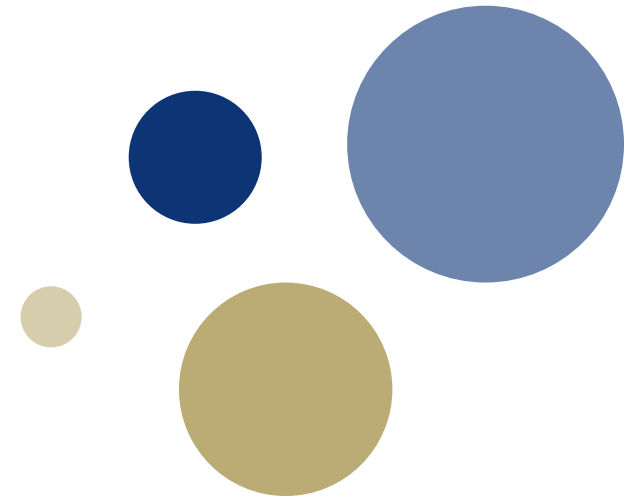




Norwegian University of  
Science and Technology



# **Do the LHAASO Galactic diffuse emission data require a contribution from unresolved sources?**

Presented by: Vittoria Vecchiotti

Based on work done in collaboration with: G. Peron, E. Amato, G. Morlino, S. Menchiari, F. L. Villante and G. Pagliaroli

# Outline:

1. Galactic gamma-ray diffuse emission model;
2. Source model, compatibility with LHAASO KM2A measurements, unresolved source contribution;
3. Comparisons with LHAASO diffuse emission measurements.

*Cataldo et al. Astrophys.J. 904 (2020)*

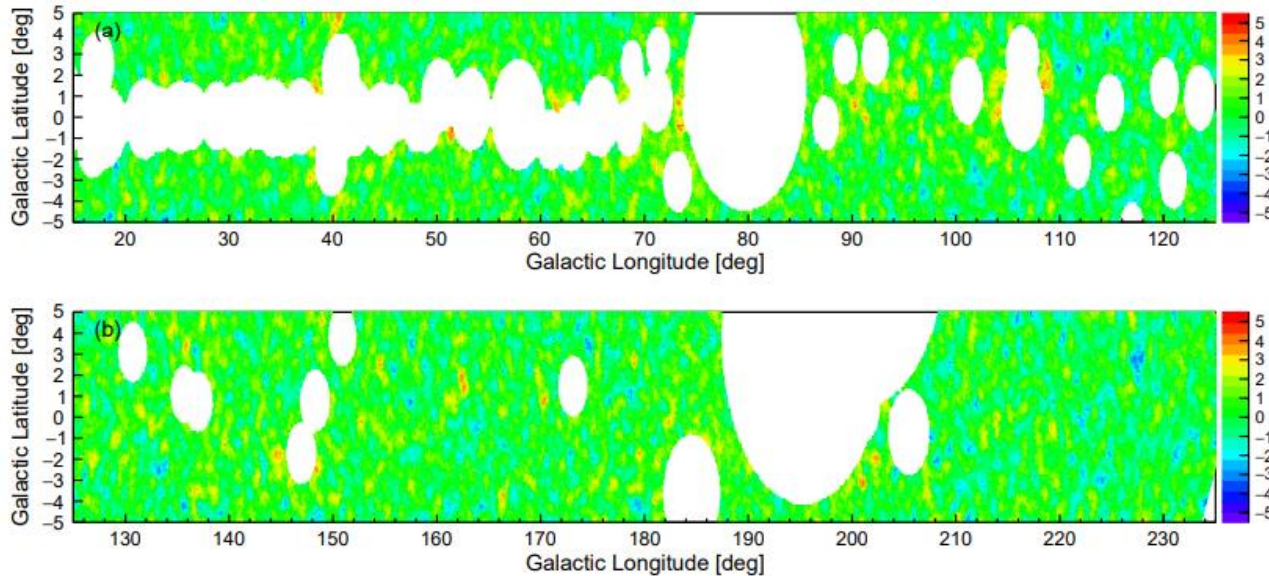
*Vecchiotti et al (2024), in preparation*

# LHAASO diffuse emission measurements:

$$\varphi_{\gamma,\text{diff}}^{\text{LHAASO}}$$



The LHAASO collaboration provided a measurement of the Galactic diffuse  $\gamma$ -ray emission in the energy range 10 TeV to 1 PeV in two sky regions by masking the contribution of known sources. *Z. Cao et al. 2023, Phys. Rev. Lett, 131*



**Do the LHAASO Galactic diffuse emission data are contaminated unresolved sources?**

$$\varphi_{\gamma,\text{diff}}^{\text{LHAASO}} = \varphi_{\gamma,S}^{\text{UnRes}} + \varphi_{\gamma,\text{diff}}$$

Population study H.E.S.S.

Models:  
Assumptions on the CR spatial and energy distributions, cross-section, and ISM.

# Diffuse gamma-ray emission:

1. Differential inelastic cross section of pp interaction.

2. Interstellar gas distribution in the Galaxy

$$\varphi_{\gamma}(E_{\gamma}, \hat{n}_{\gamma}) = \int_{E_{\gamma}}^{\infty} dE \frac{d\sigma(E, E_{\gamma})}{dE_{\gamma}} \int_0^{\infty} dl \varphi_{CR}(E, \bar{r}_{\odot} + l\hat{n}_{\gamma}) n_H(\bar{r}_{\odot} + l\hat{n}_{\gamma})$$

3. Cosmic-ray energy and spatial distribution

# 1. Cross section:

$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

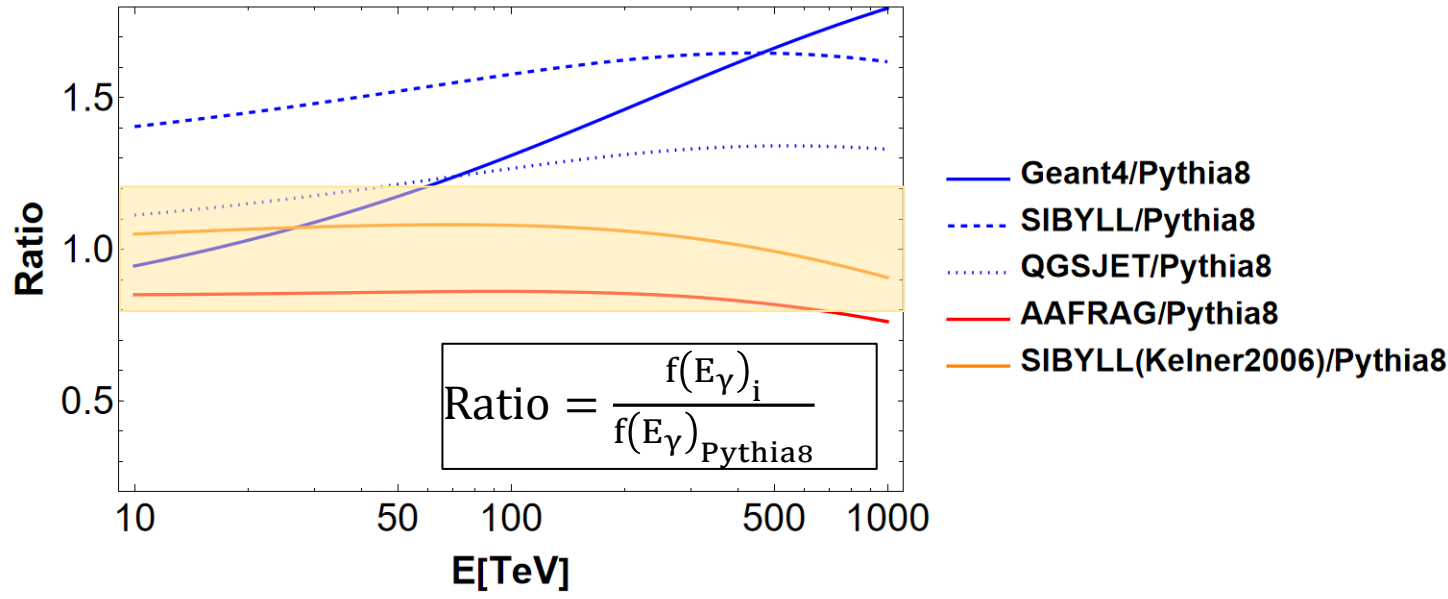
We compare:

- different parameterization based on different MC codes SIBYLL, QGSJET, Pythia8, and Geant4 (*Kafexhiu et al 2014*);
- SIBYLL (*Kelner et al 2006*);
- AAFRAG based on QGSJET-II-04m (*M. Kachelriess et al 2022*)

