

Revealing the production mechanism of high-energy neutrinos from NGC1068

TeVPA 2024

Abhishek Das

Advisor: Dr. Kohta Murase (PSU, Yukawa)

Collaborator: Dr. B. Theodore Zhang (Yukawa)



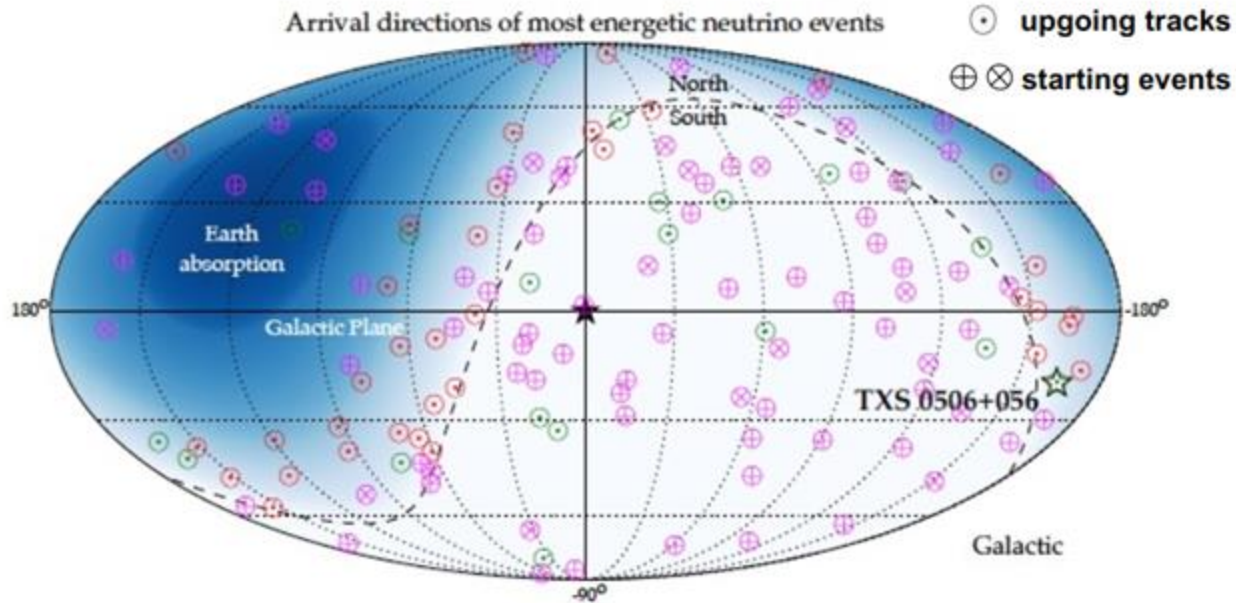
PennState



Chicago 2024

High-energy neutrino sky

Slide courtesy: Kohta Murase

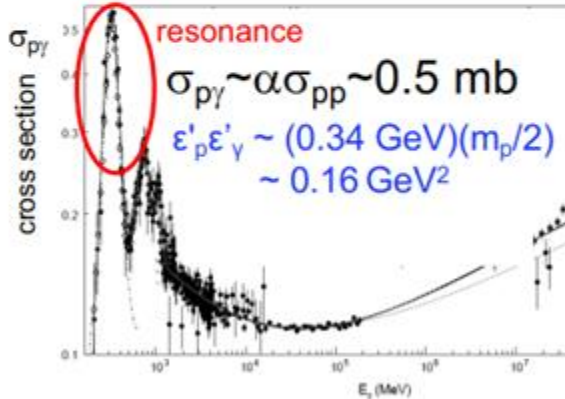
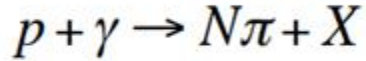
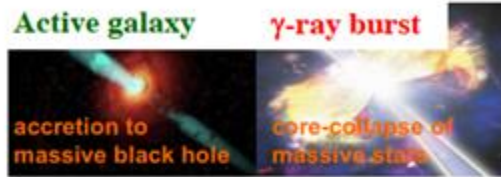


consistent w. **isotropic** distribution/**extragalactic** origins

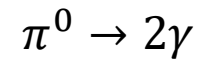
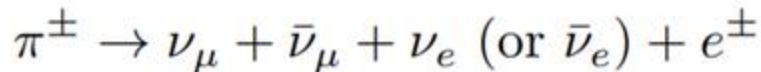
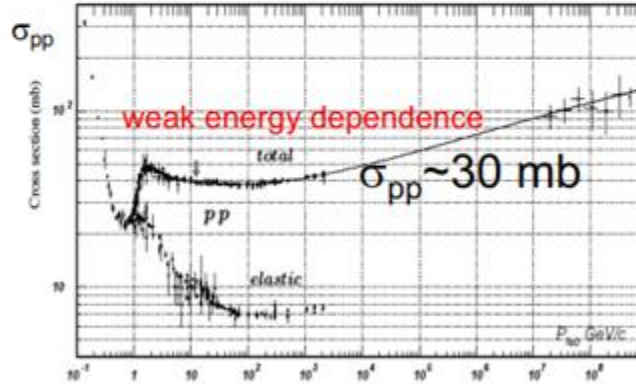
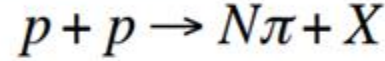
#Galactic contribution: ~10% (IceCube 23 Science)

