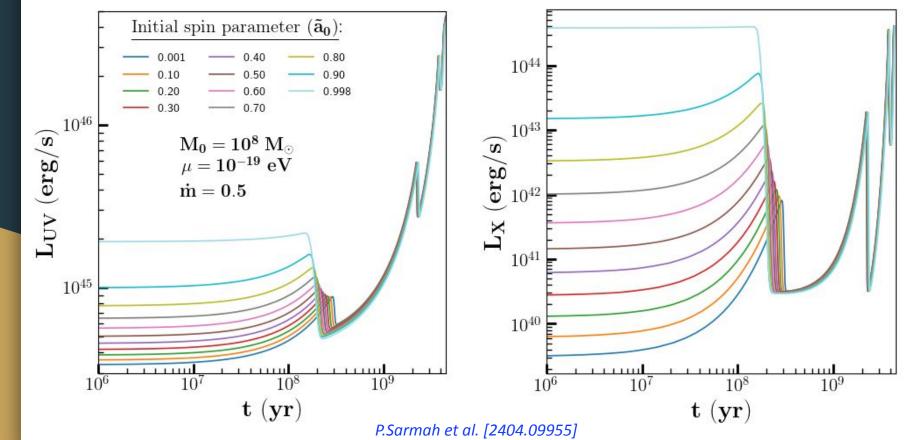


Luminosity in various bands

Using Novikov-Thorne model of the accretion disk, get the spin-dependent flux $F_{\lambda}(\tilde{a}, r)$

$$L_{\rm X} = \int_{10^{-4}}^{0.01} F_{\lambda} d\lambda,$$
$$L_{\rm UV} = \int_{0.01}^{0.4} F_{\lambda} d\lambda,$$
$$L_{\rm Vis-IR} = \int_{0.4}^{100} F_{\lambda} d\lambda,$$

