



Modulation of Neutrino-Induced Radio Signals by Evolving Polar Ice

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TeV Particle Astrophysics - TeVPA
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UHEN Detection in Ice

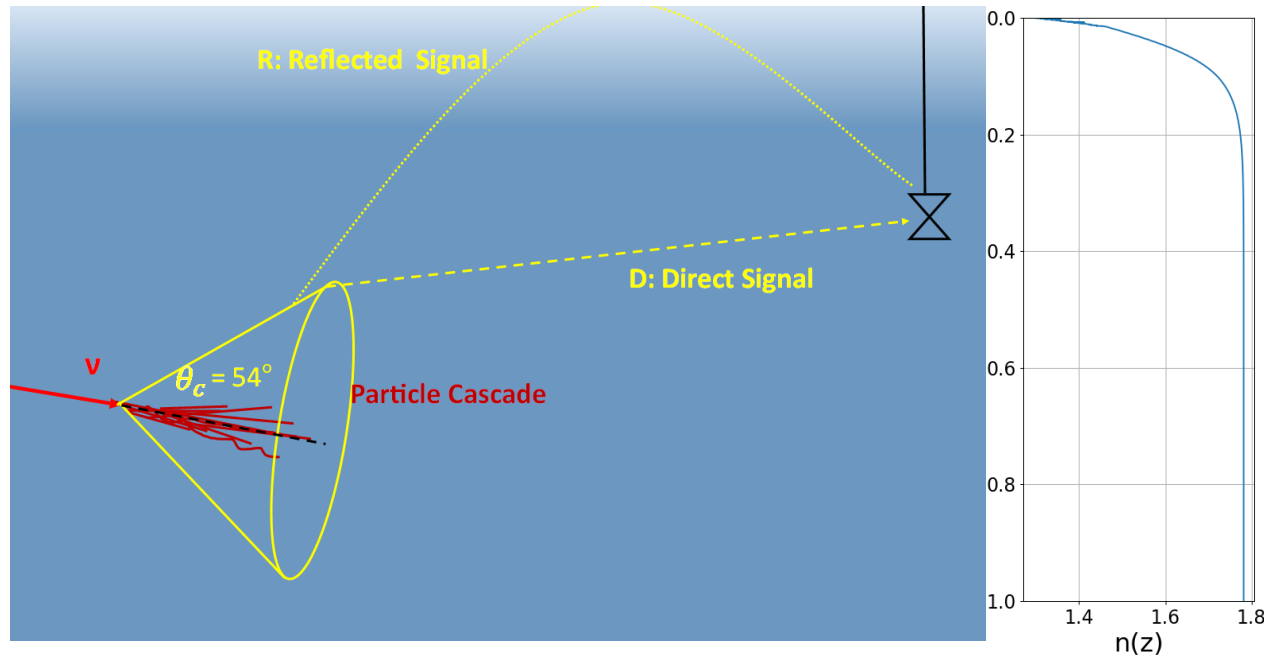
Polar ice sheets are ideal detection media for ultra-high energy neutrinos (**UHE- ν**) $E_\nu > 10$ PeV

- **Transparent to radio waves in the MHz – GHz regime** - attenuation length $L_\alpha = O(1 \text{ km})$ at 100 MHz to 500 MHz
- Allows for large volumes of ice to be instrumented
- Antennas are also far cheaper than optical modules

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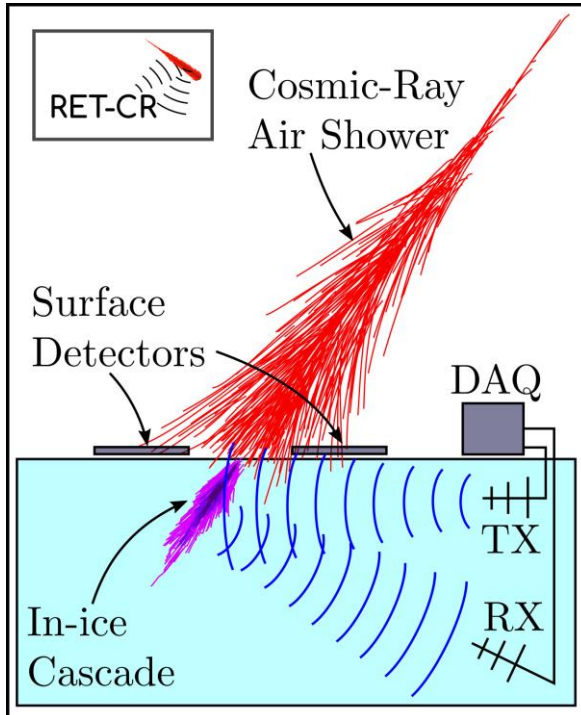
Askaryan emission

- Build-up of excess negative charge in the medium
- Coherent radio emission produced at the Cherenkov angle θ_c to the cascade direction defined by the refractive index ($\theta_c = 56$)
- Experiments searching for in-ice Askaryan: RICE, ARA, RNO-G, ARIANNA, ANITA, PUEO, IceCube-Gen2 Radio

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Radar-Echo method:

- Ionization trail persists briefly after UHE- ν cascade \rightarrow acts as a reflective object for in-ice radar.
- Method is currently being tested in the field using UHE cosmic rays: Radar Echo Telescope for Cosmic Rays (RET-CR)
- Verification of the method in nature will facilitate the development of a future RET-N
- See Dylan Frikken's talk - HE Astro / Gravitational Waves II (08/26)

Firn Layer

RET-CR & In-Ice Askaryan Detectors have antennas within or immediately below the ‘firn layer’

Firn: intermediate stage between fresh fallen snow and glacial ice

- 100 – 150 m deep in Greenland and Antarctic ice caps
- Assuming constant accumulation and temperature:

$$\rho(z) = \rho_i + \Delta\rho e^{-kz}$$

- Densification rate k changes due to different dominant process

- $0 < z < z_{550}$ – ‘Shallow Firn’
- $z_{550} < z < z_{800}$ – ‘Deep Firn’
- $z > z_{800}$ – ‘Glacial Ice’

- The refractive index was found empirically to be linearly proportional to the firn density [Kovacs et al. 1994]:

$$n = 1 + A(\rho[g/cm^3])$$

Snow: settling and grain growth
200 kg m⁻³

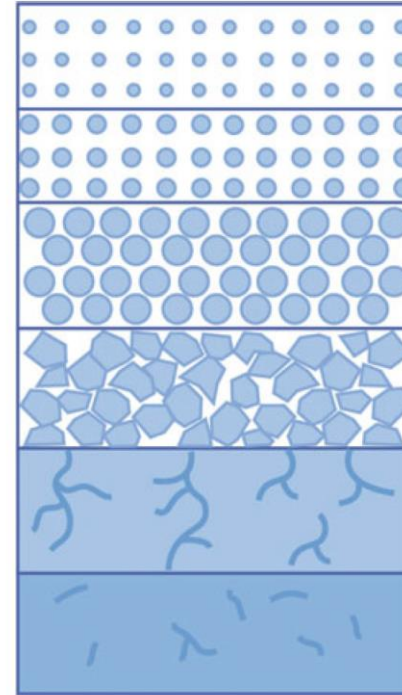
Firn: settling and grain growth
400 kg m⁻³

Firn: settling and grain growth
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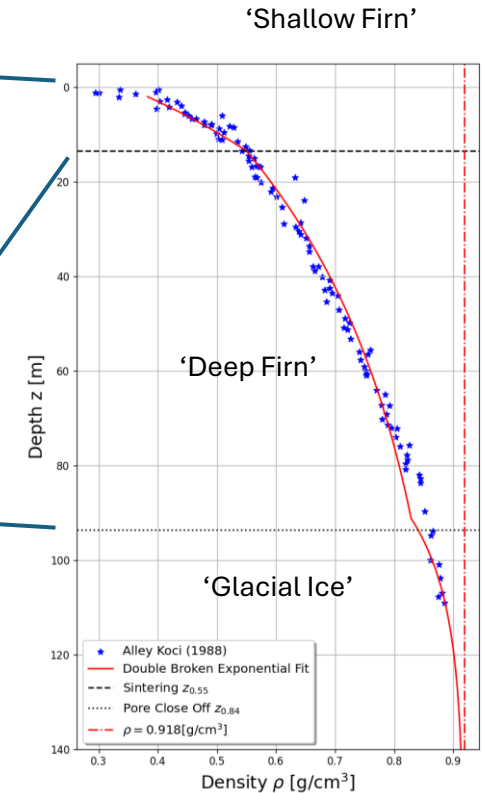
Sintering
800 kg m⁻³

Pore close-off
830 kg m⁻³

Glacier ice
900 kg m⁻³



Buizert & Helsen, *Glaciers and Ice Sheets in the Climate System*, Chapter 11 (Springer)



Firn density profile at Site A, Greenland (1988) – Alley & Koci (1988) – *Annals of Glaciology*.

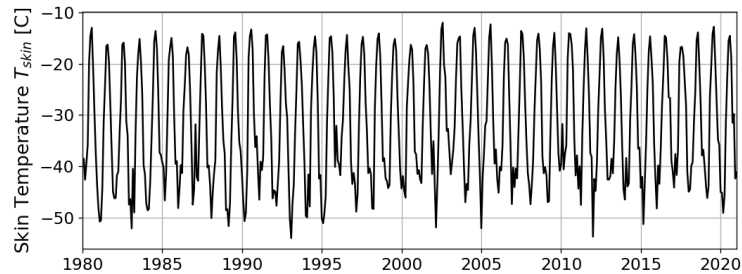
Firn Variation and Evolution

This simple description of Firn is complicated by seasonal fluctuations in temperature:

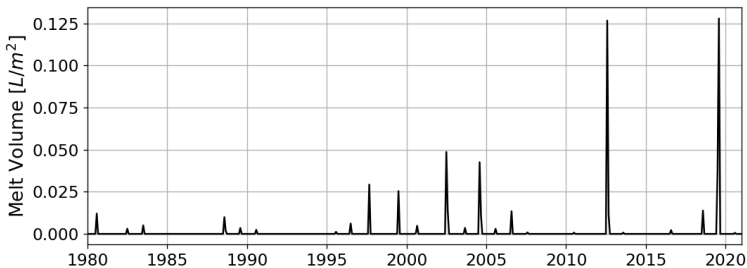
- Temperature variation leads to fluctuations in firn density
- Episodic surface melting events & rainfall

These lead to the formation of refrozen ice layers & density fluctuations → Strongest in the shallow firn layer ($z < 15$ m at Summit)

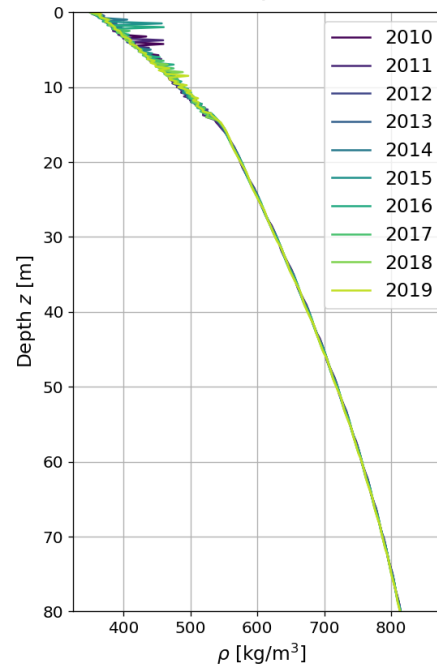
Monthly surface air temperature at Summit, Greenland