Sources of extragalactic neutrinos

Cosmic-ray Accelerators



Cosmic-ray Reservoirs



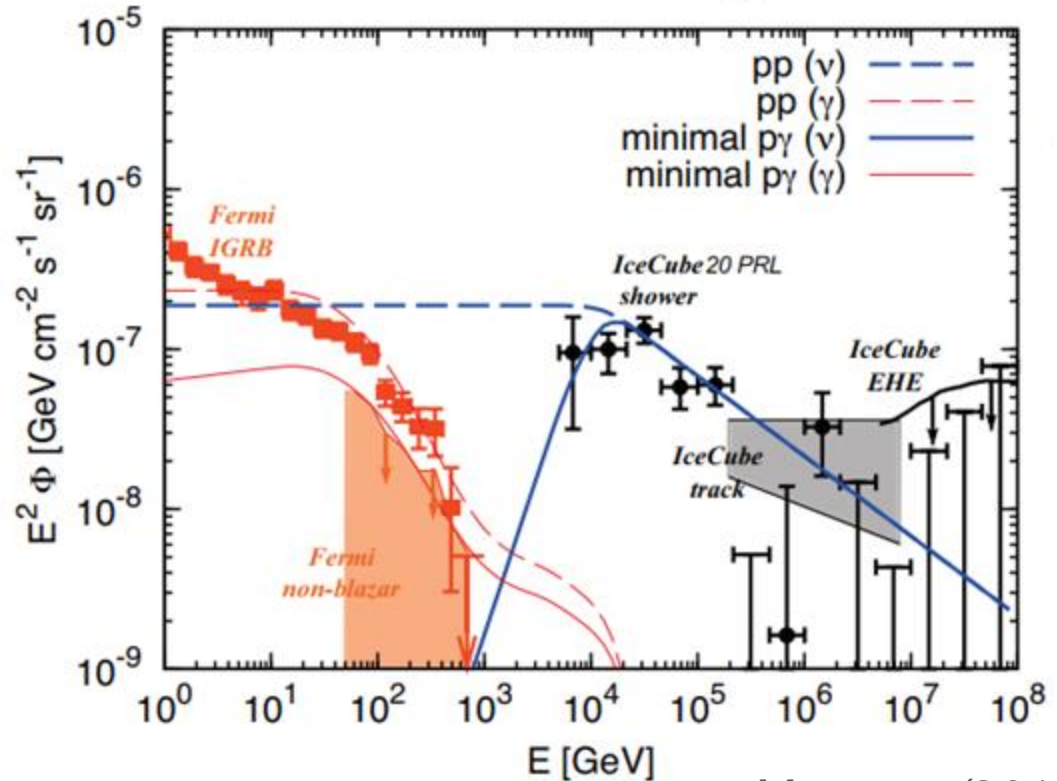
Type equation here.

Extragalactic multimessenger connection

-Fermi diffuse gamma-ray background is violated ($>3\sigma$) if sources are transparent.

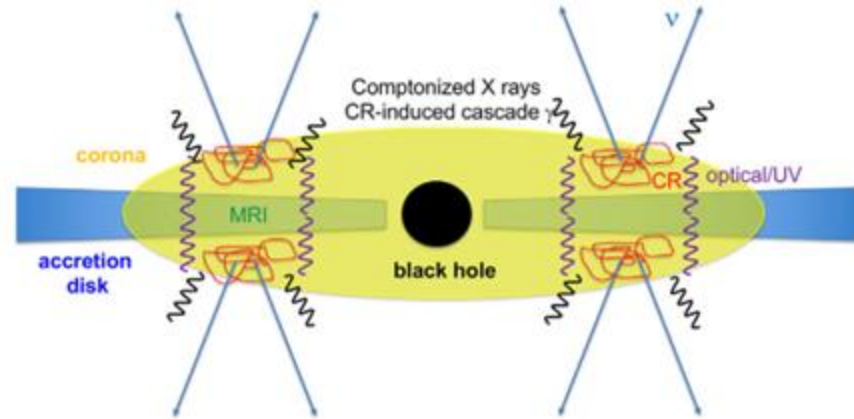
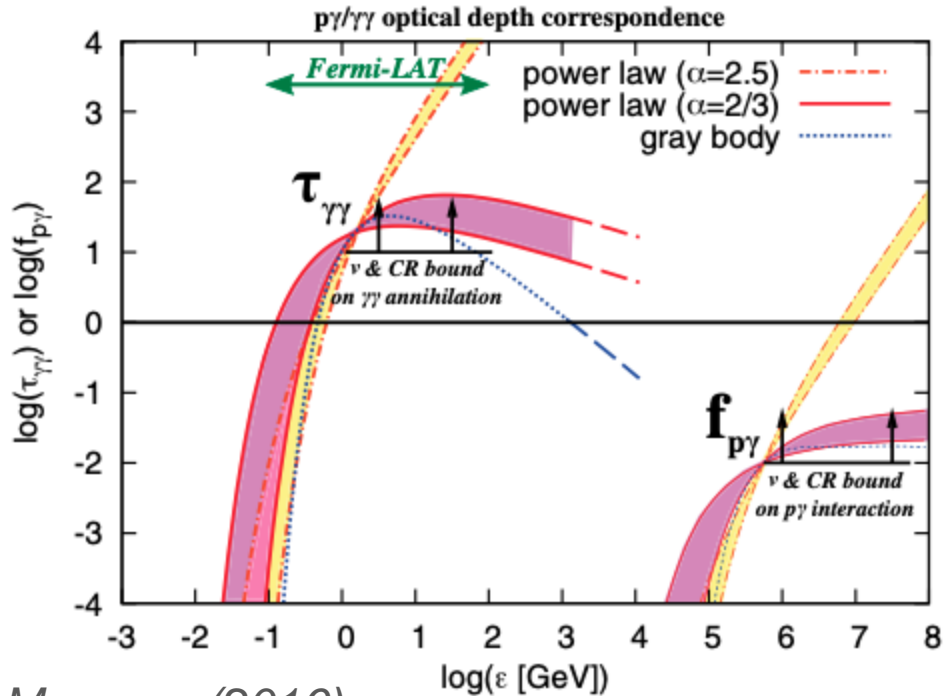
Solution:

Hidden (gamma-ray opaque) cosmic-ray accelerators



Opacity argument and prediction

$$\text{optical depth } \tau_{\gamma\gamma} \approx \frac{\sigma_{\gamma\gamma}^{\text{eff}}}{\sigma_{p\gamma}^{\text{eff}}} f_{p\gamma} \sim 1000 f_{p\gamma} \gtrsim 10$$



Prediction: Cosmic-ray acceleration site must be much closer to SMBH

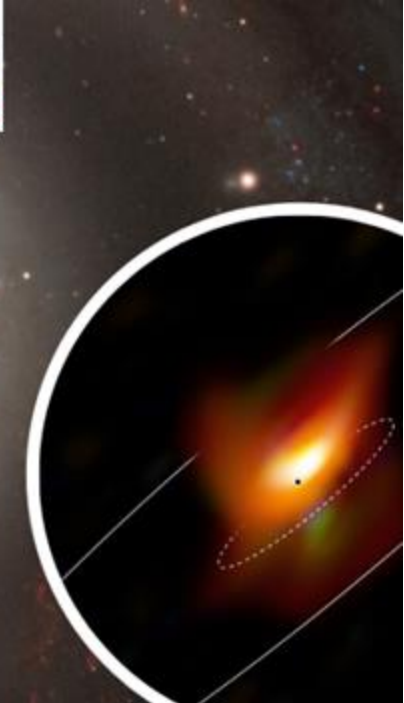
Slide courtesy: Kohta Murase

IceCube detects Neutrinos from NGC 1068

NEUTRINO ASTROPHYSICS
Evidence for neutrino emission from the nearby active galaxy NGC 1068
Science
IceCube Collaboration†
JOURNALS AAAS