Luminosity in various bands

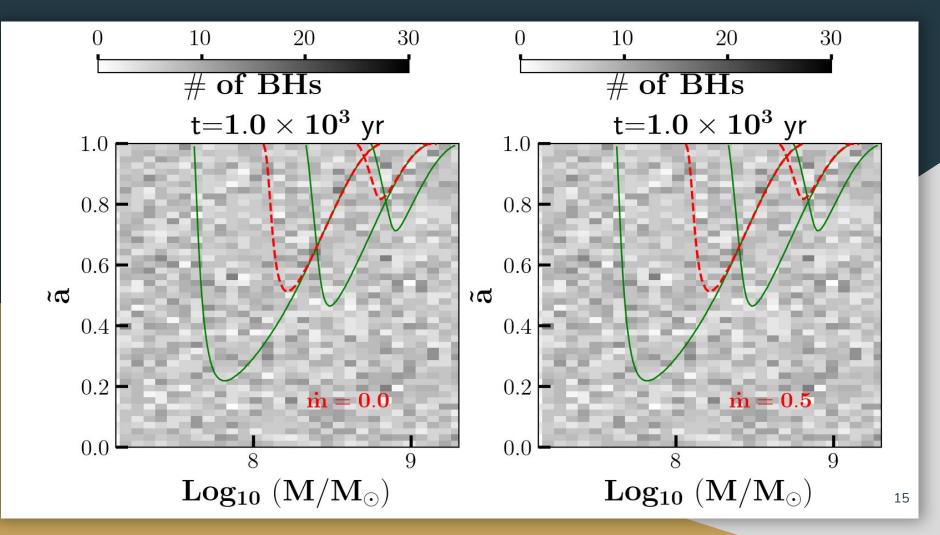


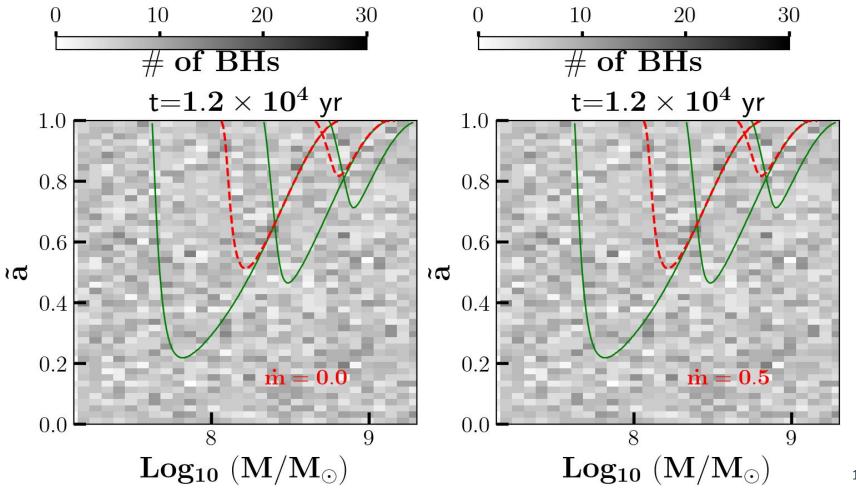


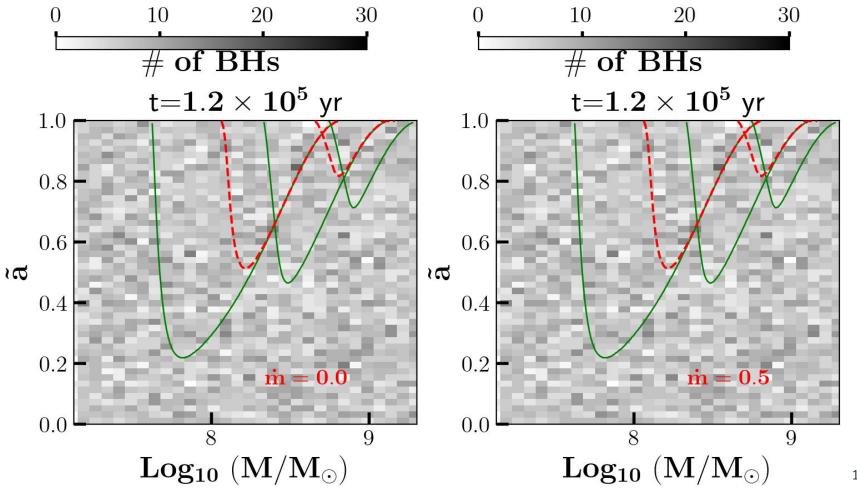


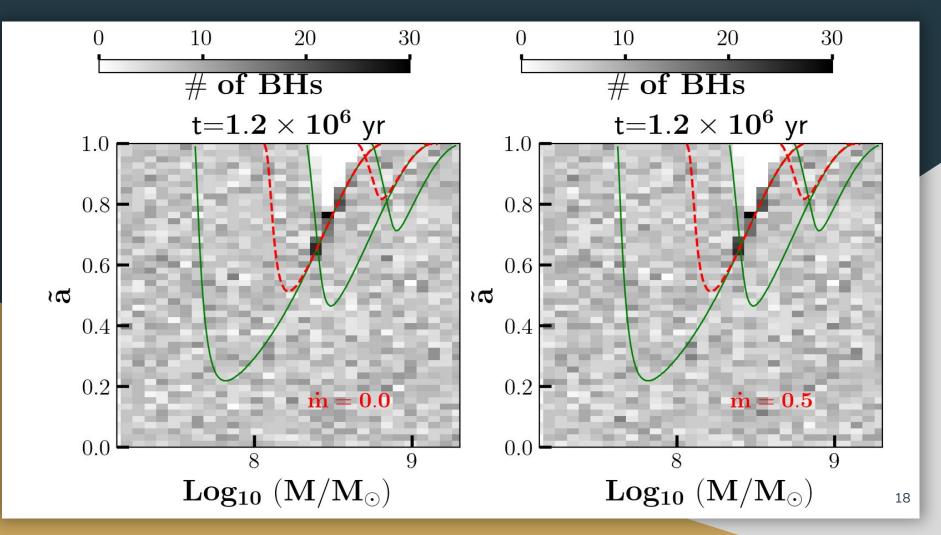
Distribution of SMBHs at the AGN core

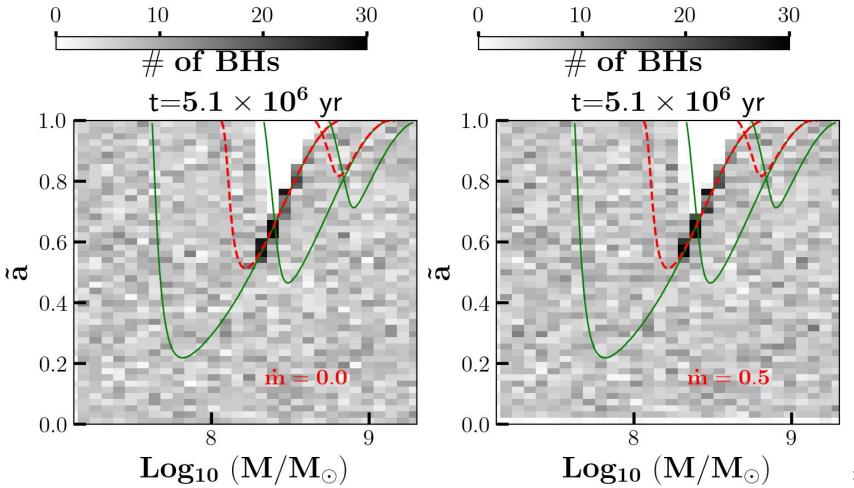