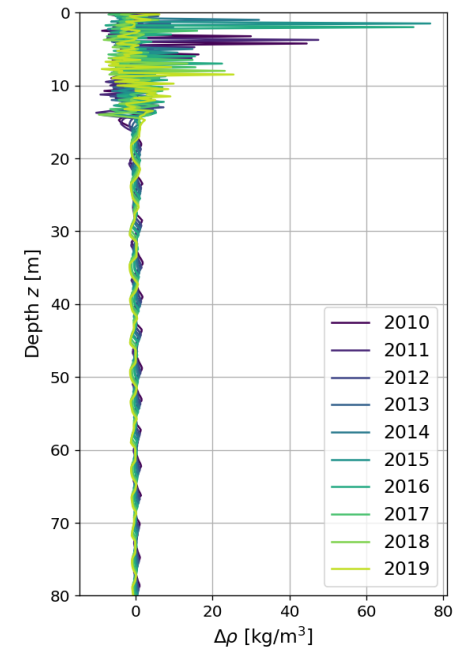
Surface Melt Events



Predicted Density from Community Firn Model



Variation of Density



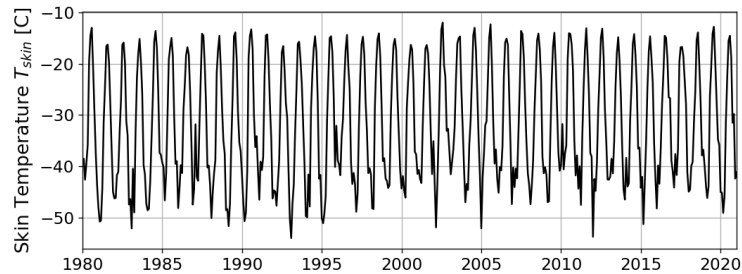
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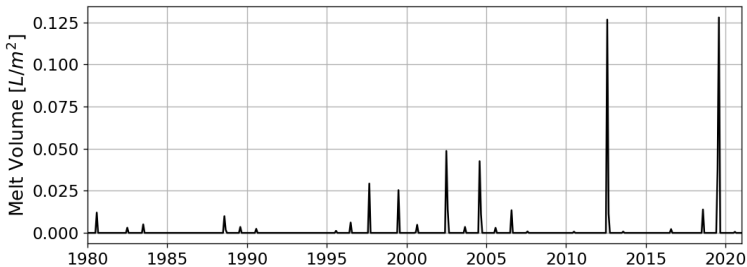
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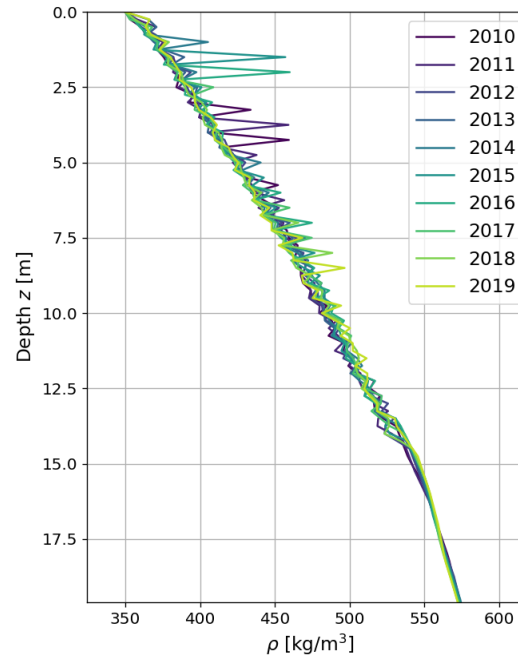
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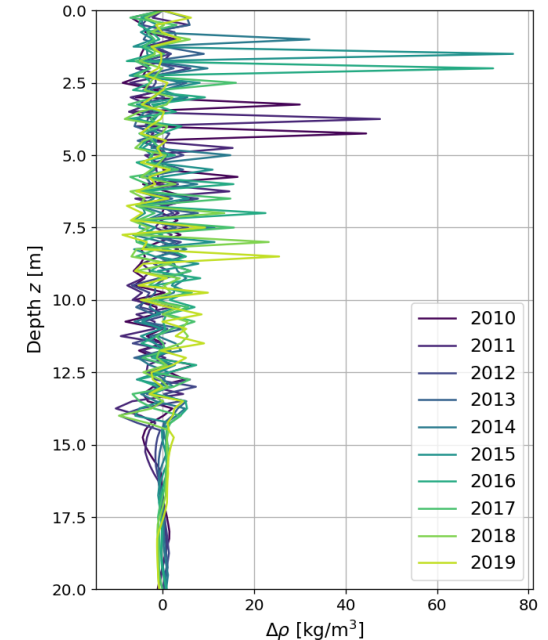
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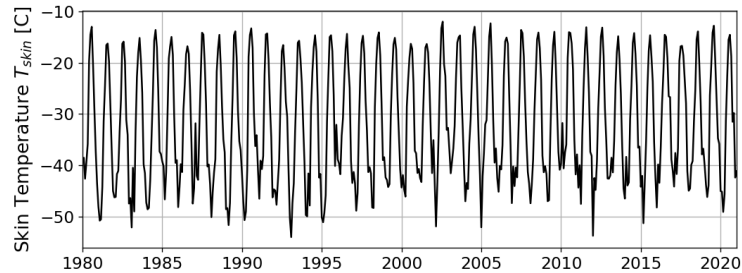
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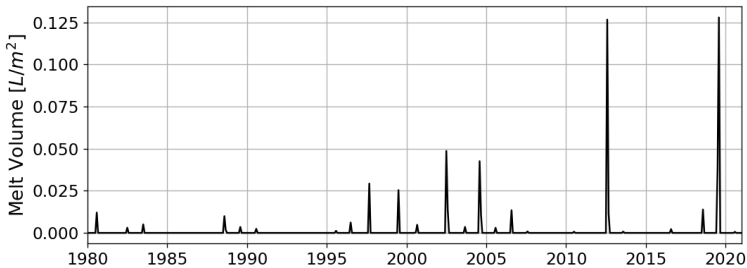
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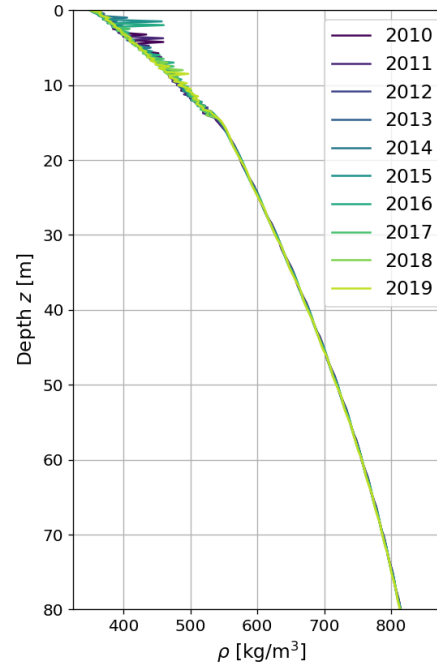
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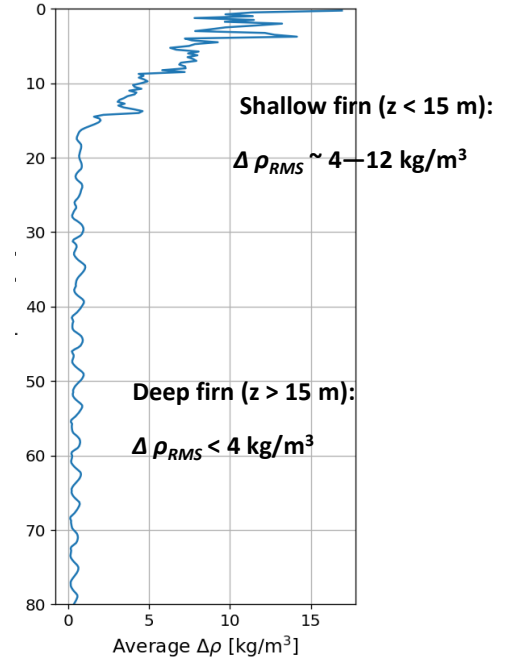
Surface Melt Events



Predicted Density from Community Firn Model



Average Variation of Density [2010-2019]



Simulation Study: Summit, Greenland

Variable firn density over time may change the properties of neutrino-induced radio signals
To investigate this effect, we use the firn layer at Summit, Greenland as a case study.

Goal: Simulate radio propagation from a deep source to a set of receivers at depths $0 \text{ m} < z < 200 \text{ m}$, for a geometry of 1km depth x 1km radius (comparable to the attenuation length!)
→ Quantify variation in out E-field trace at RX - $E_{RX}(t)$ for $n(z)$ profiles from different times

Analysis parameters:

- Peak amplitude: $E_{RX,max}$ for direct (D) and reflected/refracted (R) signal
- Time delay between D and R signal: $\Delta t_{DR} = t_R - t_D$
- Fluence: $\phi_{RX}^E = \epsilon c \int E_{RX}^2(t) dt$

Radio Simulation Codes (see backup slides for more details):

- **MEEP**: Direct solution of Maxwell's equations in a geometric grid (FDTD method)
 - Most accurate method but computationally expensive (requires grid resolution $\Delta x \leq \lambda/10$)
- **paraProp**: Parabolic-wave approximation of Maxwell's equation within cylindrical symmetry volume
 - Accurate within 'paraxial angle' to the horizontal direction → computationally efficient



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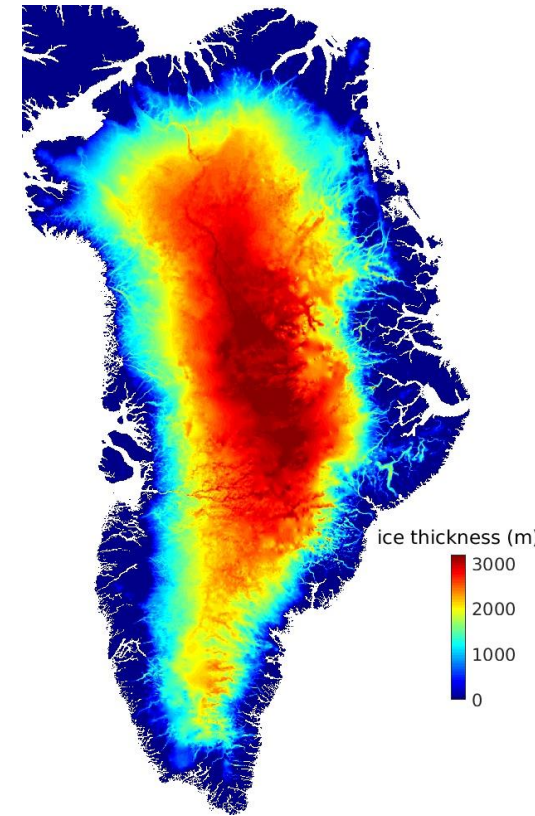
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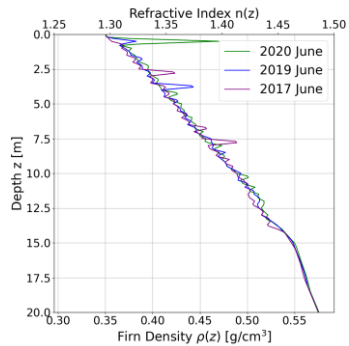