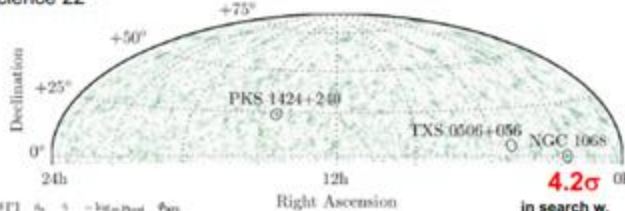


ASTRONOMY
Neutrinos unveil hidden galactic activities
By Kohta Murase^{1,2*}
An obscured supermassive black hole may be producing high-energy cosmic neutrinos

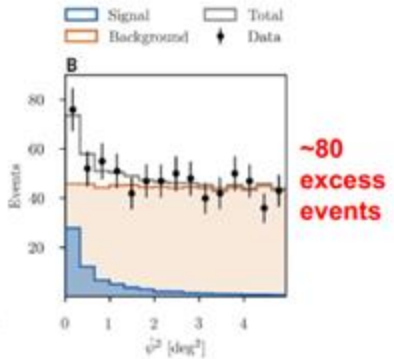
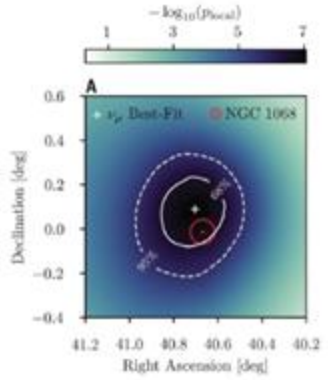


IceCube Collaboration+ Science 22

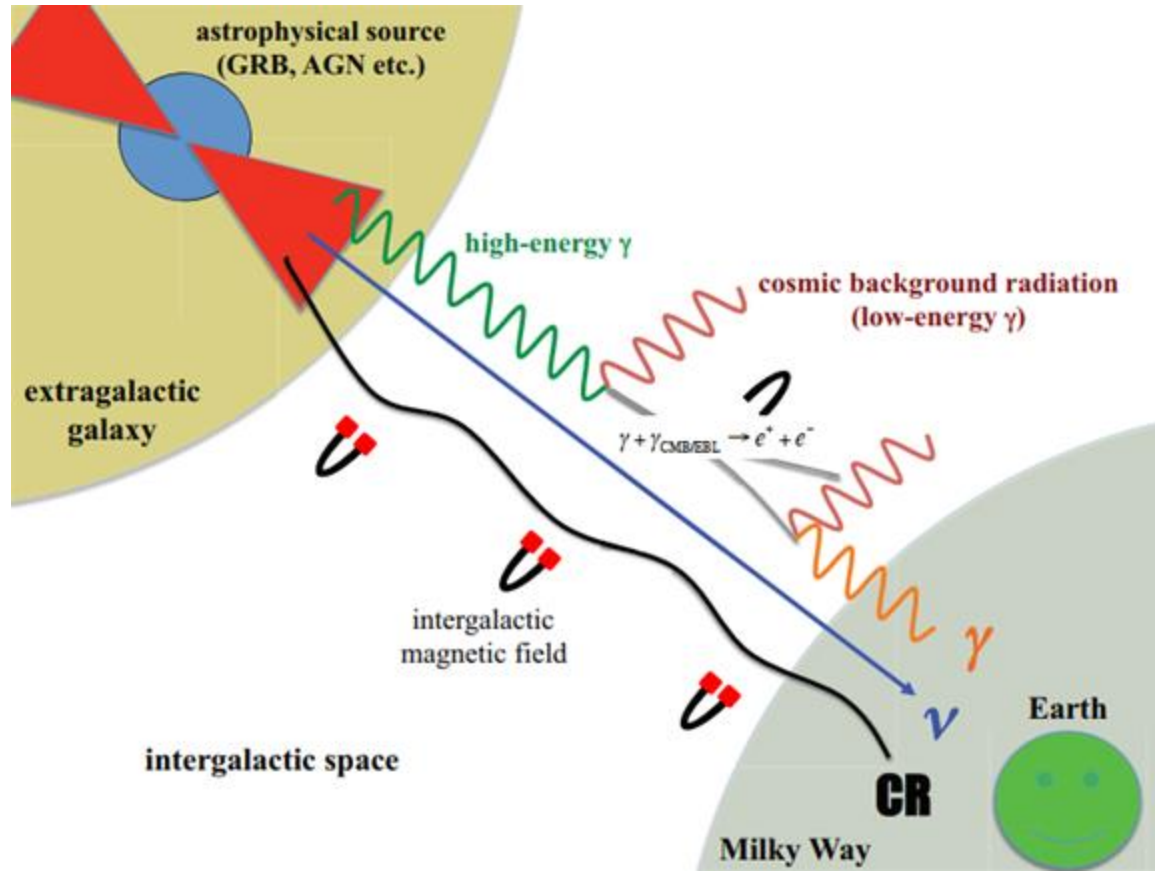
2011-2020
10 year track data



Source Name	Source Type	α [°]	δ [°]	z	$\log_{10}(P_{\text{total}})$	P_{sig}	
NGC 1068	SBRGAGN	40.67	-4.01	79	3.2	7.0 (5.2 σ)	9.6
PKS 1424+240	BL	216.26	21.80	77	3.5	4.0 (3.7 σ)	11.4
TXS 0506+056	BLFSRQ	77.36	5.70	5	2.0	3.6 (3.5 σ)	7.5



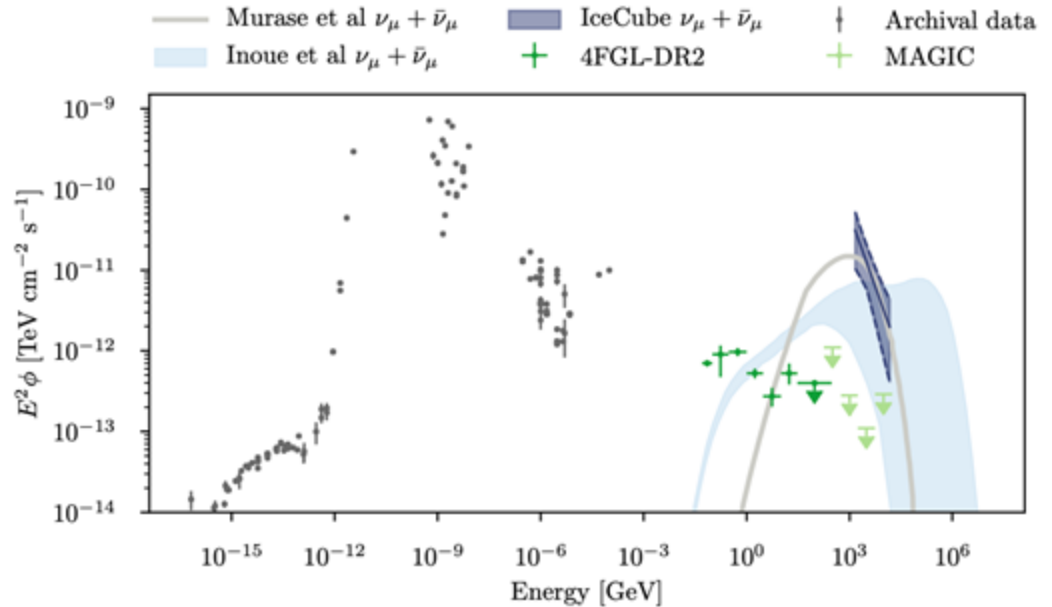
Neutrino advantage



NGC 1068 as a Hidden source

- Between 1.5 TeV and 15 TeV,
 $L_\nu = (2.9 \pm 1.1) \times 10^{42} \text{ erg/s}$
- Between 100 MeV and 200 GeV, *Fermi*-LAT reports
 $L_\gamma = 1.6 \times 10^{41} \text{ erg/s}$