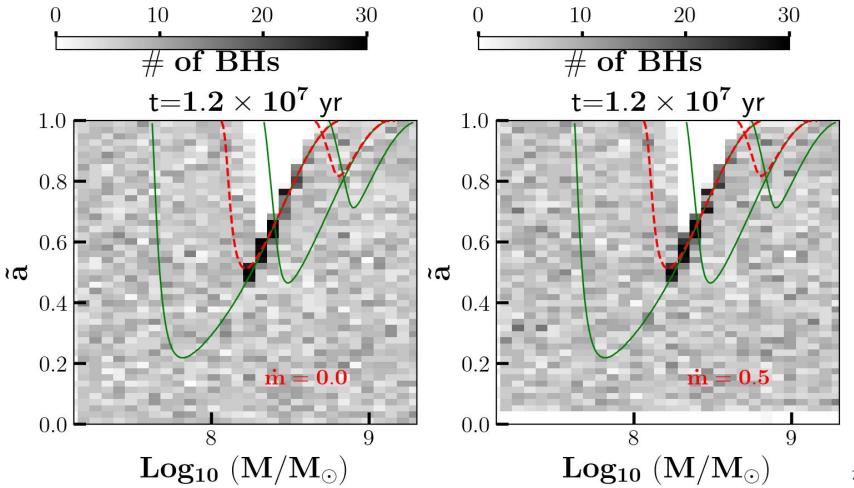


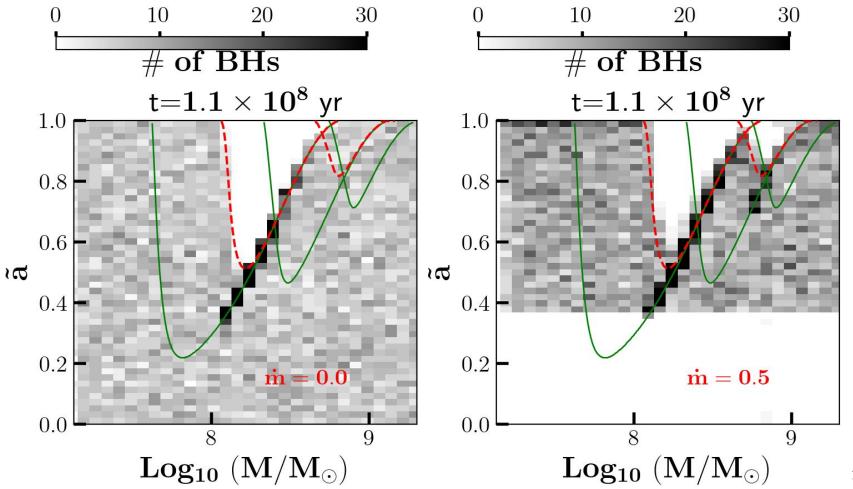


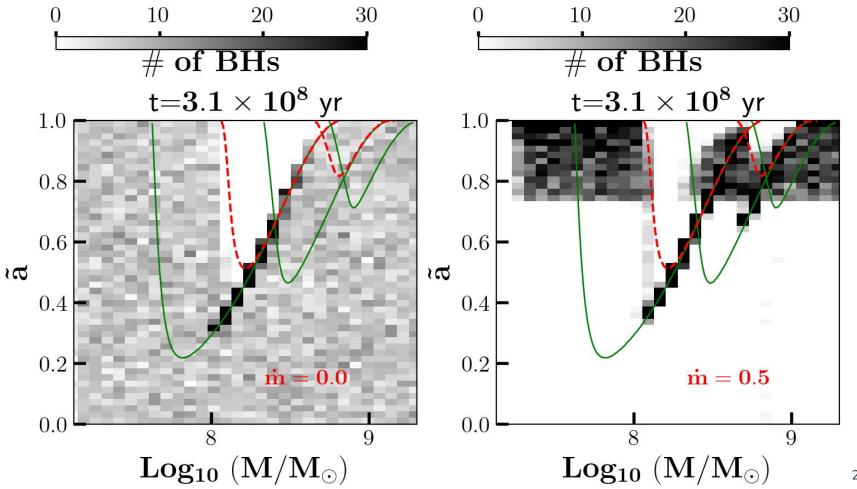


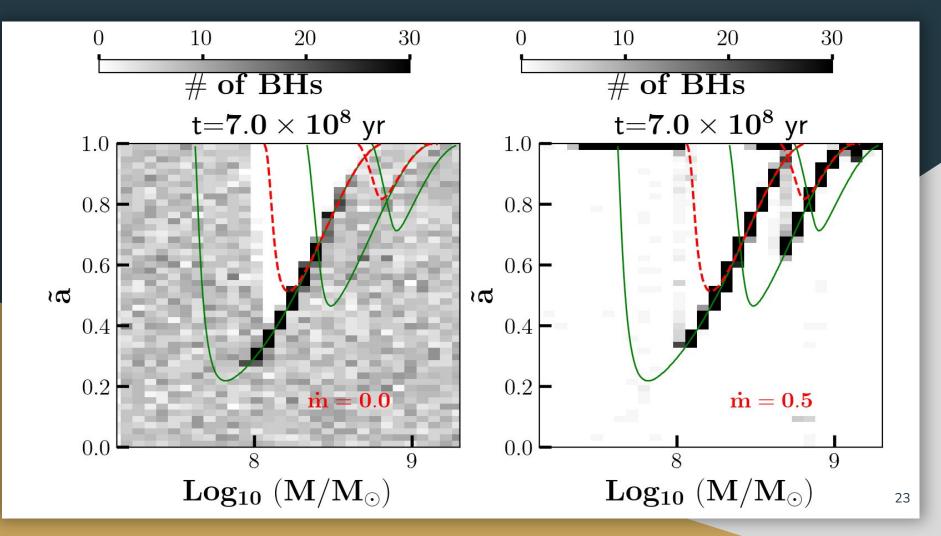










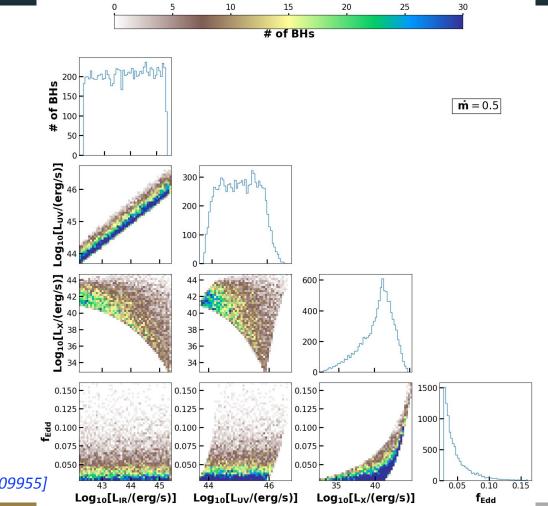




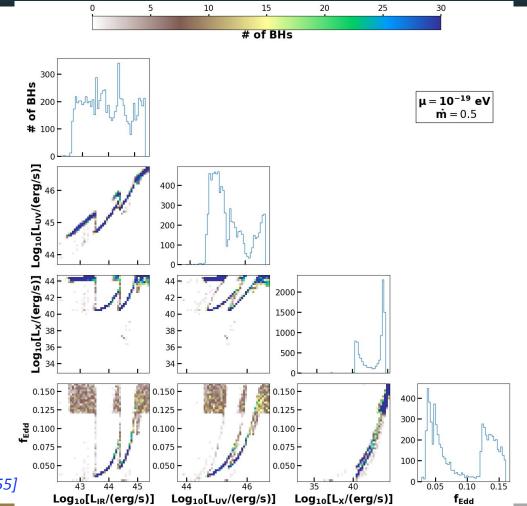


Distribution of AGN Characteristics





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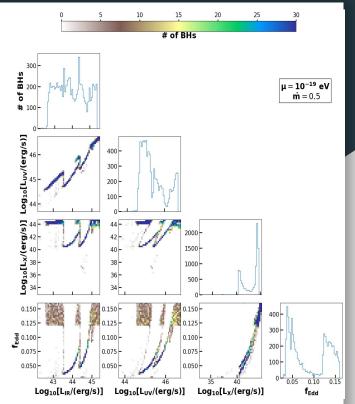
P.Sarmah et al. [2404.09955]

Summary

Accreting SMBH undergoing Superradiance at the core of AGN leads to-

Enhanced growth of scalar cloud and GW emission rate and appearance of higher modes within the age of the universe.

- **Multiple dips** in the luminosity evolution corresponding to timescales of dominant modes of superradiance.
- Observation of depletion regions in various planes of band-luminosities and f_{Edd} and accumulation of AGN along the boundaries of the depletion region.