Simulation Chain

Refractive Index:

Community Firn Model (CFM)

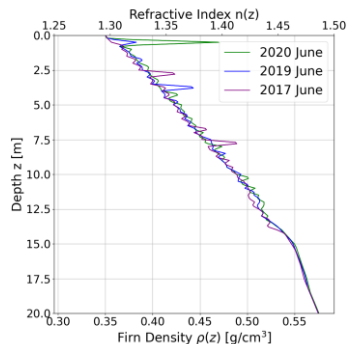
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- Depth grid size: 10 cm
- Outputs from each month from 1980 to 2020



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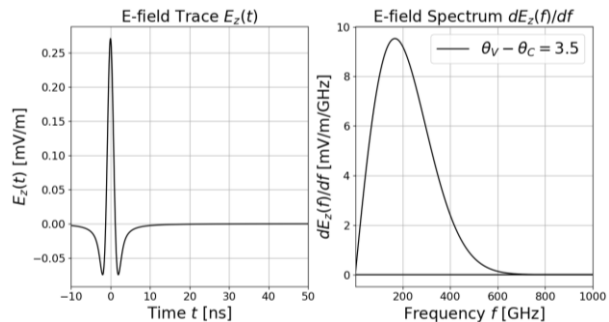
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Source signal:

Analytical model of Askaryan pulse from 10^{18} eV hadronic shower

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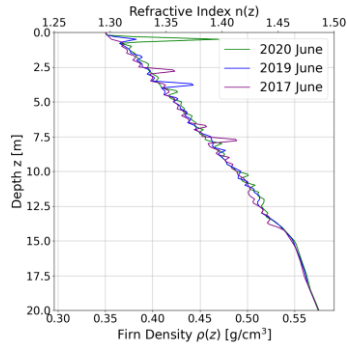


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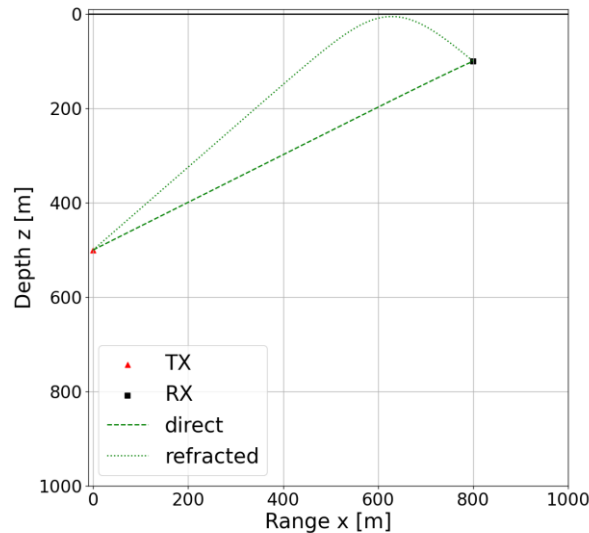
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Radio Propagation Simulation

paraProp & MEEP

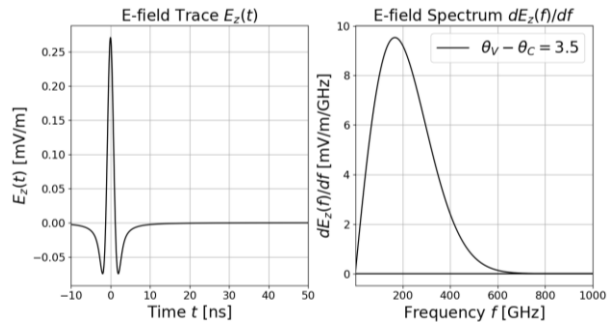
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- Propagate Askaryan pulse through throughout geometry
- 1000 m radius x 1000 m depth



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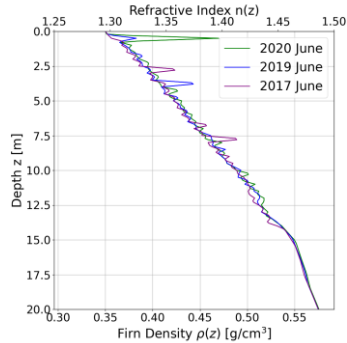


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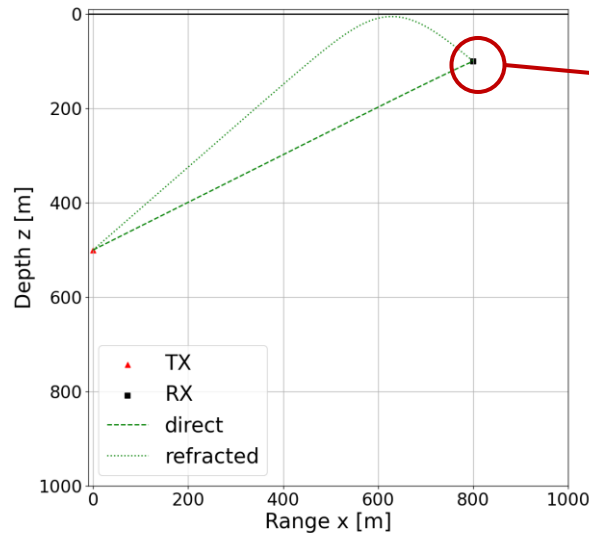
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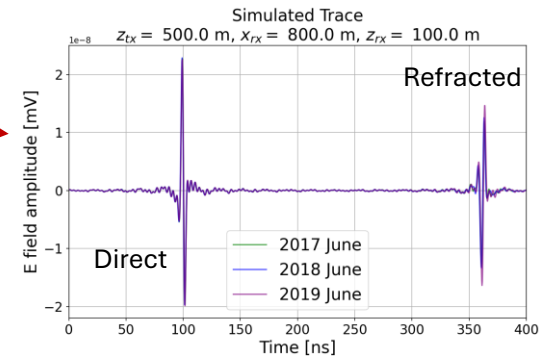
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Output

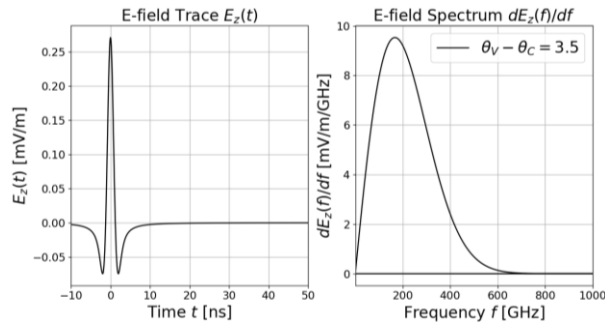
- Sample radio signal at a set of receivers (RX)
- Compare outputs for ref-index at different times
- Is possible to add antenna & detector response functions + noise



Source signal:

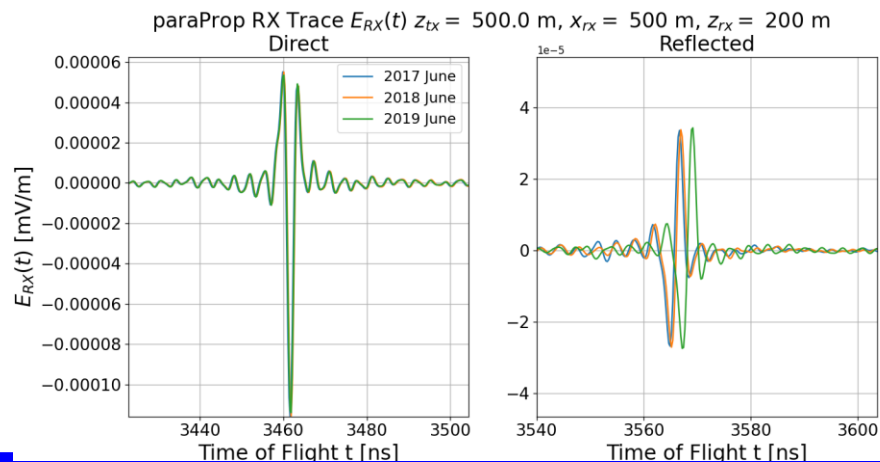
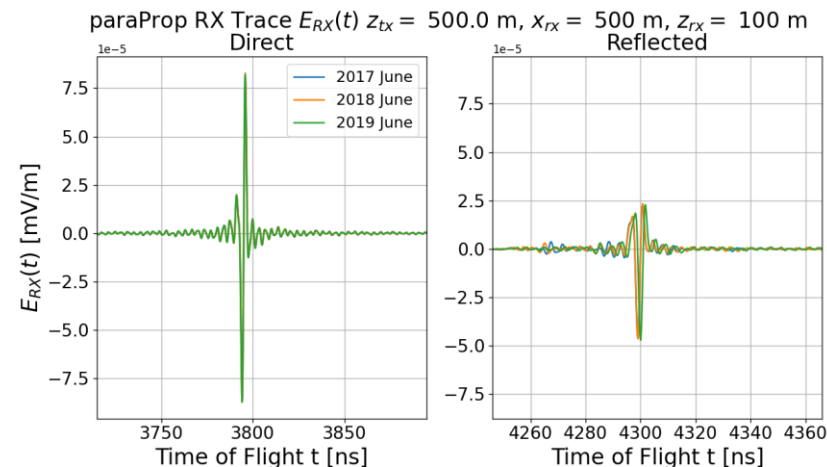
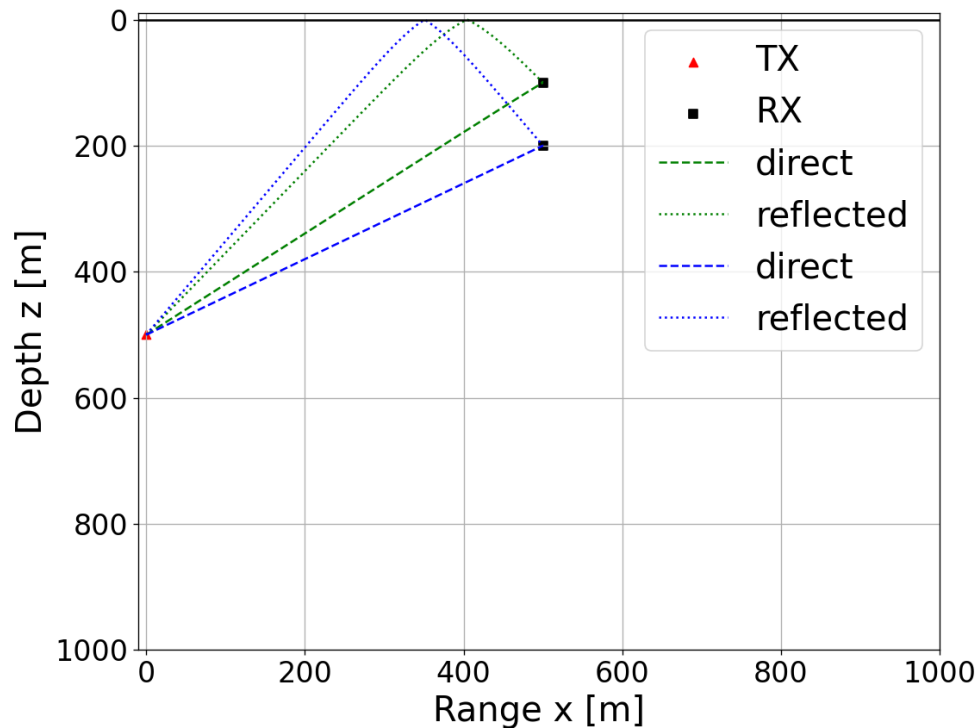
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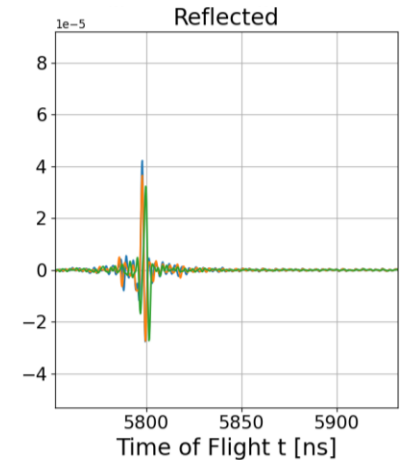
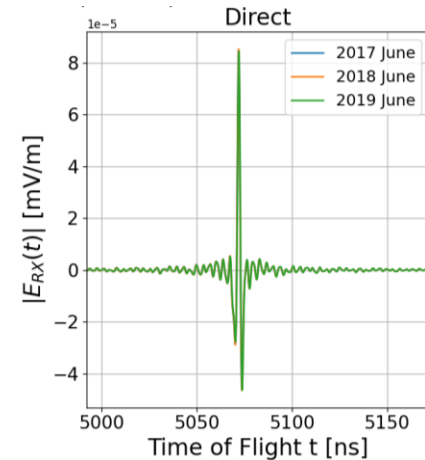
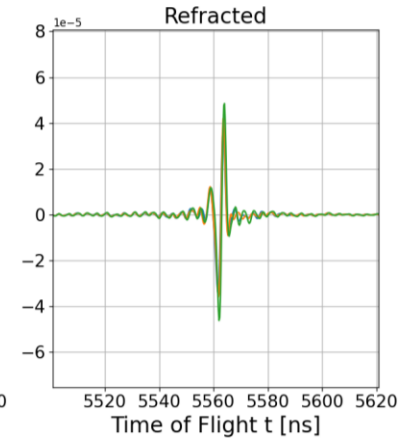
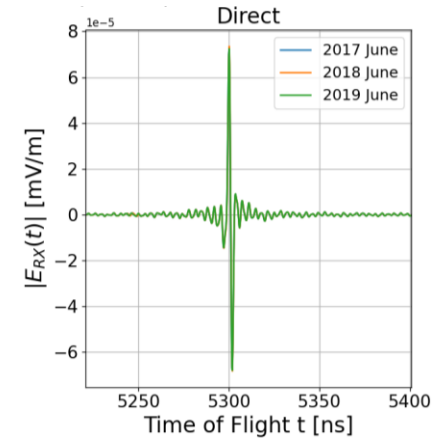
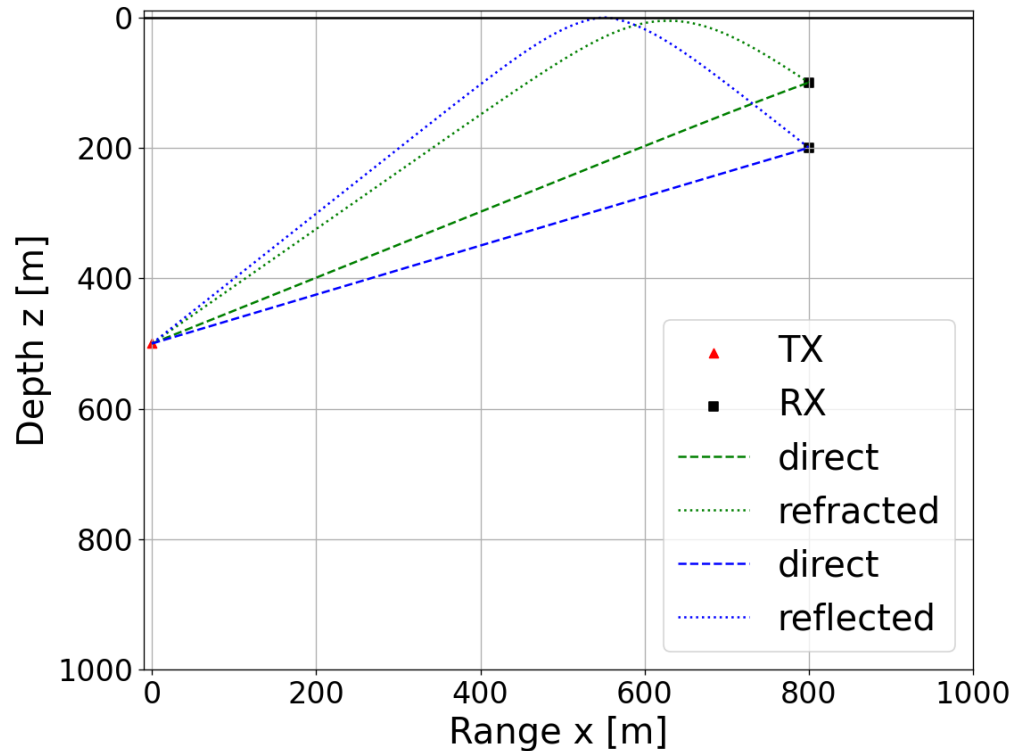
paraProp Results ($z_{TX} = 500$ m)

Paths of direct, reflected, and refracted signals from 500 m source to RX at 100 m and 200 m (calculated with NuRadioMC)
 $|\theta_v - \theta_c| = 5$



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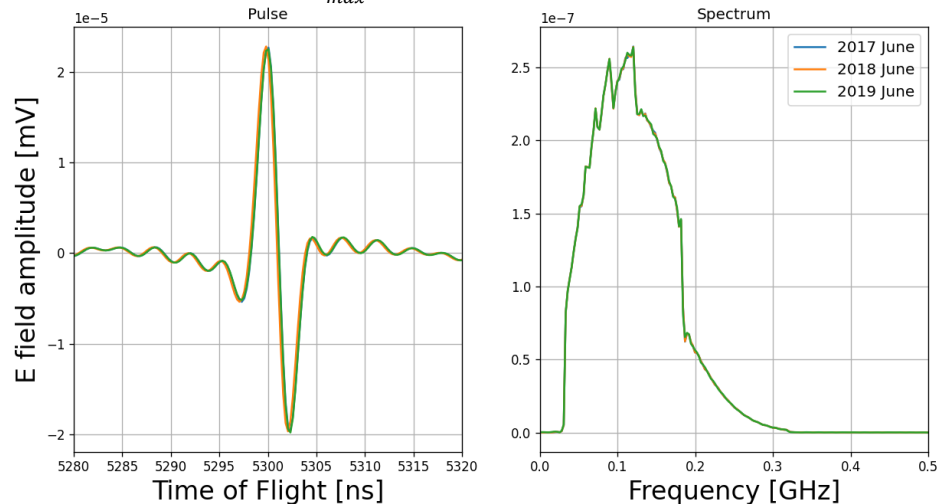
paraProp Results ($z_{TX} = 500$ m)

Only the reflected/refracted paths traverse through the shallow firn layer

- Amplitude modulation of the R-signal is an order of magnitude higher than for D-signals
- No significant variation of Δt_{DR}