Conclusion:
Emission region must be **optically thick to gamma rays**.

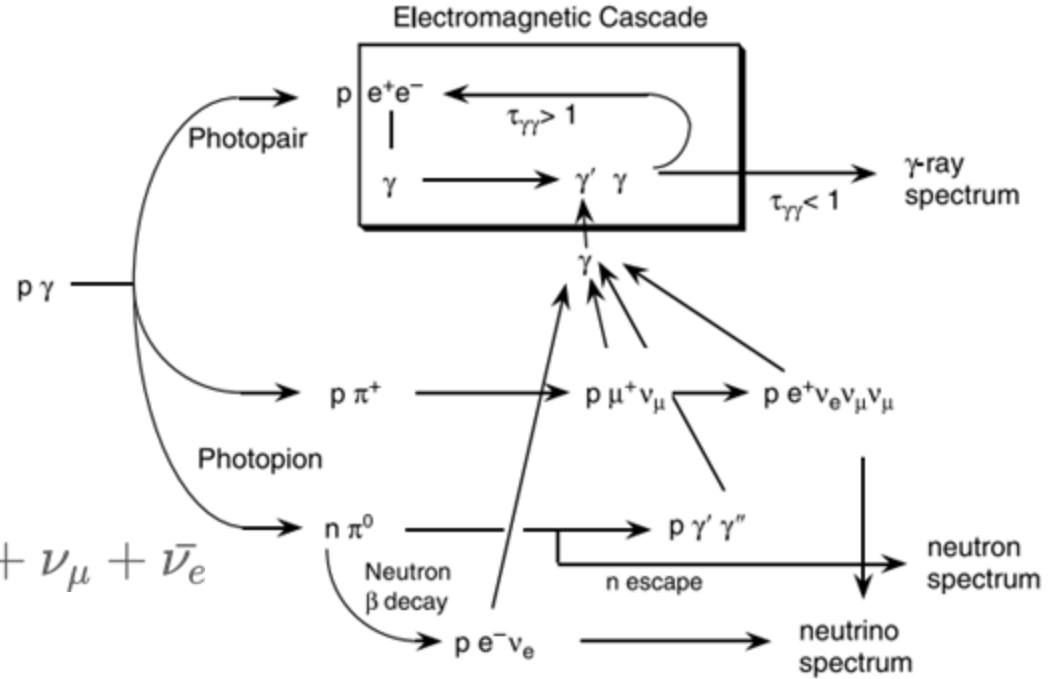
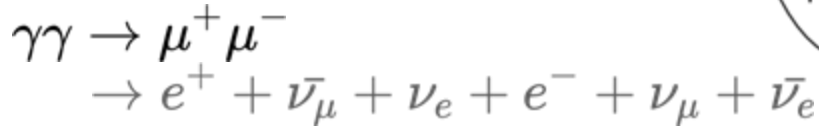


Abbasi+ (2022)




Cascade formation and observed spectra

- Ultrarelativistic CR protons interact with target photons forming **electromagnetic cascades**.

- Leptonic mechanism



Revealing the Production Mechanism of High-energy Neutrinos from NGC 1068

Abhishek Das^{1,2,3} , B. Theodore Zhang⁴ , and Kohta Murase^{1,2,3,4} 

Published 2024 August 23 • © 2024. The Author(s). Published by the American Astronomical Society.

[The Astrophysical Journal](#), [Volume 972](#), [Number 1](#)

Citation Abhishek Das *et al* 2024 *ApJ* **972** 44

DOI 10.3847/1538-4357/ad5a04



Article PDF



Article ePub

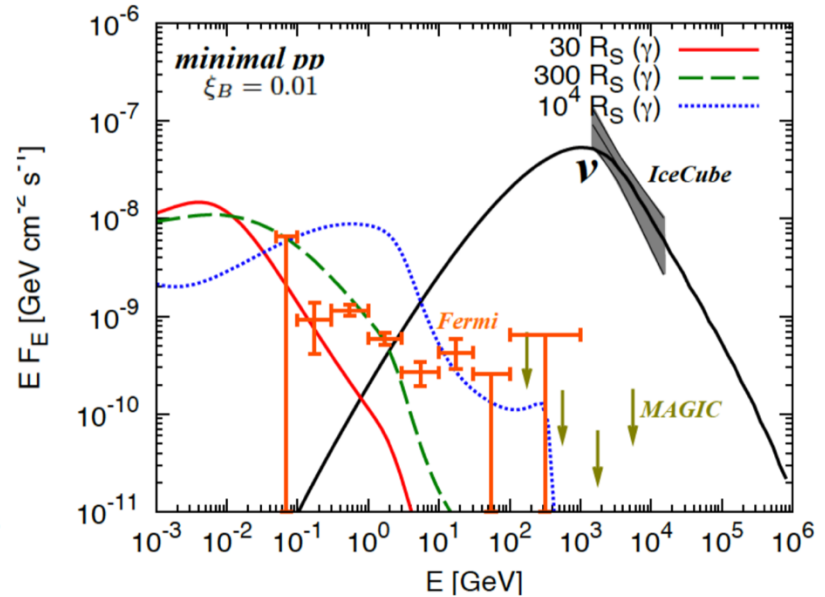
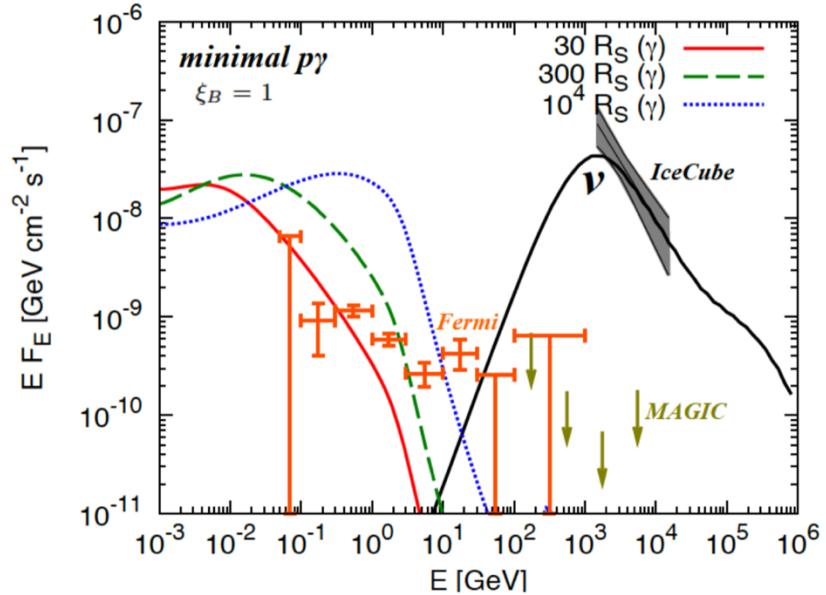
[Figures](#) ▼ [References](#) ▼