# 1. Cross section:

$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

$$\varphi_{CR}(E) = E^{-2.7}$$



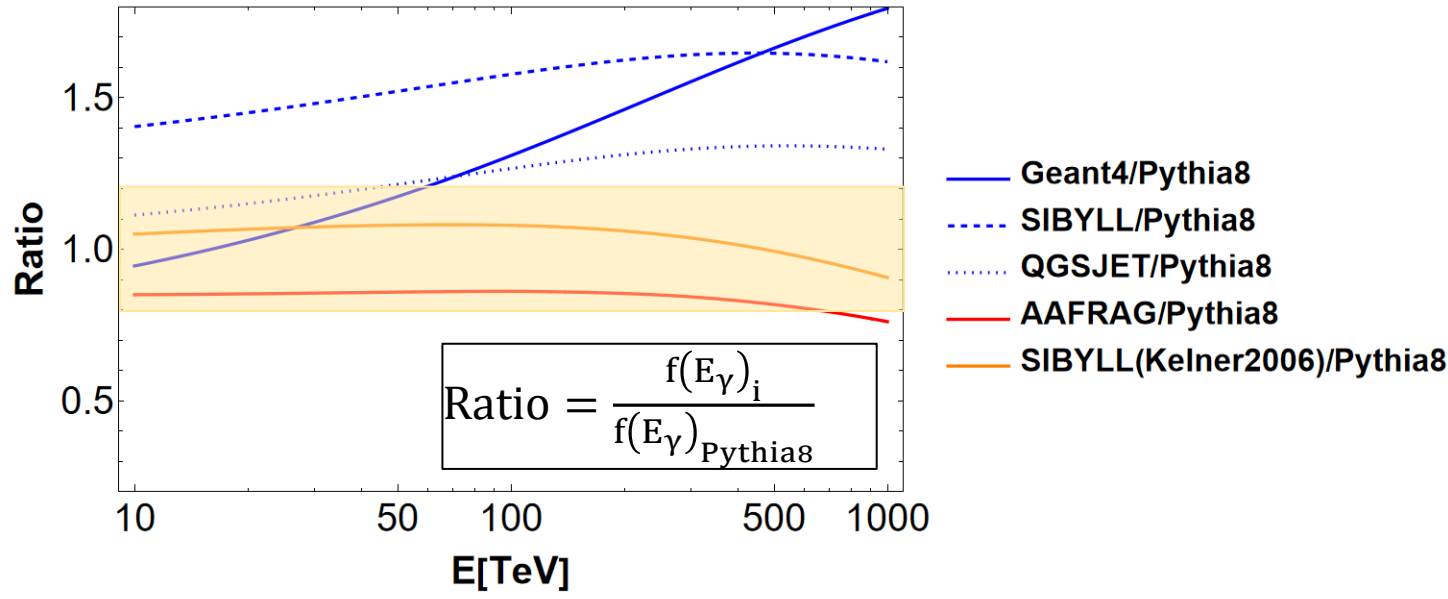
We compare:

- different parameterization based on different MC codes SIBYLL, QGSJET, Pythia8, and Geant4 (*Kafexhiu et al 2014*);
- SIBYLL (*Kelner et al 2006*);
- AAFRAG based on QGSJET-II-04m (*M. Kachelriess et al 2022*)

# 1. Cross section:

$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

$$\varphi_{CR}(E) = E^{-2.7}$$



- We compare:
- different parameterization based on different MC codes SIBYLL, QGSJET, Pythia8, and Geant4 (*Kafexhiu et al 2014*);
  - SIBYLL (*Kelner et al 2006*);
  - AAFRAG based on QGSJET-II-04m (*M. Kachelriess et al 2022*)

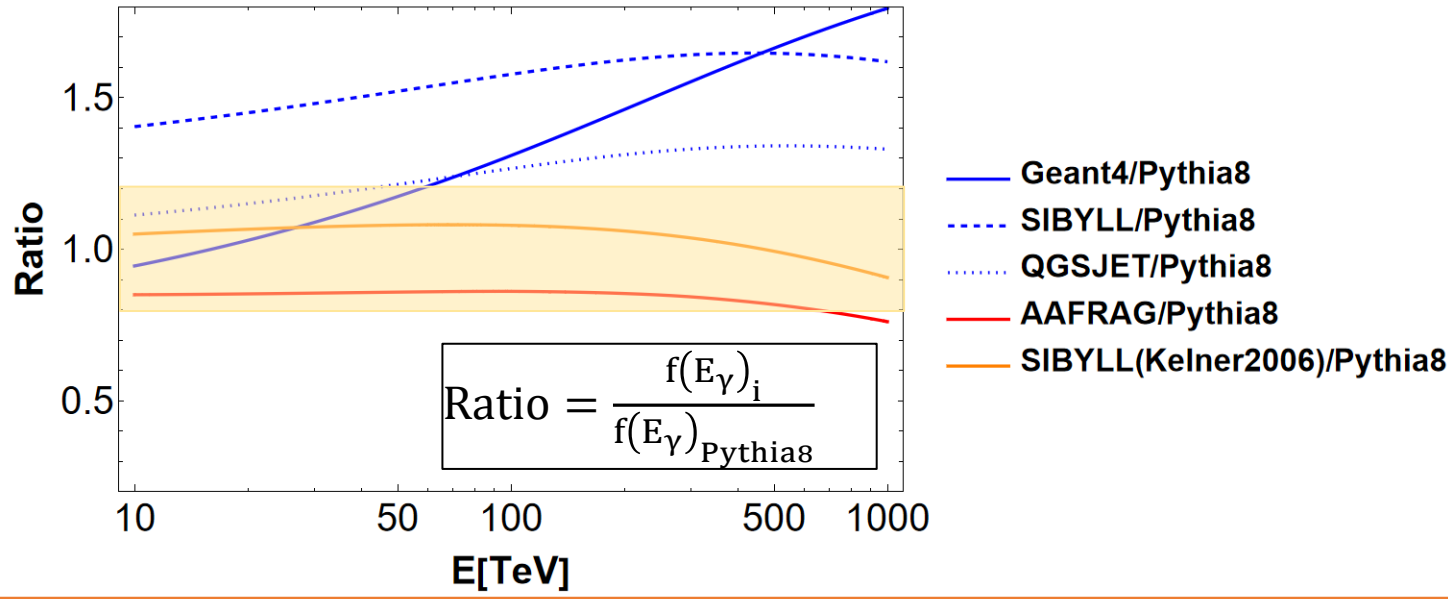
Assumptions cross section:

We take AAFRAG for the fiducial case and SIBYLL (*Kafexhiu et al 2014*) to include uncertainties.

# 1. Cross section:

$$f(E_\gamma) = \int_{E_\gamma}^{\infty} dE \frac{d\sigma(E, E_\gamma)}{dE_\gamma} \varphi_{CR}(E)$$

$$\varphi_{CR}(E) = E^{-2.7}$$



We compare:

- different parameterization based on different MC codes SIBYLL, QGSJET, Pythia8, and Geant4 (*Kafexhiu et al 2014*);
- SIBYLL (*Kelner et al 2006*);
- AAFRAG based on QGSJET-II-04m (*M. Kachelriess et al 2022*)

Assumptions cross section:

We take AAFRAG for the fiducial case and SIBYLL (*Kafexhiu et al 2014*) to include uncertainties.

# 2. ISM:

$$n_c = \frac{\int_{\Delta\Omega} d\Omega \int_0^\infty dl n_H (\bar{r}_{Sun} + l \hat{n}_\gamma)}{\Delta\Omega}$$

- Dust (masked):  $n_{c,inner} = 7.85068 \times 10^{21} \text{ cm}^{-2}$   
PLANCK  
*(Aghanim et al. 2016)*
- Galprop (masked):  $n_{c,inner} = 9.47769 \times 10^{21} \text{ cm}^{-2}$   
 $n_{c,outer} = 8.72069 \times 10^{21} \text{ cm}^{-2}$

Galprop provides ~ 20% more target than the dust in the inner region and ~ 50 % in the outer region

Assumptions ISM:

We take Galprop for the fiducial case and the Dust to include uncertainties.