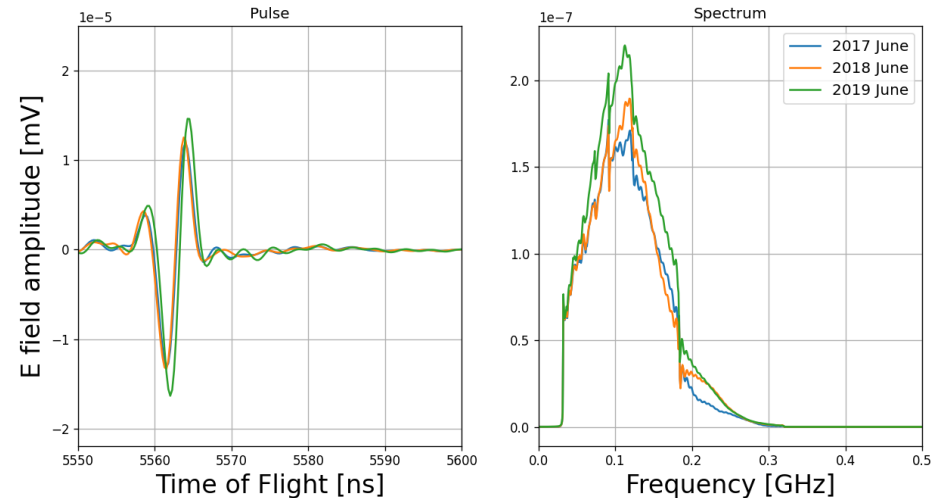
Direct signal trace at $z_{RX} = 100$ m

- $\frac{\Delta E_{max}}{E_{max}} < \sim 10^{-3}$

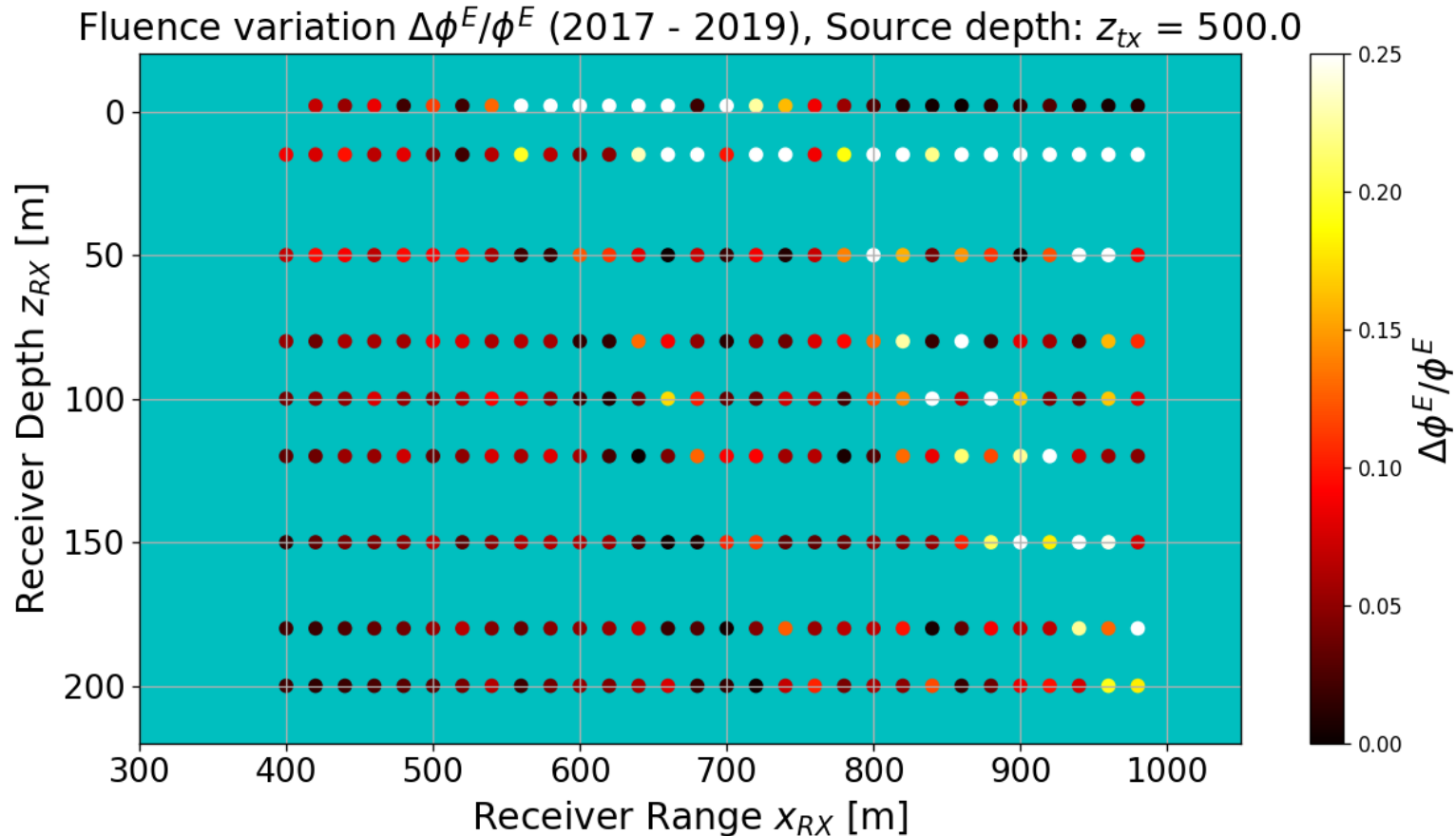


Refracted trace at $z_{RX} = 100$ m

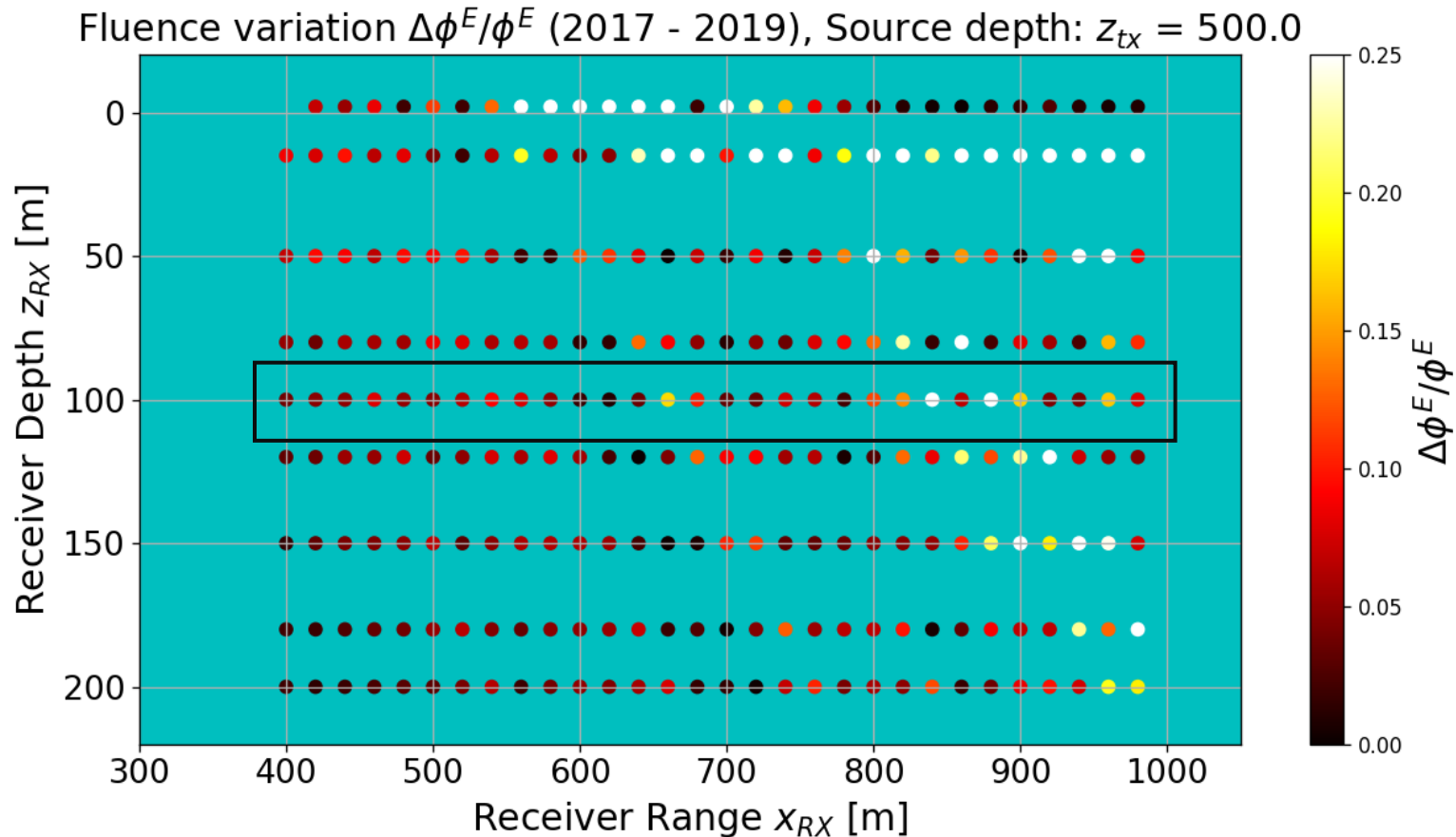
- $\frac{\Delta E_{max}}{E_{max}} < \sim 10^{-1}$



Variation of Fluence (2017 vs 2019)



Variation of Fluence (2017 vs 2019)



Variance in RF parameters

Electric field fluence Φ^E is proportional to the UHE- ν energy E_ν

Time frame: 2015-2020, $|\theta_\nu - \theta_c| = 5$

paraProp:

- Reflected: $\frac{\Delta\Phi^E_R}{\Phi^E_R} > \sim 0.1$ ($x > 600$ m)

- Direct: $\sim 10^{-3} < \frac{\Delta\Phi^E_D}{\Phi^E_D} < \sim 10^{-2}$

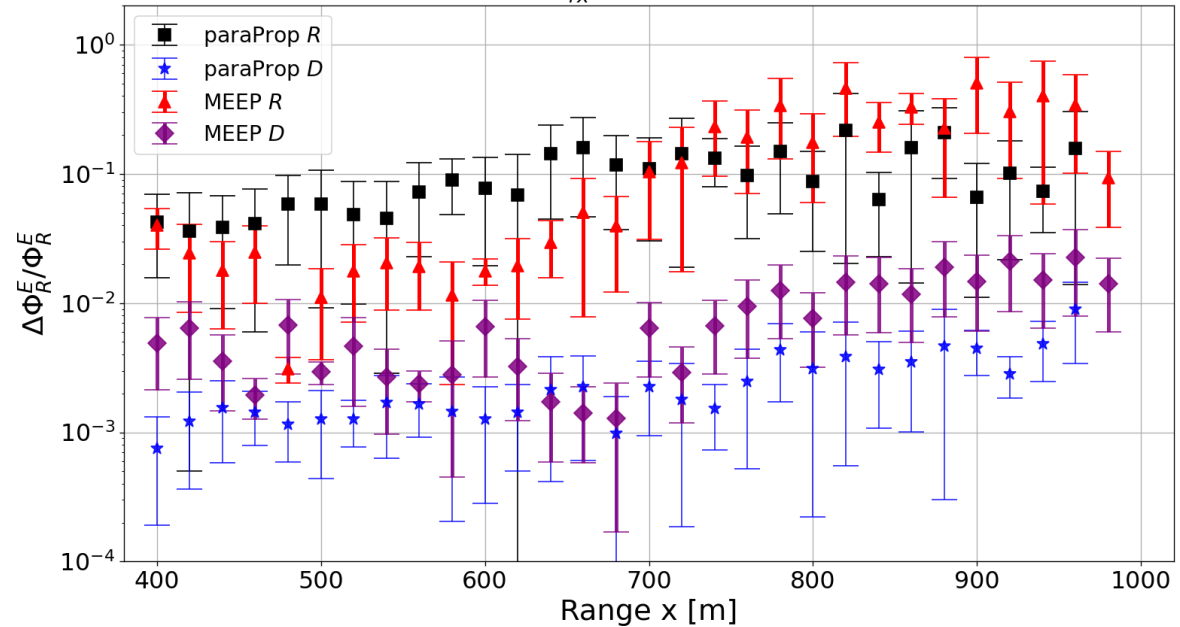
Meep:

- Reflected: $\frac{\Delta\Phi^E_R}{\Phi^E_R} > \sim 0.1$ ($x > 700$ m)

- Direct: $\sim 10^{-2} < \frac{\Delta\Phi^E_D}{\Phi^E_D} < \sim 10^{-1}$ ($x > 800$ m)

Fluence Residuals $\Delta\Phi^E/\Phi^E$ paraProp & MEEP

$Z_{rx} = 100.0$ m

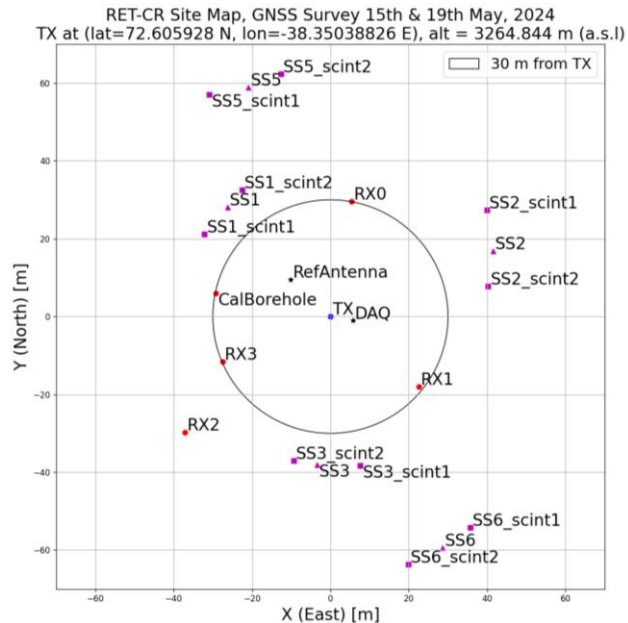


Preliminary finding: Under likely signal geometries, shallow firm fluctuations produce systematic uncertainty in the fluence of reflected and refracted signals

→ Results in systematic error for neutrino energy and arrival direction reconstruction.