▼ [Article and author information](#)

Abstract

The detection of high-energy neutrino signals from the nearby Seyfert galaxy NGC 1068 provides us with an opportunity to study nonthermal processes near the center of supermassive black holes. Using the IceCube and latest Fermi-LAT data, we present general multimessenger constraints on the energetics of cosmic rays and the size of neutrino emission regions. In the photohadronic scenario, the required cosmic-ray luminosity should be larger than $\sim 1\%–10\%$ of the Eddington luminosity and the emission radius should be $\leq 15R_S$ in low- β plasma and $\leq 3R_S$ in high- β plasma. The leptonic scenario overshoots the NuSTAR or Fermi-LAT data for any emission radii we consider, and the required gamma-ray luminosity is much larger than the Eddington luminosity. The beta-decay scenario also violates not only the energetics requirement but also gamma-ray constraints, especially when the Bethe–Heitler and photomeson production processes are consistently considered. Our results rule out the leptonic and beta-decay scenarios in a nearly model-independent manner and support hadronic mechanisms in magnetically powered coronae if NGC 1068 is a source of high-energy neutrinos.

Hadronic scenario

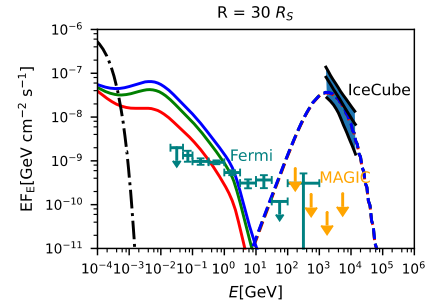
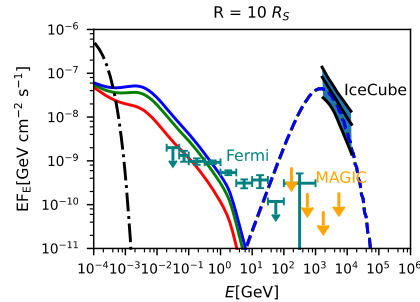
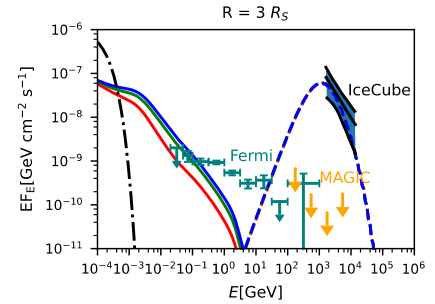
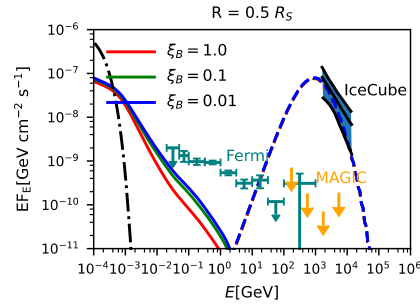
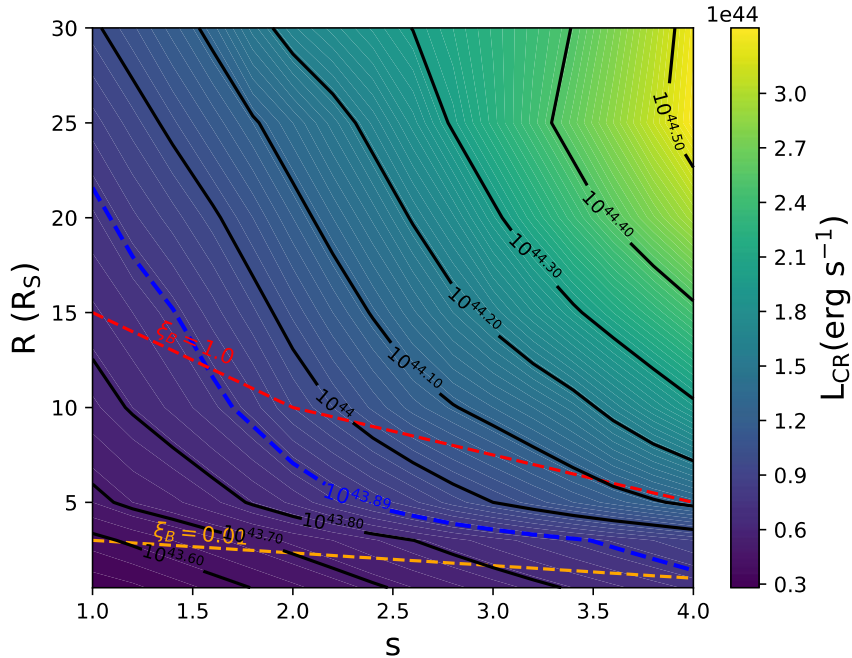


Constraint on emission radius: $R < (30 - 100) R_S$

Hadronic scenario results

$$\beta = \frac{8\pi n_p k_B T_p}{B^2} \approx \frac{\tau_T G M_{\text{BH}} m_p}{\sqrt{3} \zeta_e \sigma_T R^2 U_\gamma} \xi_B^{-1} \approx \left(\frac{\tau_T}{\sqrt{3} \zeta_e \lambda_{\text{Edd}}} \right) \xi_B^{-1}$$

$\tau_T \sim 0.1-1$ for X-ray corona, $\lambda_{\text{Edd}} \sim 0.5$
 $\xi_B > \sim 0.1$ leads to $\beta < \sim 1$