### 3. Cosmic ray distribution:

$$\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$$

### 3. Cosmic ray distribution:

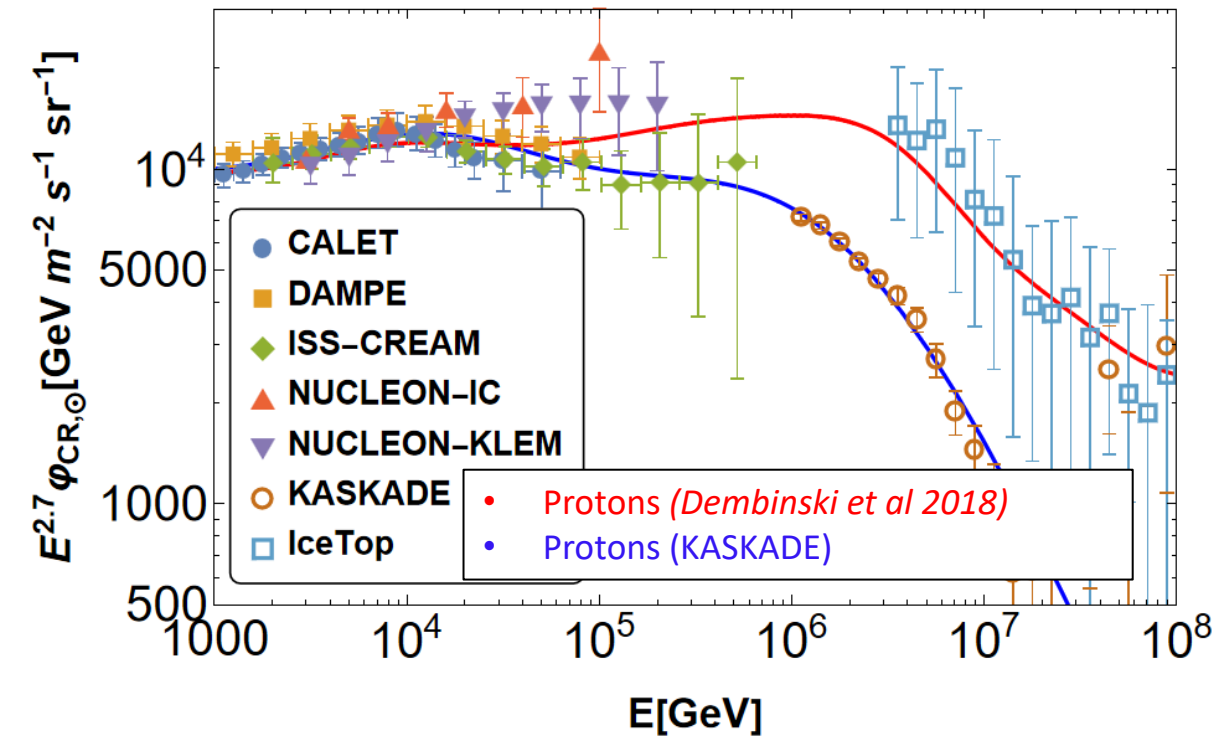
$$\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$$



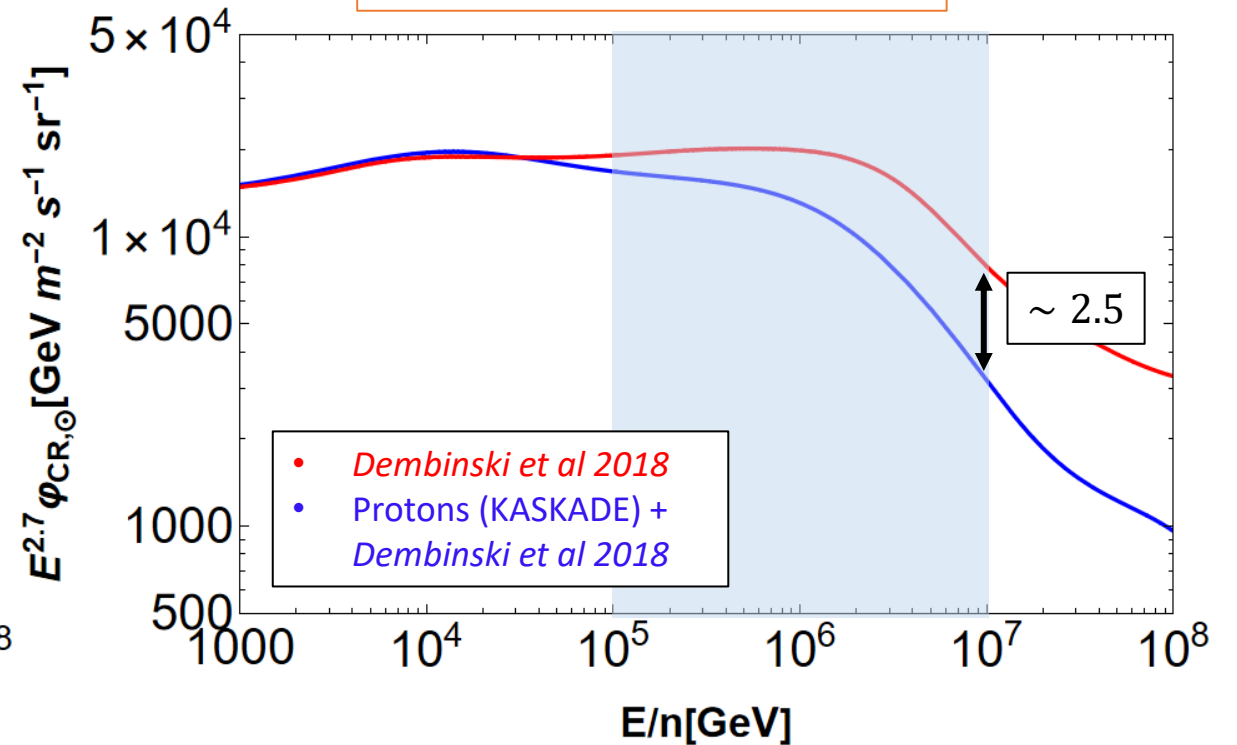
Assumption CR spectrum:

We take the data-driven CR spectrum from *Dembinski et al 2018* for the fiducial case and **Protons (KASKADE) + *Dembinski et al 2018*** to include uncertainties.

Protons



All Particle Spectrum



### 3. Cosmic ray distribution: $\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$

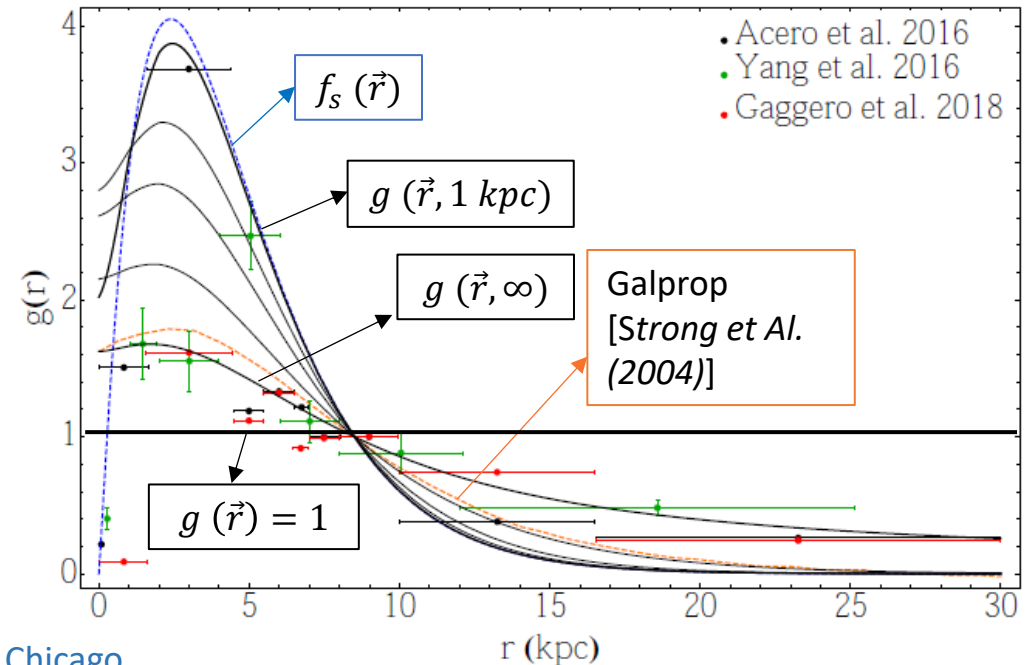


Assumption CR spectrum:

We take the data-driven CR spectrum from **Dembinski et al 2018** for the fiducial case and **Protons (KASKADE) + Dembinski et al 2018** to include uncertainties.