Measuring firn properties at Summit

- During the May 2024 RET-CR deployment, boreholes were made using a coring drill
- Firn cores were extracted down to a depth of 13 m → used to measure the refractive index profile at the RET site in two ways:
 1. Direct gravimetric measurement of core density ρ
 2. Open-ended coaxial probe in contact with firn → measure the relative permittivity ϵ_r from the reflected radio energy (method described in backup slides!)



Measuring firm properties at Summit

Measurement procedure

1. Slide core through shaping-rig
2. Press firm core against coaxial probe
3. Measure dielectric properties using coaxial probe method
4. Cut core – adjustable length (usually 10 cm)
5. Measure & log the core segment weight
6. Repeat

Firn Segment

Scales



Shaping Rig

Coaxial Probe



Results: Permittivity & Density

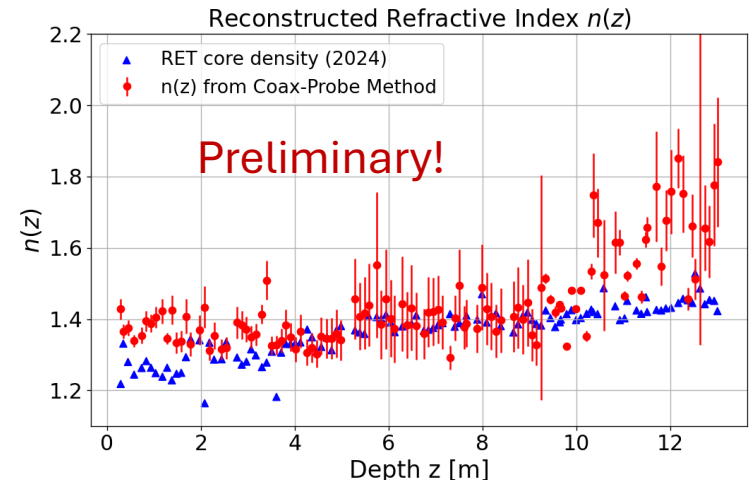
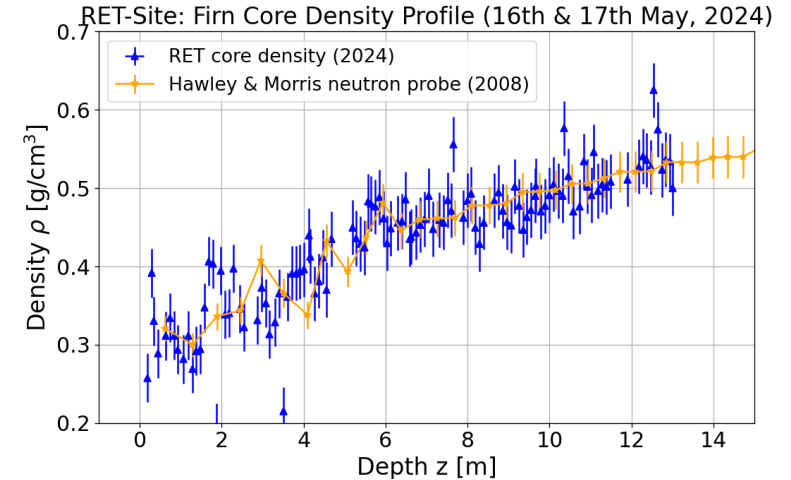
Reconstructed Density:

- Refractive index and density related empirically $n = 1 + 0.845 \rho$
- Firn density was broadly consistent with previous measurements made at Summit
 - Hawley & Morris estimated the firn density using neutron scattering measurements in 2008 to 30 m depth
- Evidence of ice layers at $z = 7.65, 10.4, \& 12.5$ m

Reconstructed permittivity (coax probe method): $n = \sqrt{\epsilon_r}$

Analysis ongoing!

- Averaged between 500 MHz and 800 MHz
- Large uncertainty due to frequency variation in the reconstruction (described in backup slides) \rightarrow not yet explained
 - High conductivity of the firn \rightarrow frequency-dependent permittivity reconstruction
 - Conductivity of test/calibration materials uncertain
- Correlation between density and permittivity with significant variations



Seasonal modulation of radio signals due to ice properties:

- ***The polar regions are warming rapidly:*** understanding firn evolution and its modulation of Askaryan and Radar-Echo signals will be important for UHE- ν searches
- Under likely signal geometries, shallow firn fluctuations produce systematic uncertainty in the fluence of R-signals, and possibly for neutrino energy reconstruction as well

In-situ measurements of ice properties

- The density profile at the Summit site is broadly consistent with previous measurements – with evidence of recent melting events
- Correlation of reconstructed permittivity with density measurements
 - Caution about results: calibration likely incomplete, frequency dependence likely unphysical

Future work:

- Simulating larger geometries – 3 km (depth) x 3km (radius)
- Analysis of TX to RX radio propagation at RET site \rightarrow further insight into ice properties
- Examine seasonal radio modulation at South Pole, Antarctica

Thanks for your attention!

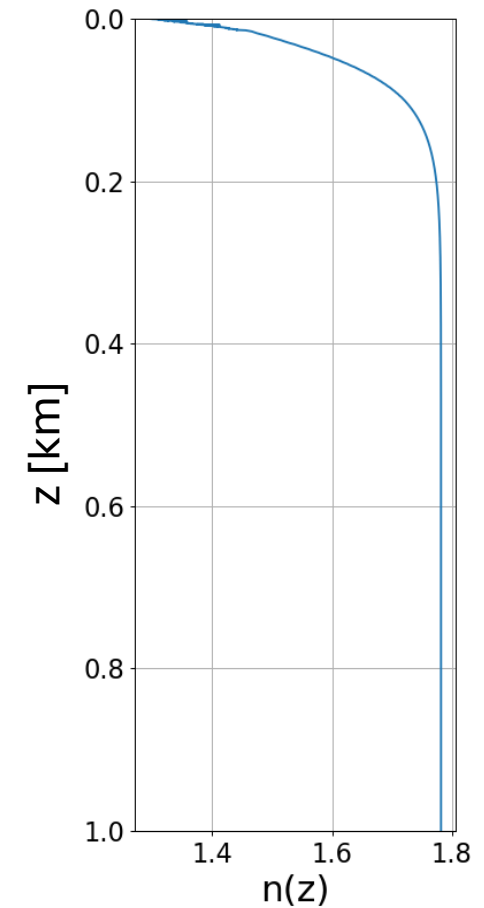
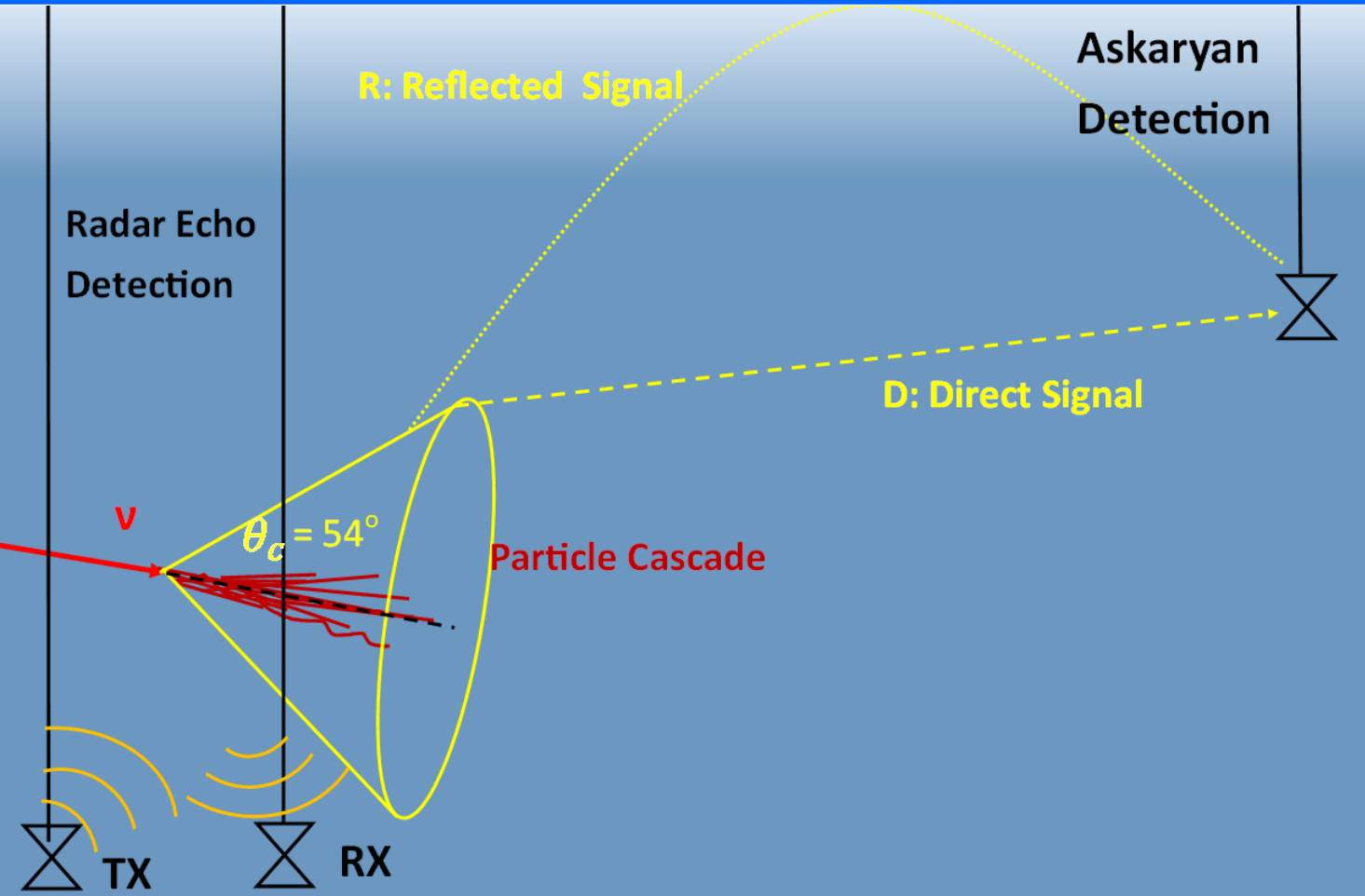