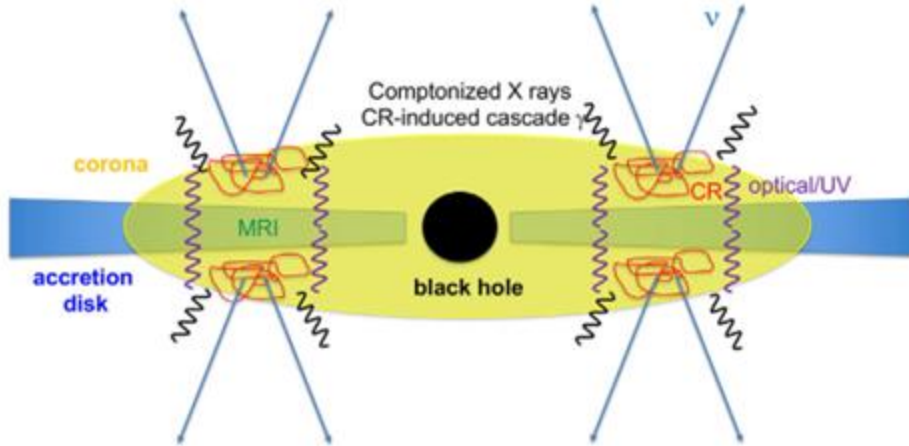
$R \leq 15 R_S$ for low- β plasma (synchrotron dominated)

$R \leq 3 R_S$ for high- β plasma (IC dominated, smaller than ISCO for non-rotating BH)

Emission region and acceleration mechanism

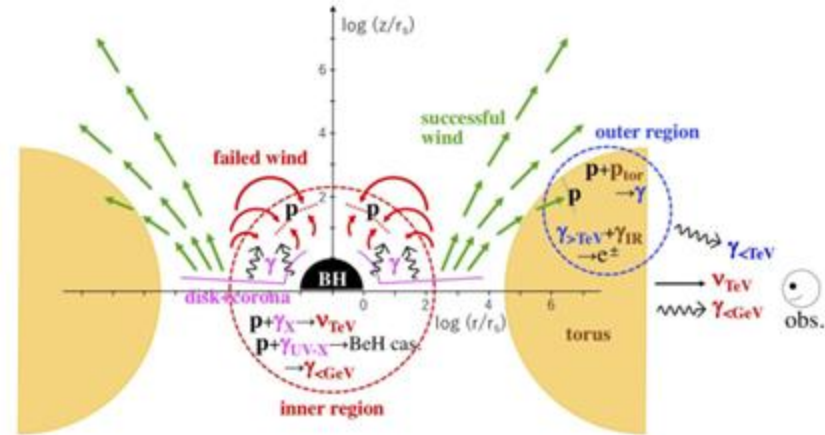
Magnetic corona model favored by low- β plasma

disk-corona



Murase+ (2020)

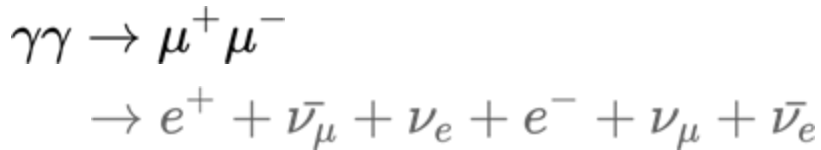
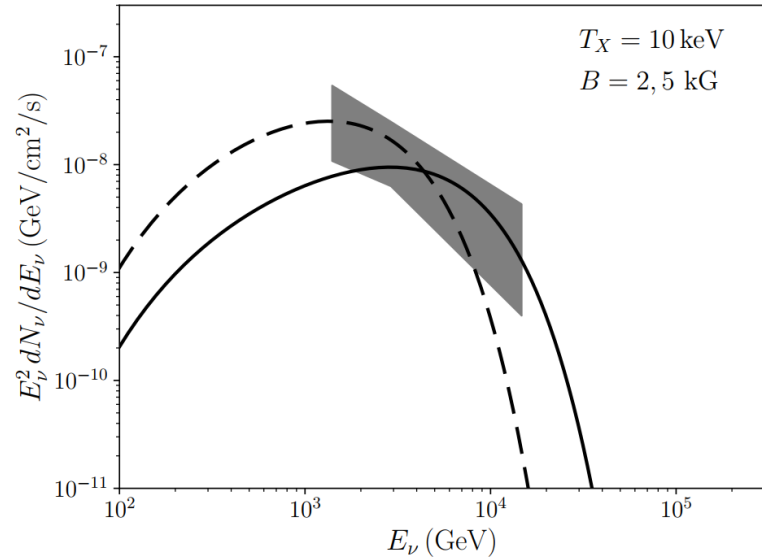
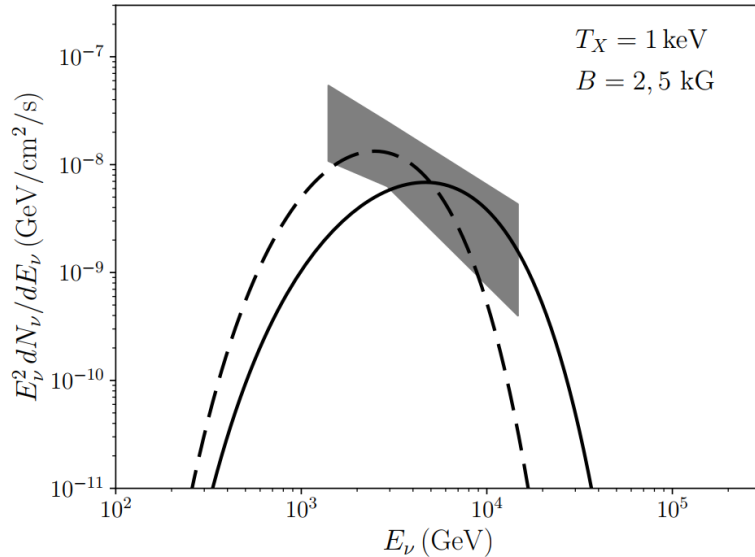
failed winds/ accretion shock