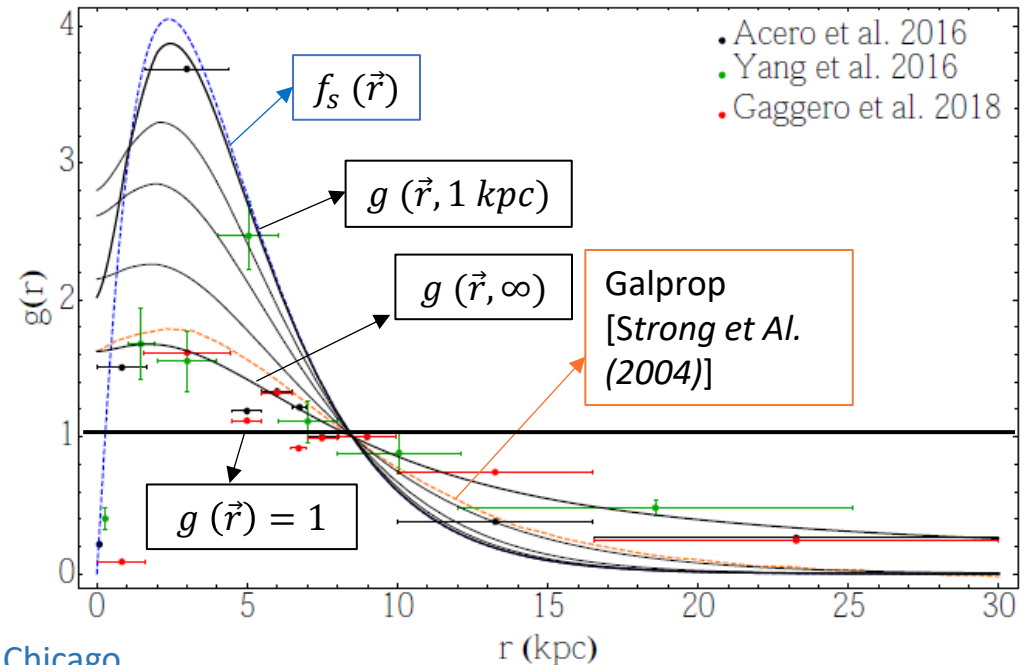
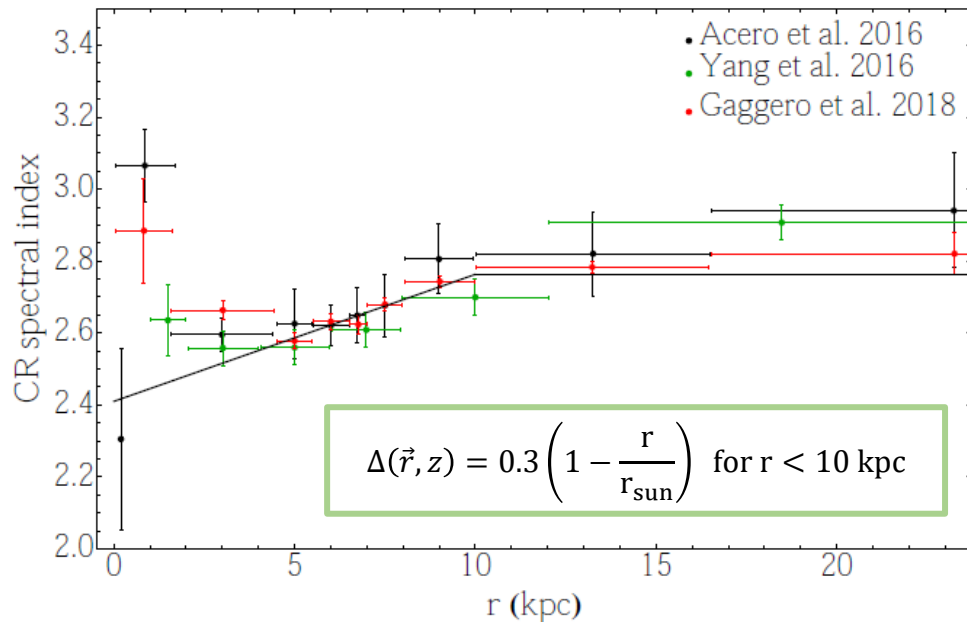
$g(r)$  is determined by the distribution of the CR sources  $f_s(\vec{r})$  (proportional to the SNR number density by Green et al. (2015), and by the propagation of CR in the Galactic magnetic field.



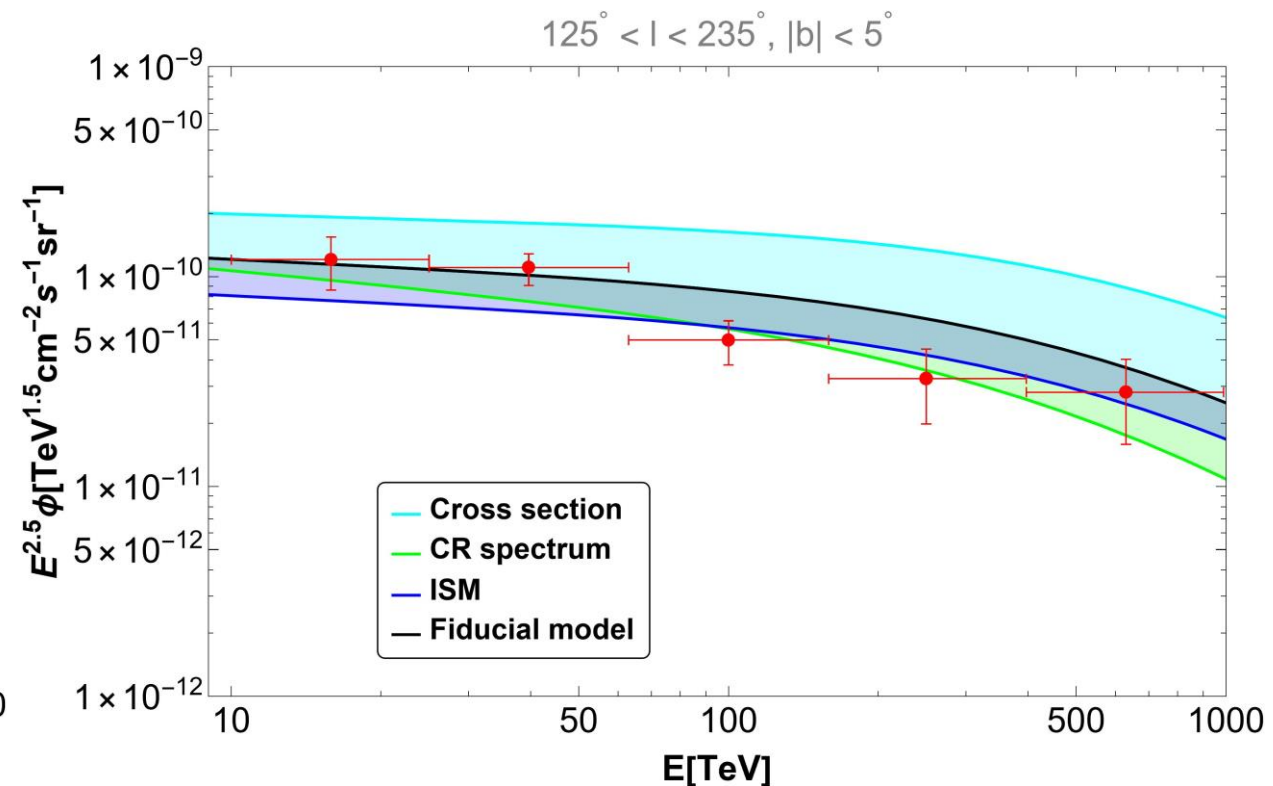
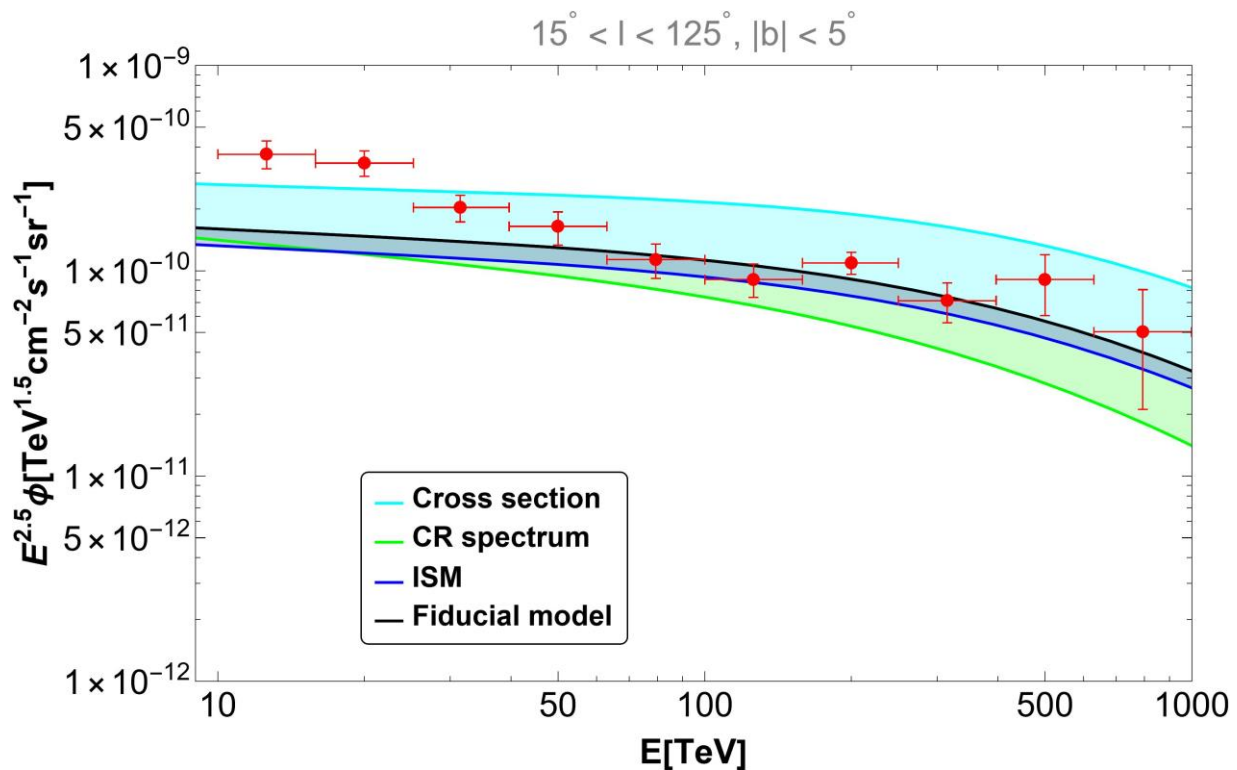
### 3. Cosmic ray distribution: $\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$