RET May 2024 Deployment Team

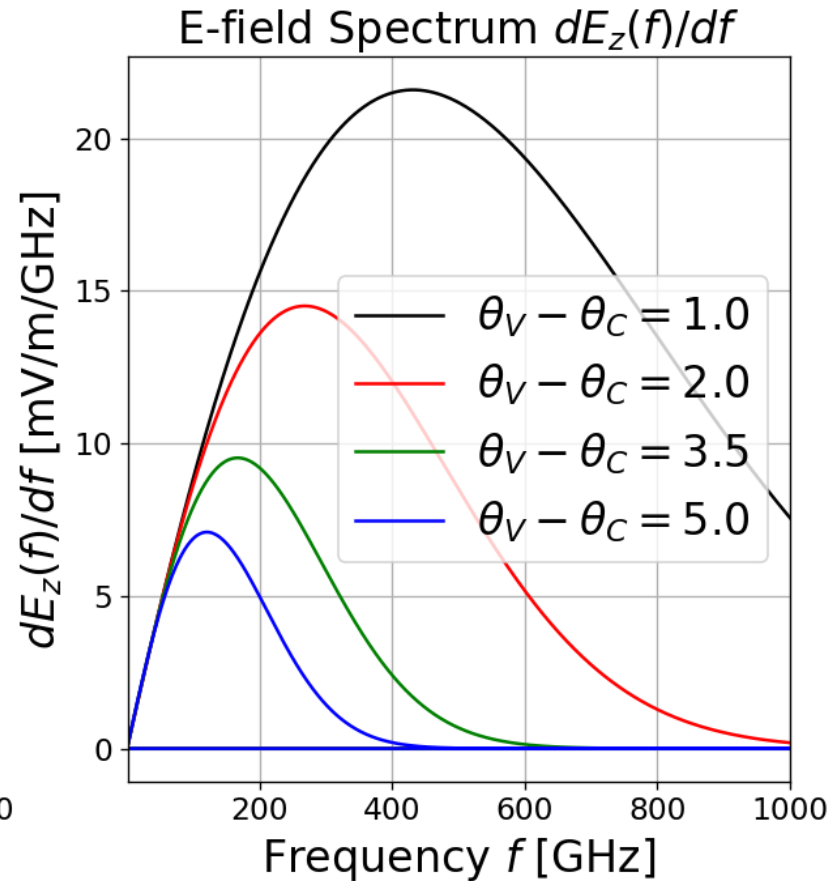
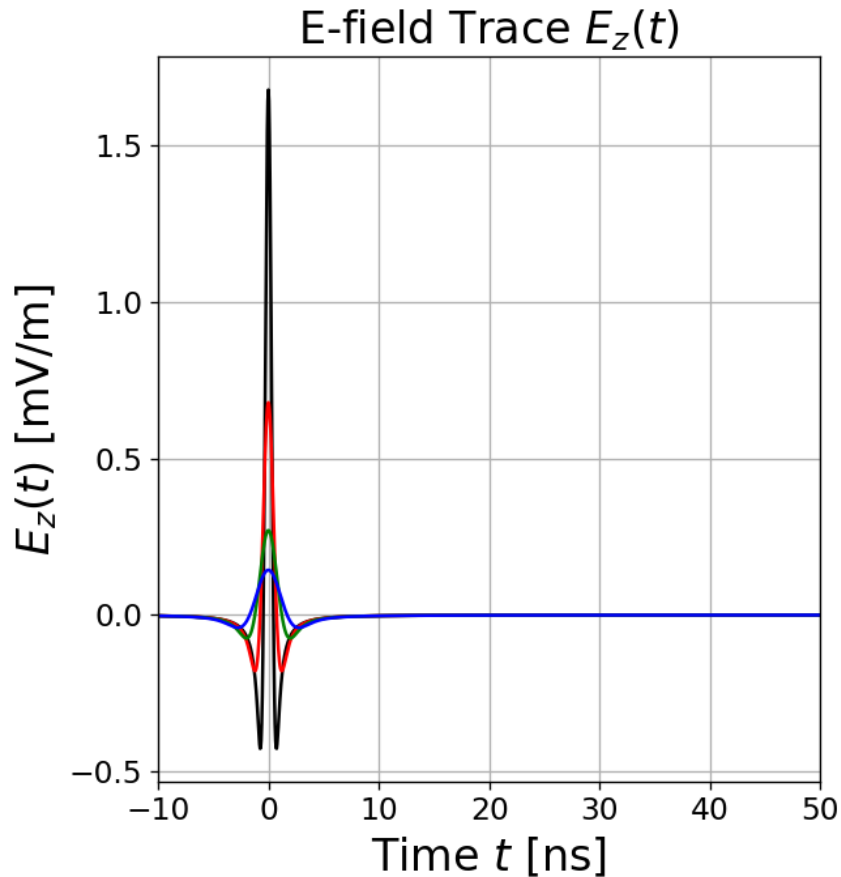
Backup Slides



UHEN Detection in Ice



Askaryan Signal



Challenge: Large domain size & need for high spatial resolution to accurately resolve high frequencies -> Cylindrically symmetric medium 1 km radius x 1km depth, frequency range 1 MHz < f < 1000 MHz

MEEP - Finite Difference Time Domain (FDTD) method

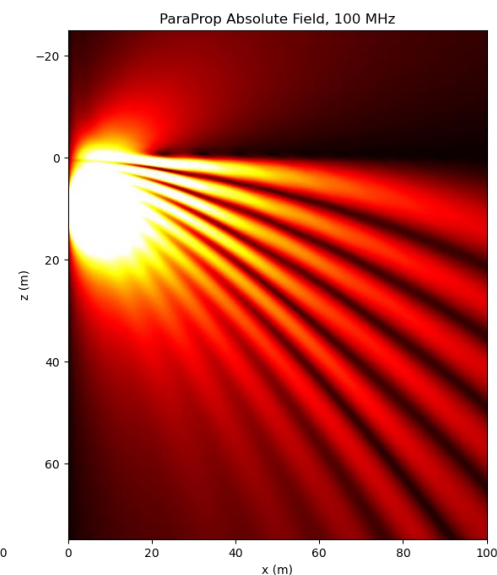
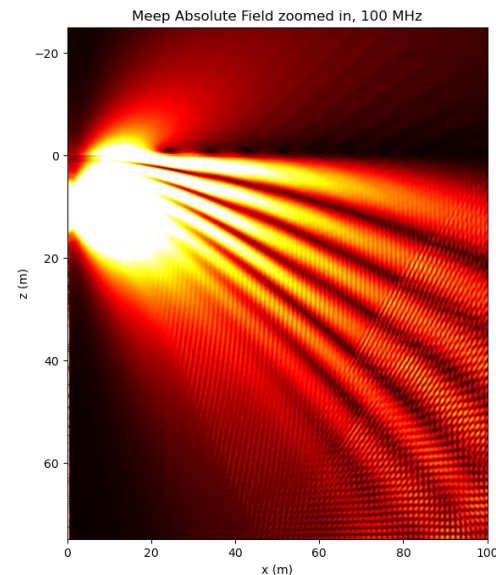
- Solves Maxwell's equations inside discrete cells
- Computationally expensive:
 - Cell size $\Delta x \leq \lambda/10$ i.e. 8 cm for f = 300 MHz in vacuum
 - Time resolution is related to cell size:

$$\Delta t \leq \frac{1}{c_{medium} \sqrt{(\Delta x^2 + \Delta y^2 + \Delta z^2)}}$$

- Simulations for f > 500 MHz are exceedingly expensive

paraProp - Parabolic Equation:

- Parabolic wave approximation of Maxwell's equations in a cylindrically symmetric medium
- Amplitude and phase residual errors low (<0.01) relative to FDTD within the 'paraxial angle' ($\theta < 45$ deg)
- Computationally efficient -> Tractable for volumes > 1 km and f > 1 GHz



Continuous wave RF emission at f = 100 MHz from a shallow dipole antenna – simulated with MEEP and paraProp

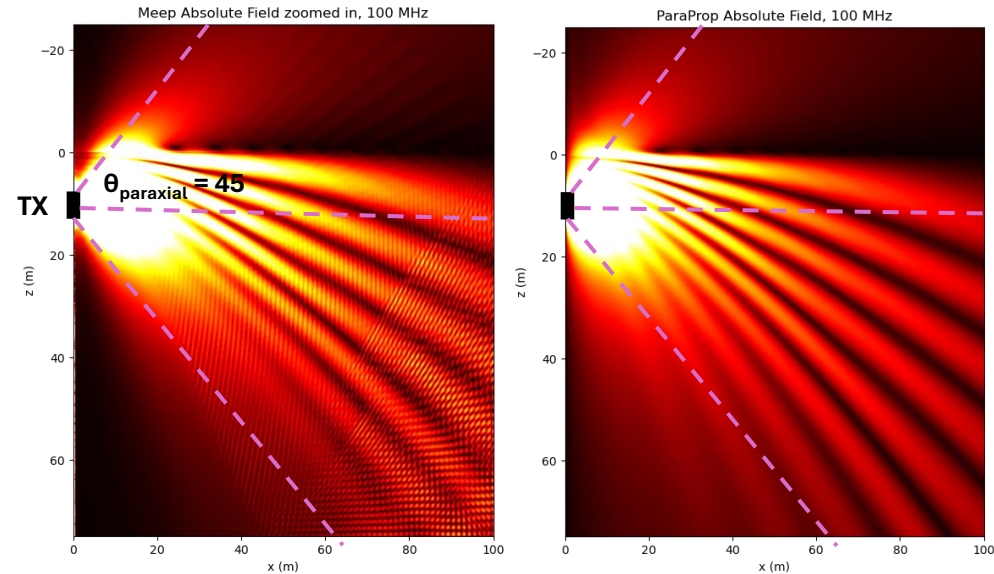
Challenge: Large domain size & need for high spatial resolution to accurately resolve high frequencies -> Cylindrically symmetric medium 1 km radius x 1km depth, frequency range 1 MHz < f < 1000 MHz

MEEP - Finite Difference Time Domain (FDTD) method

- Solves Maxwell's equations inside discrete cells
- Computationally expensive:
 - Cell size $\Delta x \leq \lambda/10$ i.e. 8 cm for f = 300 MHz in vacuum
 - Time resolution is related to cell size:
$$\Delta t \leq \frac{1}{c_{medium} \sqrt{(\Delta x^2 + \Delta y^2 + \Delta z^2)}}$$
 - Simulations for f > 500 MHz are exceedingly expensive

paraProp - Parabolic Equation:

- Parabolic wave approximation of Maxwell's equations in a cylindrically symmetric medium
- Amplitude and phase residual errors low (<0.01) relative to FDTD within the 'paraxial angle' ($\theta_{paraxial} < 45$ deg)
- Computationally efficient -> Tractable for volumes > 1 km and f > 1 GHz



Continuous wave RF emission at f = 100 MHz from a shallow dipole antenna – simulated with MEEP and paraProp

In-Situ Ice Property Measurements

Density Measurements:

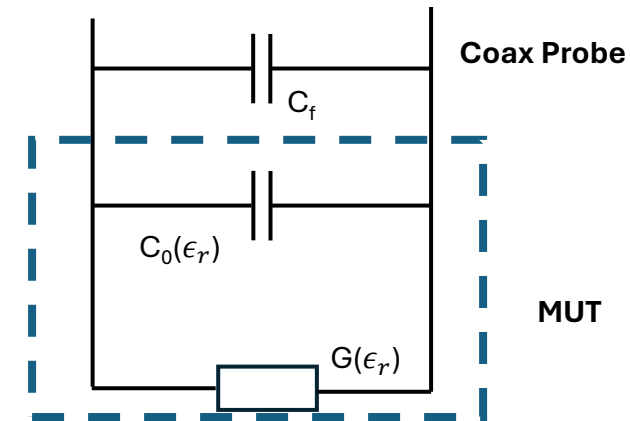
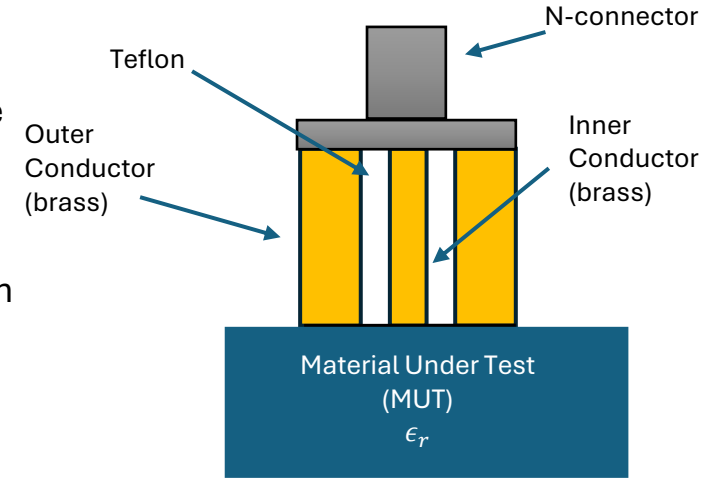
- Gravimetric: measure the mass of a firm sample with a definite volume
- Neutron Probe: measure the return scatter of neutrons from a radio-isotope