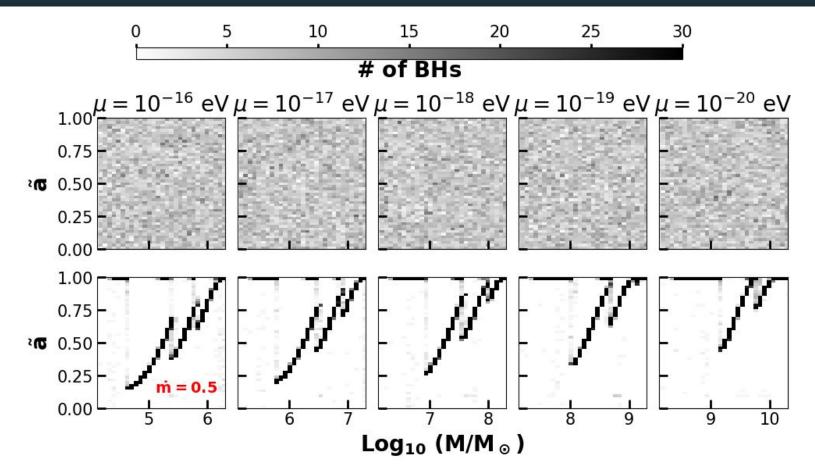


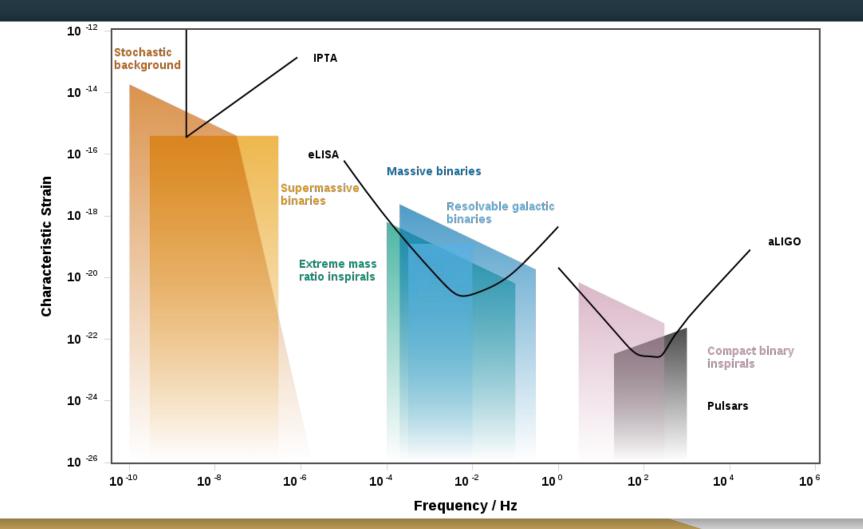


Thank you!

Questions? Comments? Suggestion?







Superradiance in a nutshell

• The metric around a rotating BH parameterized in terms of BH mass M and spin a= ã M, ã dimensionless spin parameter

 $\Box \Phi + \mu^2 \Phi = 0$ $\Phi = S_{\mu}(\Theta) \psi(\mathbf{r}) / \mathbf{r} \exp(-i\omega t + im\varphi)$

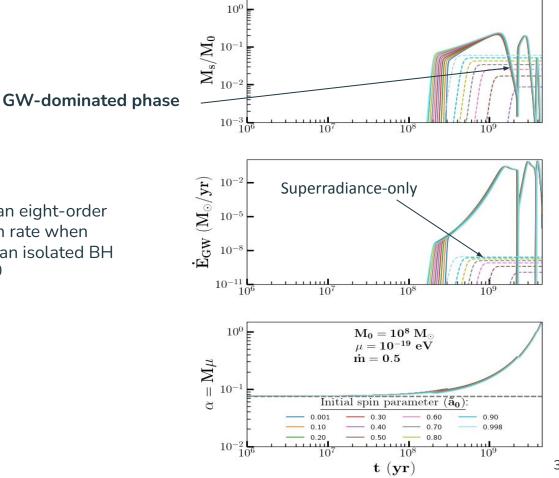
• Energy eigenvalue $\sim \omega_{\rm R} + i (m\Omega - \omega_{\rm R})$

Observational signatures of Superradiance

• Interesting signatures of gravitational wave emission emitted from the annihilation of scalars in the cloud around the BH, *Arvanitaki et al. 2015b*

• Scalar cloud affecting the black hole images, *Davoudiasl & Denton 2019, Saha et al.* 2022

Depletion region in Regge plane i.e. spin versus mass plane of the BH, *Brito et al.* 2014

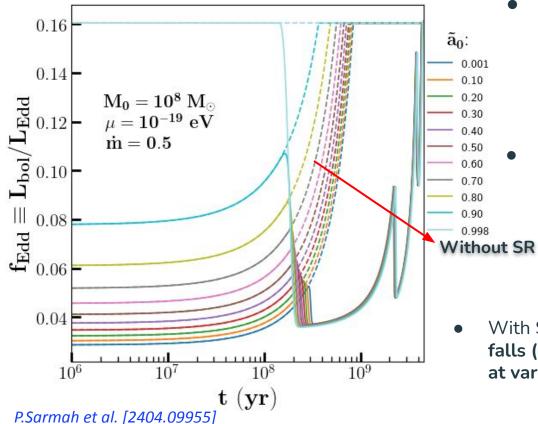


• **GW-dominated phase:** observe an eight-order increase in the peak GW emission rate when accretion is present compared to an isolated BH $dE_{GW}/dt \sim (Ms/M)^2 \alpha^{4l+10}$

Yoshino H., Kodama H.'14

P.Sarmah et al. [2404.09955]

Eddington Ratio



• **sudden drops** at the time-scales corresponding to various modes of superradiant growth.