- ★ Assumption CR spectrum: We take the data-driven CR spectrum from ***Dembinski et al 2018*** for the fiducial case and **Protons (KASCADE) + *Dembinski et al 2018*** to include uncertainties.
- ★  $g(r)$  is determined by the distribution of the CR sources  $f_s(\vec{r})$  (proportional to the SNR number density by Green et al. (2015), and by the propagation of CR in the Galactic magnetic field.
- ★ **2 cases: with and without spatially dependent CR spectral index** (from the analysis of the FermiLAT data at  $\sim 20$  GeV  
*Acero et al. (2016), Yang et al. (2016), Gaggero et al. (2018)*)

$$h(E, \vec{r}) = \left( \frac{E}{20 \text{ GeV}} \right)^{\Delta(\vec{r})}$$



# Comparison with LHAASO (standard diffusion):



- LHAASO data can be explained by the “truly” diffuse emission in the outer region and in the inner region above  $\sim 30$  TeV;  
The contribution from unresolved sources must be negligible.

# Pulsar Wind Nebulae population:

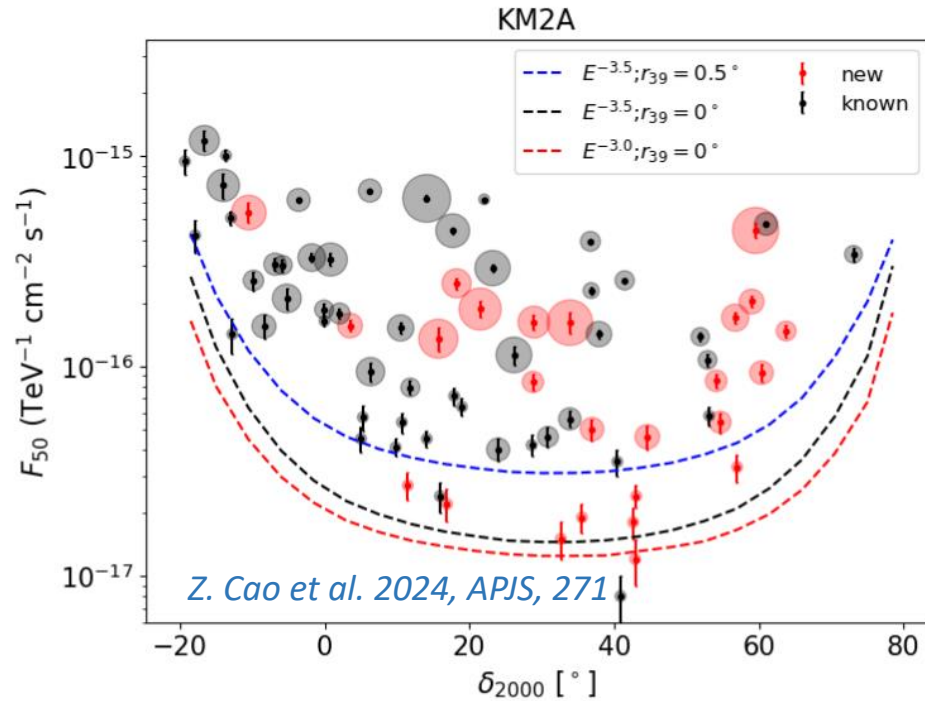
Cataldo et al. *Astrophys.J.* 904 (2020); Pagliaroli et al, *Universe*, 9,881 (2023)

We built a synthetic population of PWNe using the best fit of the maximum luminosity in the energy range 1 – 100 TeV:  $L_{max} = 2.2 \times 10^{35}$  erg/s and the spin-down timescale:  $\tau_{sd} = 2.9$  kyrs derived from fitting the brightest sources of the HGPS.

## Assumptions:

- Latitude, longitude, and radius are extracted from the Lorimer distribution that scales as  $\exp(-|z|/H)$  with  $H = 0.05$  kpc (it is the value that provides the best chi-square in the fit of HGPS data);
- The age of sources  $t_{age}$  is extracted uniformly in the interval  $[1, 10^6]$  yr;
- The luminosity is calculated from:  $L = L_{max} \left(1 + \frac{t_{age}}{\tau_{sd}}\right)^{-2}$
- Spectrum: power-law with exponential cut-off ( $E_{cut} = 100$  TeV), spectral index fixed to: 2.4.

# Source contributions:



“The flux sensitivity is defined as the flux normalization required to have 50% probability of detecting a source at  $5\sigma$  level”

At 50 TeV the differential threshold of point-like sources depends only mildly on the spectral assumption.

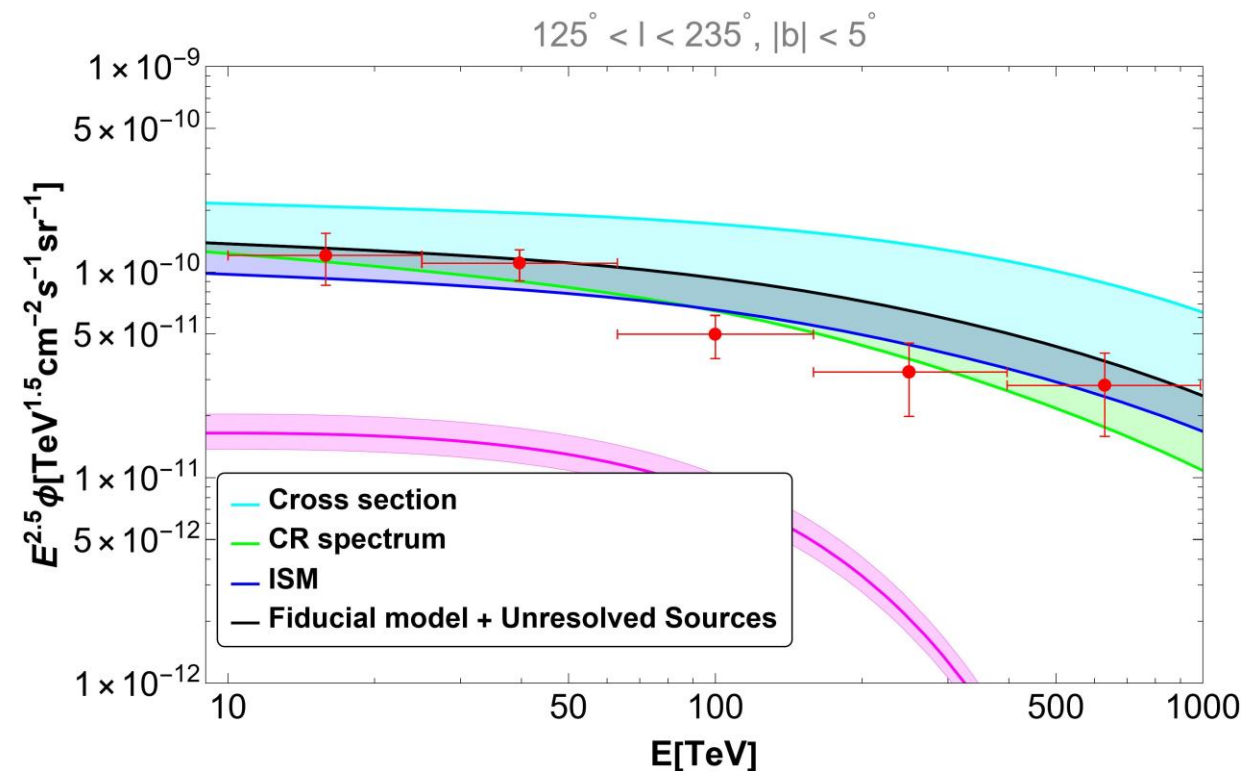
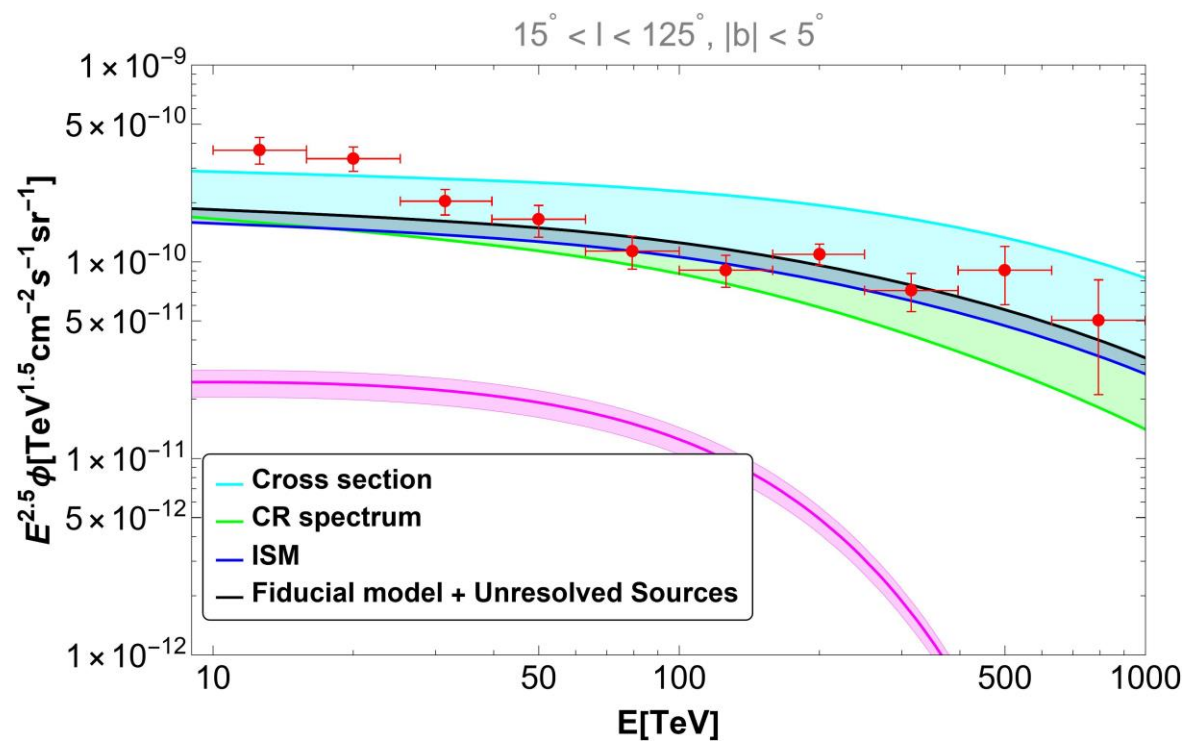
Def: Resolved sources:  $\phi_{50} > \phi_{th,50}$ ,  $\phi [TeV^{-1} cm^{-2} s^{-1}]$