S. Inoue+ (2022)

Leptonic scenario

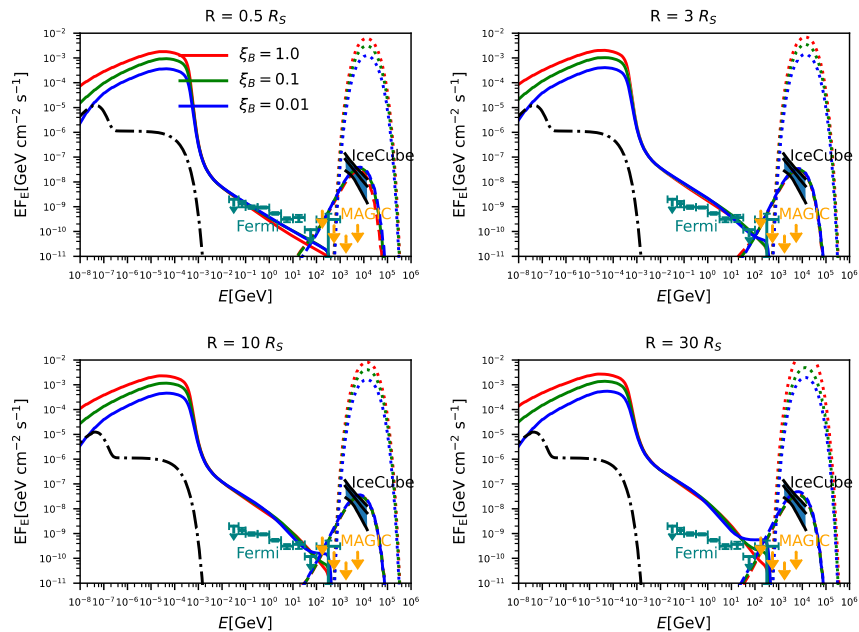
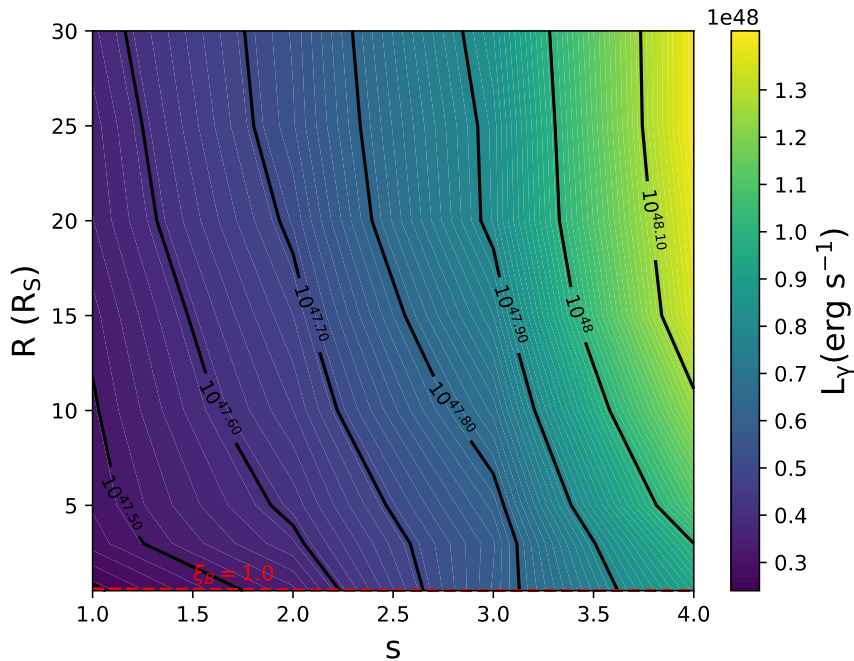
In principle, neutrinos can be produced by two-photon annihilation and muon-antimuon pair production. (Li & Waxman 2007 for neutrino production in intergalactic space)



Hooper & Plant (2023)

Leptonic scenario results

$$\frac{\tau_{\gamma\gamma \rightarrow e^+e^-}}{\tau_{\gamma\gamma \rightarrow \mu^+\mu^-}} \sim 2 \times 10^6 \left(\frac{\tilde{L}_{\text{disk}, 44.7}}{\tilde{L}_{\text{cor}, 43.3}} \right) \Lambda_1, \quad \varepsilon_\nu \frac{dL_{\varepsilon_\nu}}{d\varepsilon_\nu} \approx \frac{2}{3} \frac{\tau_{\gamma\gamma \rightarrow \mu^+\mu^-}}{\tau_{\gamma\gamma \rightarrow e^+e^-}} \varepsilon_\gamma \frac{dL_{\varepsilon_\gamma}}{d\varepsilon_\gamma}$$

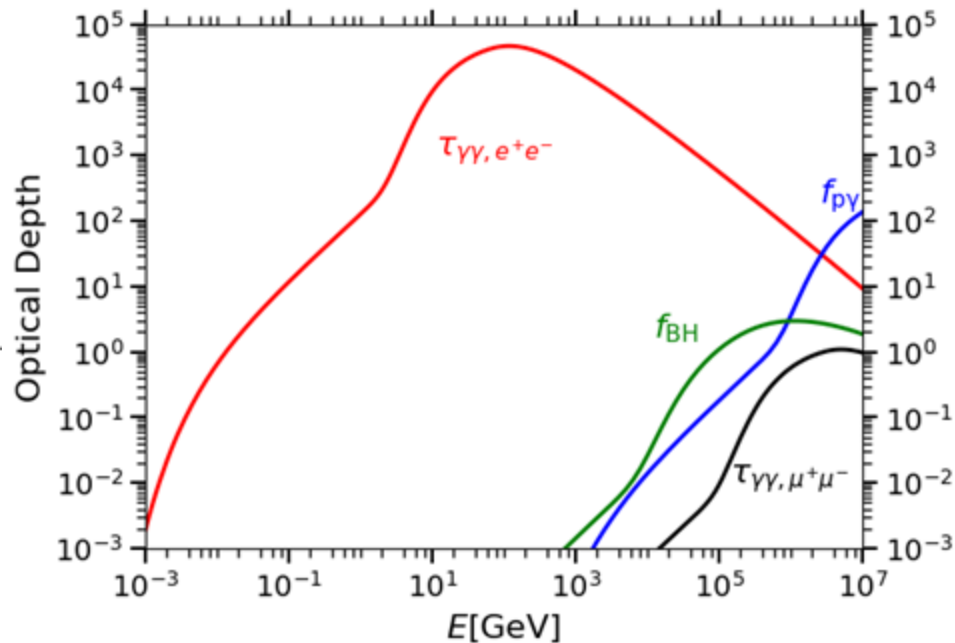


No viable parameter space found for physically motivated SED and magnetic field strengths.