 $f_{\rm Edd} = \epsilon(\tilde{a})\dot{m}.$

Without scalar field, **f**_{Edd} **monotonically increases with time due to accretion.**

With SR, no longer monotonically increasing, falls (due to SR) and rise (due to accretion) at various epochs.

Time evolution of BH + scalar cloud system

$$\begin{split} \frac{dM}{dt} &= -\sum_{nlm} 2M_s^{nlm} \omega_I^{nlm} + \dot{M}_{\rm Acc} \ , \\ \frac{dJ}{dt} &= -\sum_{nlm} \frac{2}{\mu} m M_s^{nlm} \omega_I^{nlm} + \dot{J}_{\rm Acc} \ , \\ \frac{dM_s^{nlm}}{dt} &= 2M_s^{nlm} \omega_I^{nlm} - \dot{E}_{\rm GW}^{nlm} \ , \\ \frac{dJ_s^{nlm}}{dt} &= \frac{2}{\mu} m M_s^{nlm} \omega_I^{nlm} - \frac{1}{\mu} m \dot{E}_{\rm GW}^{nlm} \ , \end{split}$$

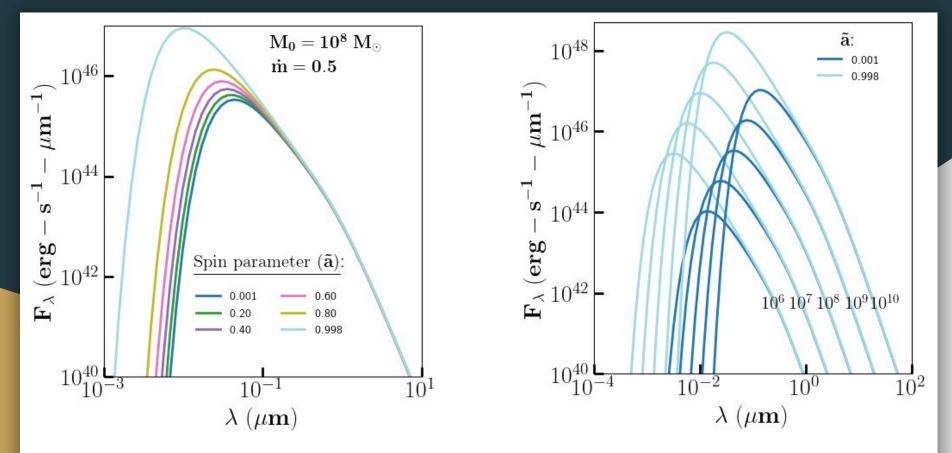
Accretion disk around Kerr BH: Novikov-Thorne model

$$F(r) = 7 \times 10^{26} \frac{\mathrm{erg}}{\mathrm{s \ cm^2}} \dot{m} \frac{M_{\odot}}{M} \left(\frac{M}{r}\right)^3 \mathcal{B}^{-1} C^{-1/2} Q$$

where B, C, Q are functions of BH spin \tilde{a} and radius r

Spectrum is obtained by integrating the flux, assuming the flux coming from local Black body

$$F_{\lambda} = 2 \int f_{\lambda}(r) r dr d\phi = 4\pi \int f_{\lambda}(r) r dr$$



 Most visible effects in the X-ray and UV band luminosities of AGNs, least effect in Vis-IR: higher energetic photons come from the inner part

P.Sarmah et al. [2404.09955]