		$N_R$	$\varphi_R$	$\varphi_{UNR}$	$\varphi_{UNR,H}$
$15^\circ < l < 235^\circ$	MC	$84_{-5}^{+7}$	$1.69_{-0.43}^{+0.62} \times 10^{-14}$	$2.82_{-0.14}^{+0.15} \times 10^{-15}$	—
	KM2A	76	$1.72 \times 10^{-14}$	—	—
$15^\circ < l < 125^\circ$	MC	$73_{-8}^{+5}$	$1.32_{-0.33}^{+0.37} \times 10^{-14}$	$2.56_{-0.16}^{+0.14} \times 10^{-15}$	$2.23_{-0.36}^{+0.34} \times 10^{-16}$
	KM2A	59	$1.44 \times 10^{-14}$	—	—
$125^\circ < l < 235^\circ$	MC	$12_{-2}^{+3}$	$2.82_{-1.1}^{+1.8} \times 10^{-15}$	$2.53_{-0.35}^{+0.46} \times 10^{-16}$	$2.08_{-0.34}^{+0.49} \times 10^{-16}$
	KM2A	16	$2.74 \times 10^{-15}$	—	—

Results:

- The predicted number and flux of resolved sources are compatible with the KM2A quantities within  $2\sigma$ ;
- The unresolved source flux is suppressed by 91 % and 18 % in the inner and outer regions, respectively.

# Comparison with LHAASO (standard diffusion):



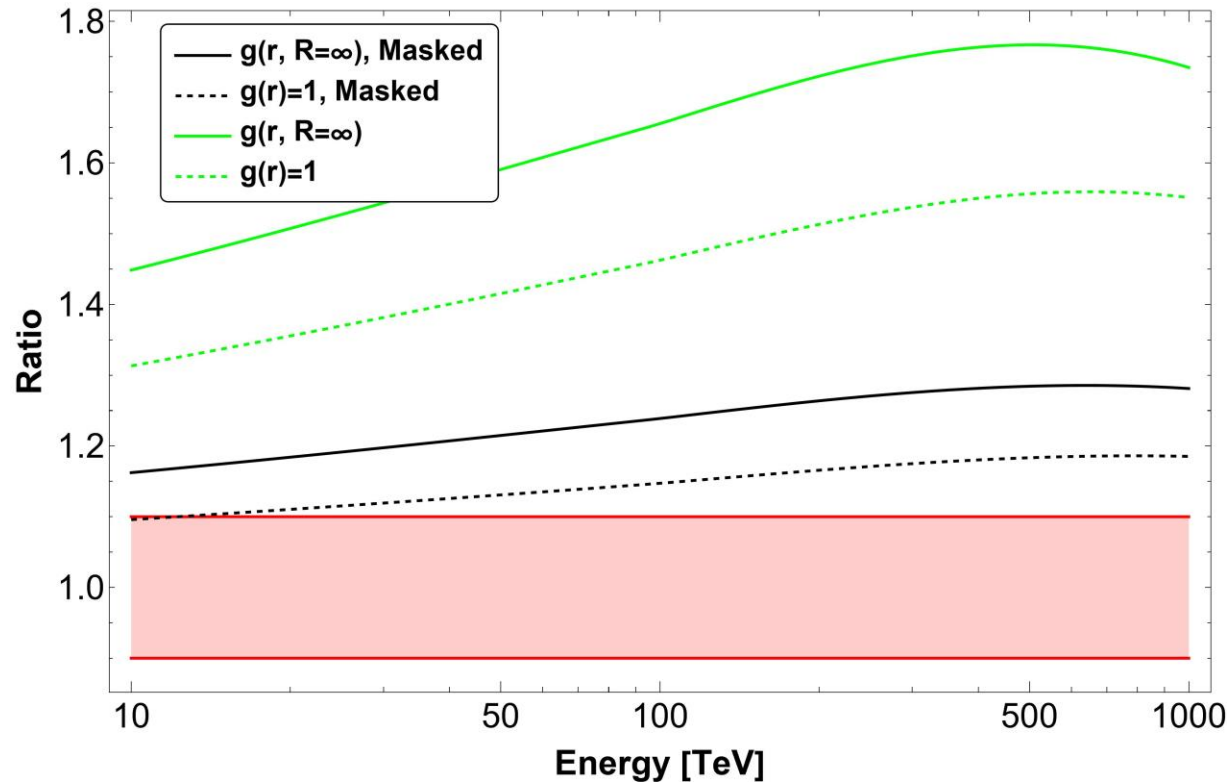
- The masks cancel out most of the unresolved source contributions. Unresolved sources contribute  $\sim 15\%$  of the fiducial model at 50 TeV in both regions;



# Effect of LHAASO masks on the hardening:

Def: hardening= spatially dependent CR spectral index

Inner Region:  $15^\circ < l < 125^\circ$



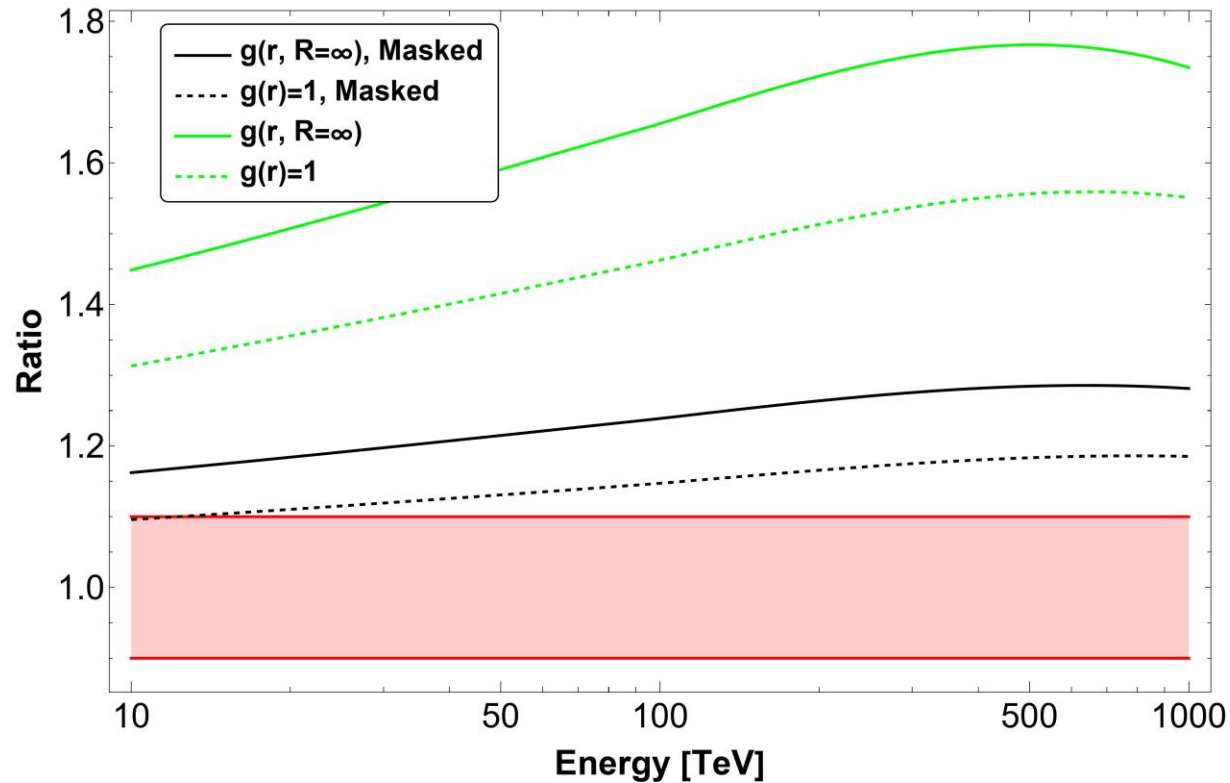
$$\text{Ratio} = \frac{\Phi_{\text{hardening}}}{\Phi_{\text{standard}}}$$