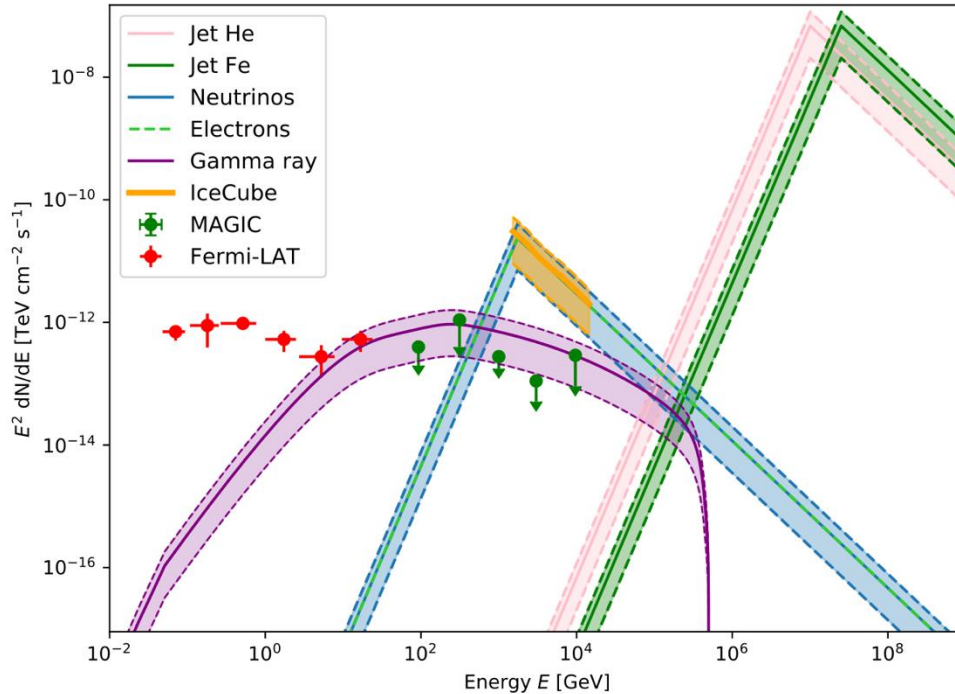
Optical Depth calculations

$$\begin{aligned}\epsilon_{\gamma\gamma, \hat{e}^+e^-} &= m_e^2 c^4 / \epsilon_0 \\ &\simeq 0.13 \text{GeV} (\epsilon_0 / 2 \text{keV})^{-1}\end{aligned}$$

$$\begin{aligned}\epsilon_{\gamma\gamma, \hat{\mu}^+\mu^-} &= m_\mu^2 c^4 / \epsilon_0 \\ &\simeq 5.6 \times 10^3 \text{GeV} (\epsilon_0 / 2 \text{keV})^{-1}\end{aligned}$$



Beta Decay scenario

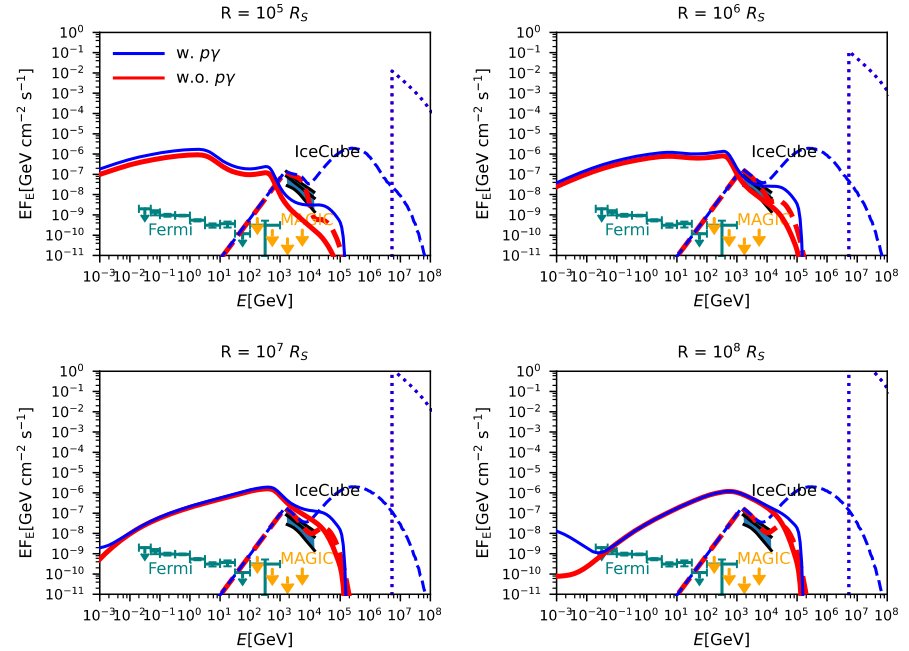
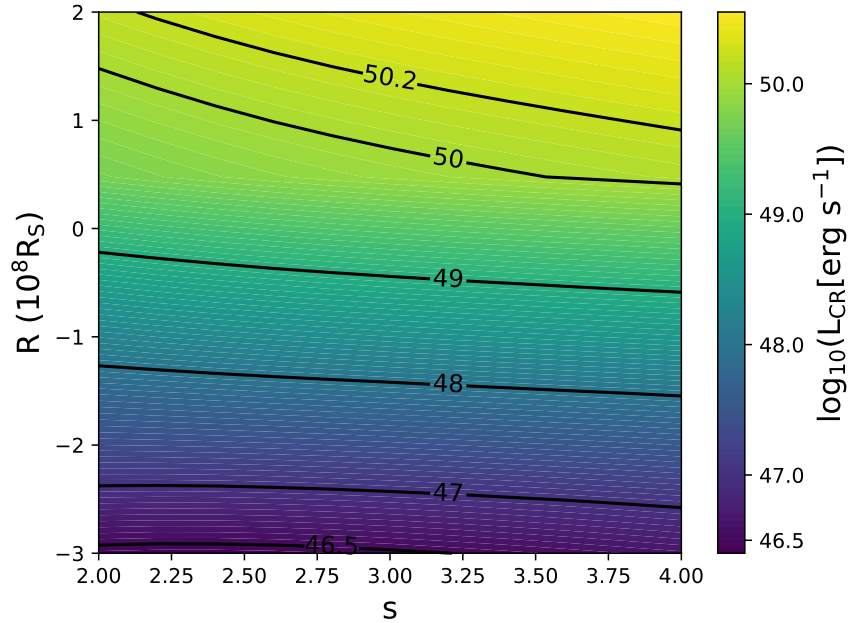


Photodisintegration of heavy nuclei accelerated by the jets produces neutrons, which undergo beta decay to form neutrinos.

Electrons from beta decay undergo IC scattering with CMB and EBL photons to produce gamma rays.

Emission region is located outside the dust torus if nuclei are accelerated in jets (Pe'er et al. 2009).

Beta Decay scenario results



No viable parameter space found upon considering Bethe-Heitler pair production (threshold for electron-positron production is lower than that for nuclear decay).

Summary & Conclusions

- Neutrino production from AGN strongly implies the acceleration of cosmic ray protons. Exotic models of neutrino production do not work if other relevant processes are included (violates limits set by EM observations).
- X-ray and gamma-ray observations set strong limits on the emission region. $R \leq 15 R_S$ for low- β plasma and $R \leq 3 R_S$ for high- β plasma (smaller than ISCO for non-spinning black hole).
- Emission region must be very close to the SMBH (magnetized corona). Disk-corona model (stochastic acceleration) is favored while failed line-driven wind model (shock acceleration) is ruled out by energetics.

Thank you!