Dielectric Measurements

- Coaxial Probe Impedance Measurement
 - Open-coaxial probe in contact with a dielectric material (including firm and ice)
 - Measure the complex S_{11} (reflection) parameter (amplitude and phase)
 - Approximation of of coaxial probe impedance using equivalence circuit:

$$i \omega (C_f + \epsilon_r(\omega)\chi) + G(\epsilon_r(\omega))^{2.5} = \frac{1}{Z_0} \frac{1 - S_{11}}{1 + S_{11}}$$

- Calibration using (3+) materials with known permittivity -> allows you to measure an unknown material



Radio propagation measurements

Density Measurements:

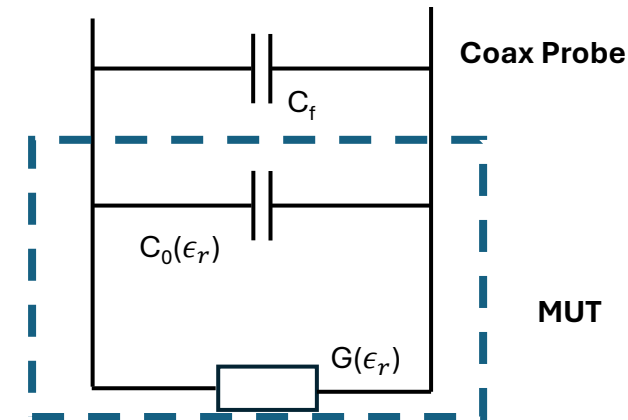
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In-Situ Ice Property Measurements

S_{11} Values

Example (right): Measurements of dielectric in the laboratory (KU)

S_{11} measured using miniVNA (PC controlled VNA)

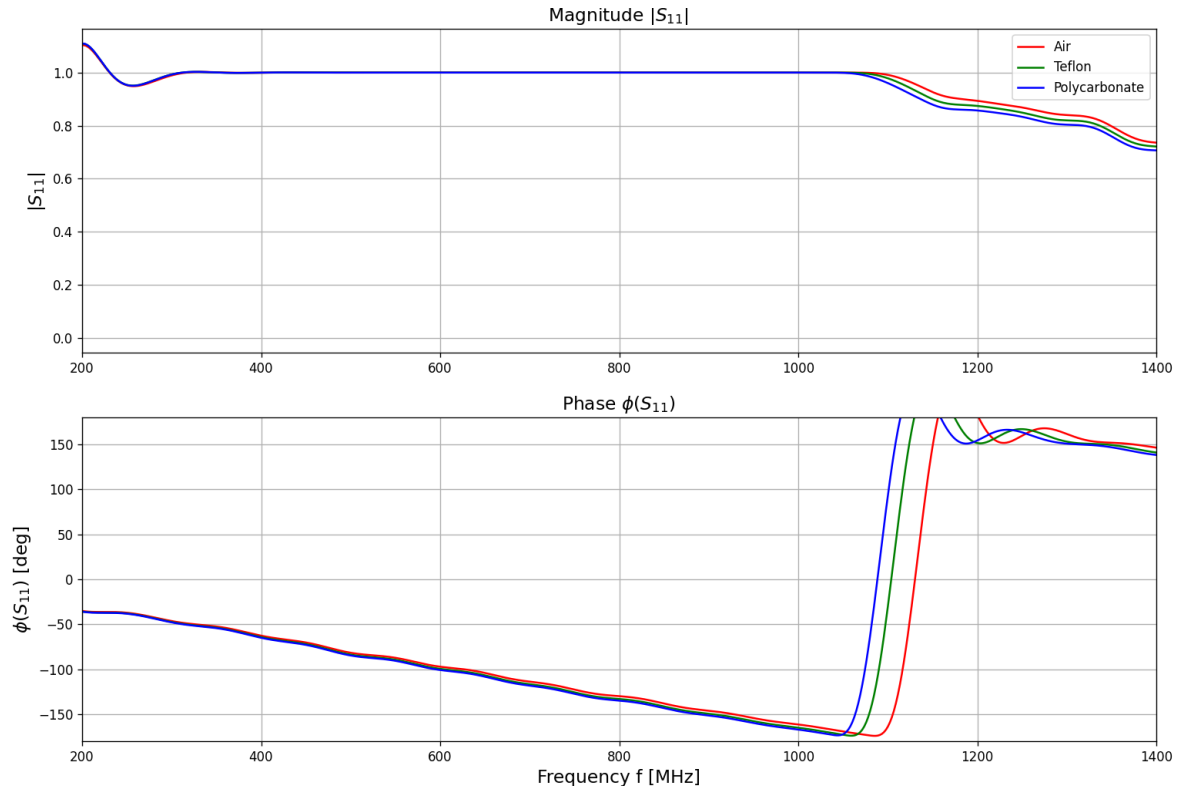
Materials used for calibration:

- Air
- Teflon (PTFE)
- Polycarbonate

Materials used for testing:

- HDPE
- Acrylic
- Polypropylene

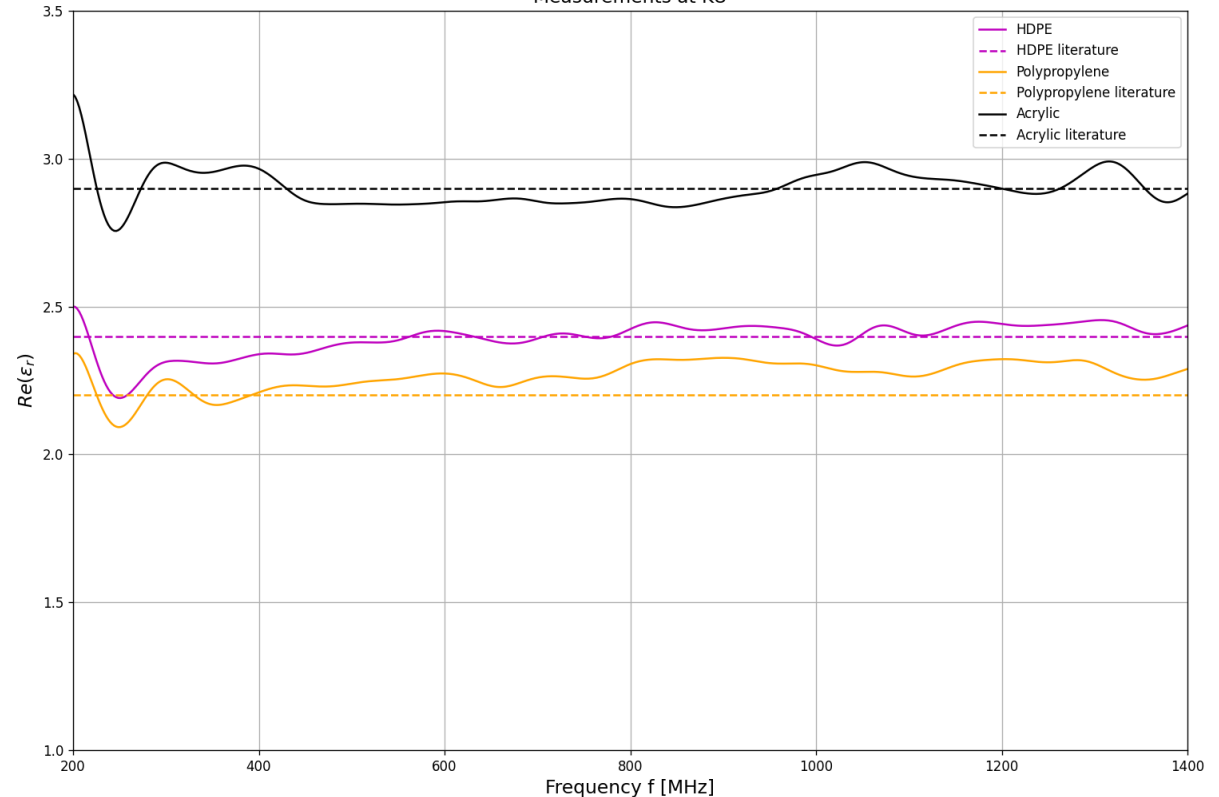
The real part of permittivity is most sensitive to the phase of S_{11}



In-Situ Ice Property Measurements

Reconstruction of ϵ_r

Measurements at KU



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S_{11} measured using miniVNA (PC controlled VNA)

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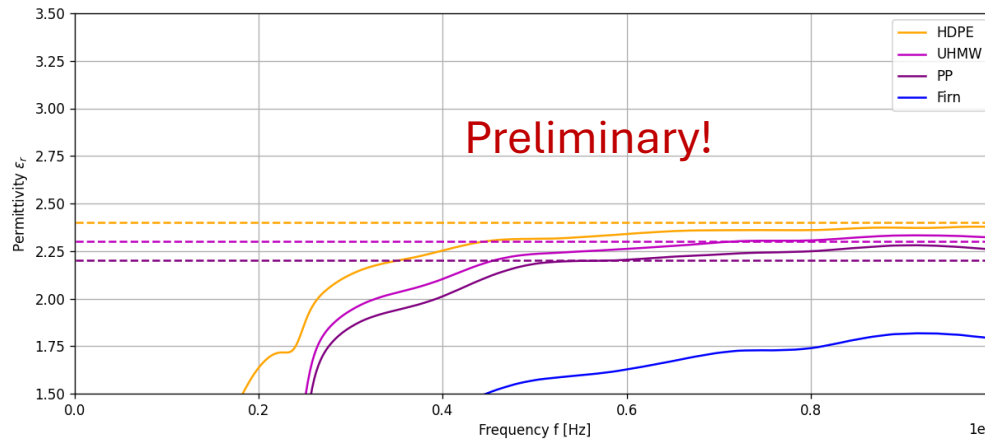
Permittivity reconstruction

Measurements in the field:

- Performed over two days → miniVNA needed to be recalibrated 3 times
- Performed with a 5 m RF cable between miniVNA and coaxial head → instrument calibration was incomplete
- High degree of frequency dependence in reconstructed permittivity → likely not physical
- Further work needed

Calibration 0, $z = 2.18$ m

ϵ_r Cal ID 0, $z = 2.18$ m



Calibration 2, $z = 9.45$ m

ϵ_r Cal ID 2, $z = 9.45$ m