The ratio is *independent* of the cross-section, ISM and CR spectrum but it *depends* on the CR spatial distribution

# Effect of LHAASO masks on the hardening:

Def: hardening= spatially dependent CR spectral index

Inner Region:  $15^\circ < l < 125^\circ$



$$\text{Ratio} = \frac{\Phi_{\text{hardening}}}{\Phi_{\text{standard}}}$$

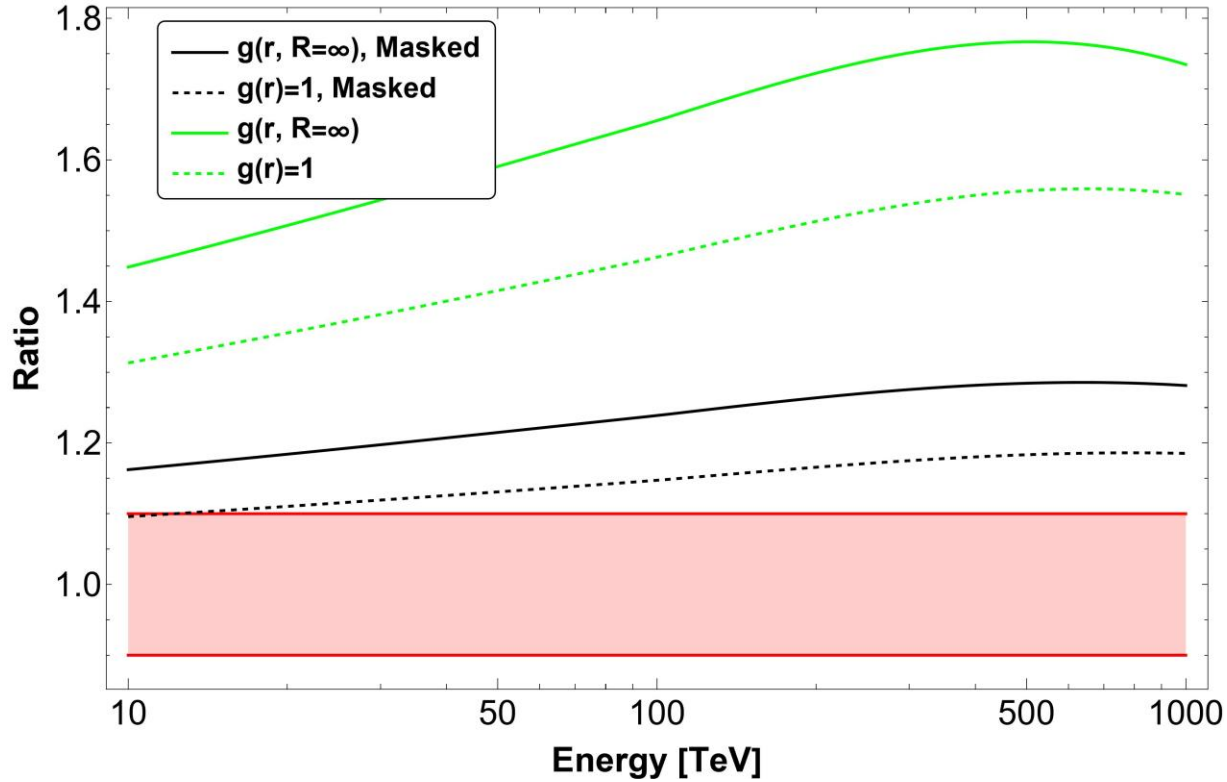
The ratio is *independent* of the cross-section, ISM and CR spectrum but it *depends* on the CR spatial distribution

- $g(r, R = \infty)$ : The  $\Phi_{\text{hardening}}$  produces 76 % more signal than  $\Phi_{\text{standard}}$  at 500 TeV ;
- $g(r) = 1$ : The  $\Phi_{\text{hardening}}$  produces 55% more signal than  $\Phi_{\text{standard}}$  at 500 TeV

# Effect of LHAASO masks on the hardening:

Def: hardening= spatially dependent CR spectral index

Inner Region:  $15^\circ < l < 125^\circ$



$$\text{Ratio} = \frac{\Phi_{\text{hardening}}}{\Phi_{\text{standard}}}$$

The ratio is *independent* of the cross-section, ISM and CR spectrum but it *depends* on the CR spatial distribution

- $g(r, R = \infty)$ : The  $\Phi_{\text{hardening}}$  produces 76 % more signal than  $\Phi_{\text{standard}}$  at 500 TeV ;
- $g(r) = 1$ : The  $\Phi_{\text{hardening}}$  produces 55% more signal than  $\Phi_{\text{standard}}$  at 500 TeV

After masking:

- $g(r, R = \infty)$ :  $\Phi_{\text{hardening}}$  produces 28 % more signal than  $\Phi_{\text{standard}}$  at 500 TeV
- $g(r) = 1$ :  $\Phi_{\text{hardening}}$  produces 18% more signal  $\Phi_{\text{standard}}$  at 500 TeV

# Conclusions:

1. The total flux and the number of sources derived in *Cataldo et al 2019* based on the HGPS are compatible with the observation of KM2A within  $2\sigma$ ;
2. The LHAASO masks cancel most of the effect due to unresolved sources in the inner region (suppressed by 91 %). In the outer region, unresolved sources already produce a negligible contribution to the diffuse emission that is further suppressed by the LHAASO mask of about 18 %;
3. The cross-section, CR spectrum, and ISM uncertainties are non-negligible. However, the LHAASO data seems compatible with the “truly” diffuse emission within uncertainties except for the 2 low energy points in the inner region which could be explained by introducing other classes of unresolved sources.
4. The LHAASO masks significantly reduce the effect of a spatial-dependent CR spectral index. As a consequence, it is challenging to test this hypothesis using LHAASO data.

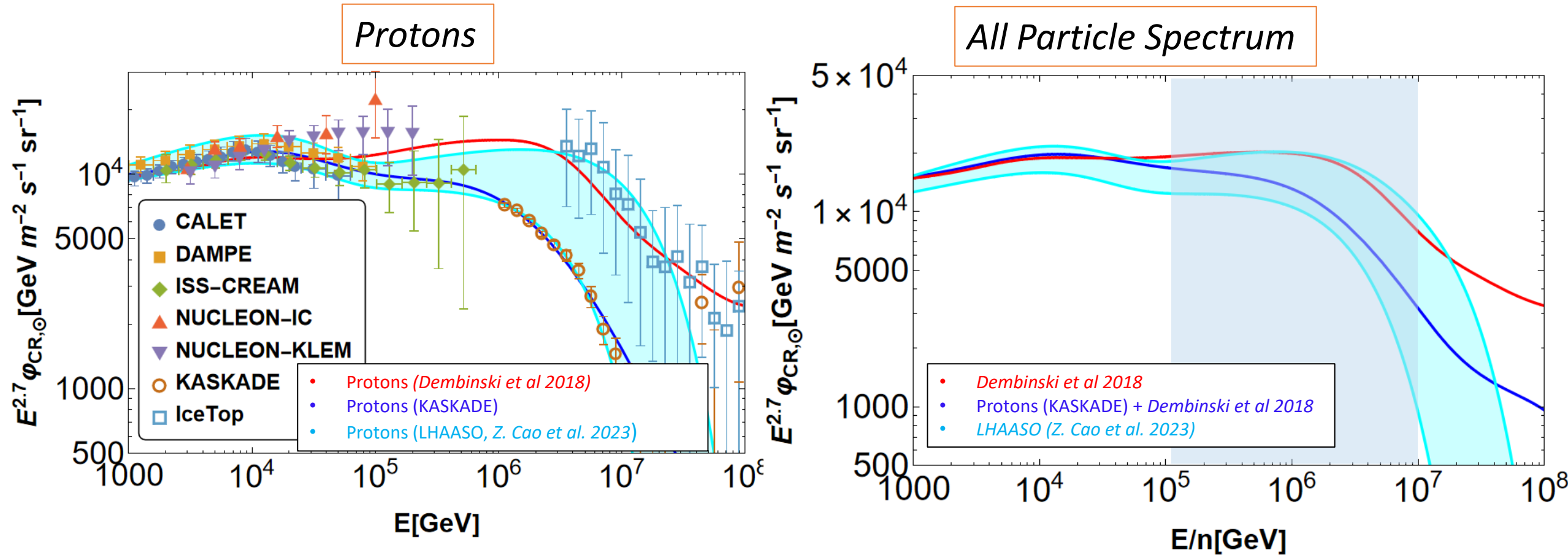
Backup slides

### 3. Cosmic ray distribution: $\varphi_{CR}(E, \vec{r}) = \varphi_{CR, Sun}(E) g(\vec{r}, R) h(E, \vec{r})$

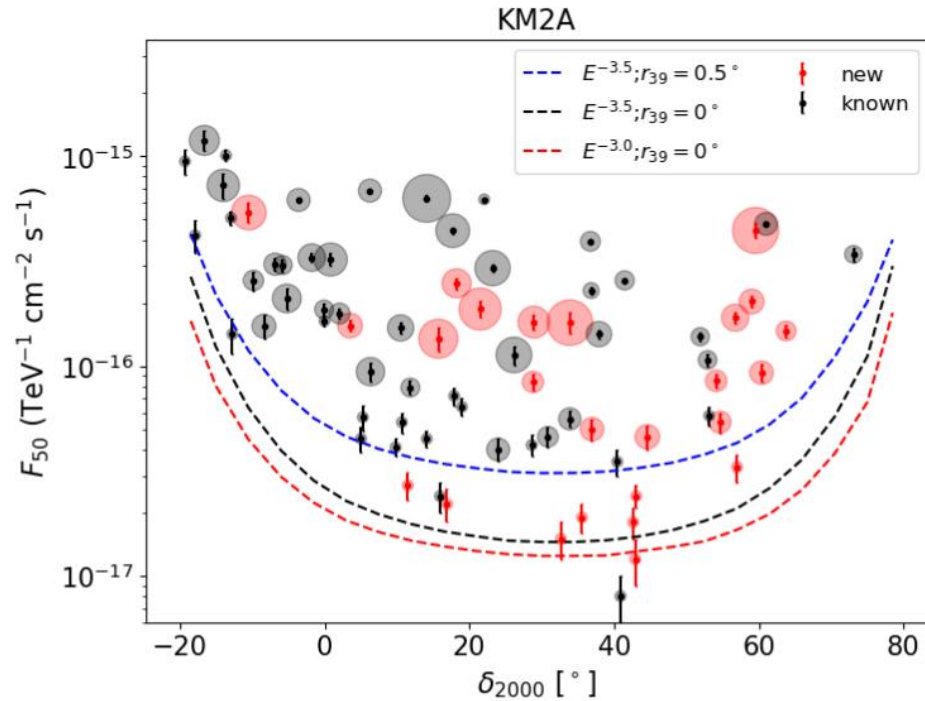


Assumption CR spectrum:

We take the data-driven CR spectrum from **Dembinski et al 2018** for the fiducial case and **Protons (KASKADE) + Dembinski et al 2018** to include uncertainties.



# Source contributions (Size 40 pc):



Best fit for size 40 pc:  
 $L_{max} = 2 \times 10^{35}$  erg/s  
 $\tau_{sd} = 4.6$  kyrs

Def: Resolved sources:  $\phi_{50} > \phi_{th,50} \sqrt{(\sigma_{psf}^2 + \sigma_s^2) / \sigma_{psf}^2}$  where  $\sigma_{psf} = 0.2^\circ$  and  $\sigma_s$  is the angular size of the source.

		$N_R$	$\varphi_R$	$\varphi_{UNR}$	$\varphi_{UNR,H}$
$15^\circ < l < 235^\circ$	MC	$69_{-5}^{+5}$	$2.28_{-0.56}^{+0.74} \times 10^{-14}$	$7.17_{-0.45}^{+0.37} \times 10^{-15}$	—
	KM2A	76	$1.72 \times 10^{-14}$	—	—
$15^\circ < l < 125^\circ$	MC	$62_{-6}^{+4}$	$1.76_{-0.45}^{+0.47} \times 10^{-14}$	$6.10_{-0.41}^{+0.37} \times 10^{-15}$	$8.19_{-1.64}^{+1.74} \times 10^{-16}$
	KM2A	59	$1.44 \times 10^{-14}$	—	—
$125^\circ < l < 235^\circ$	MC	$7_{-1}^{+2}$	$4.21_{-1.8}^{+3.4} \times 10^{-15}$	$1.05_{-0.27}^{+0.25} \times 10^{-15}$	$8.32_{-1.99}^{+2.49} \times 10^{-16}$
	KM2A	16	$2.75 \times 10^{-15}$	—	—

Results:

- The predicted number and flux of resolved sources are compatible with the KM2A quantities within  $2\sigma$  (except in the outer region);
- The unresolved source flux is suppressed by 86 % and 20 % in the inner and outer regions, respectively.

# Summary assumptions:

Assumptions for the diffuse emission fiducial model:

- CRs: *Dembinki et al 2018*;
- Gas: Galprop;
- Cross section: AAFRAG;
- CR spatial distribution of CR:  $g(\vec{r}, \infty)$

Variation with respect to the fiducial model:

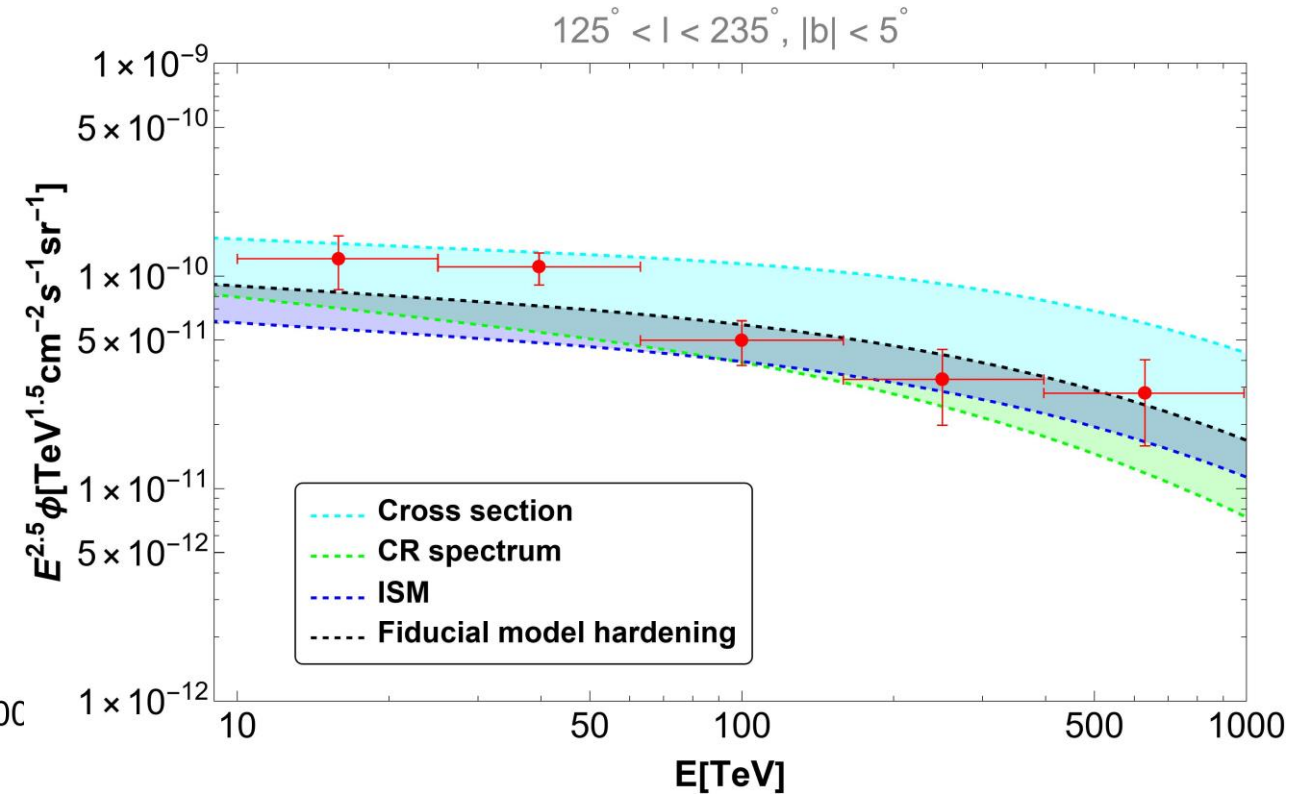
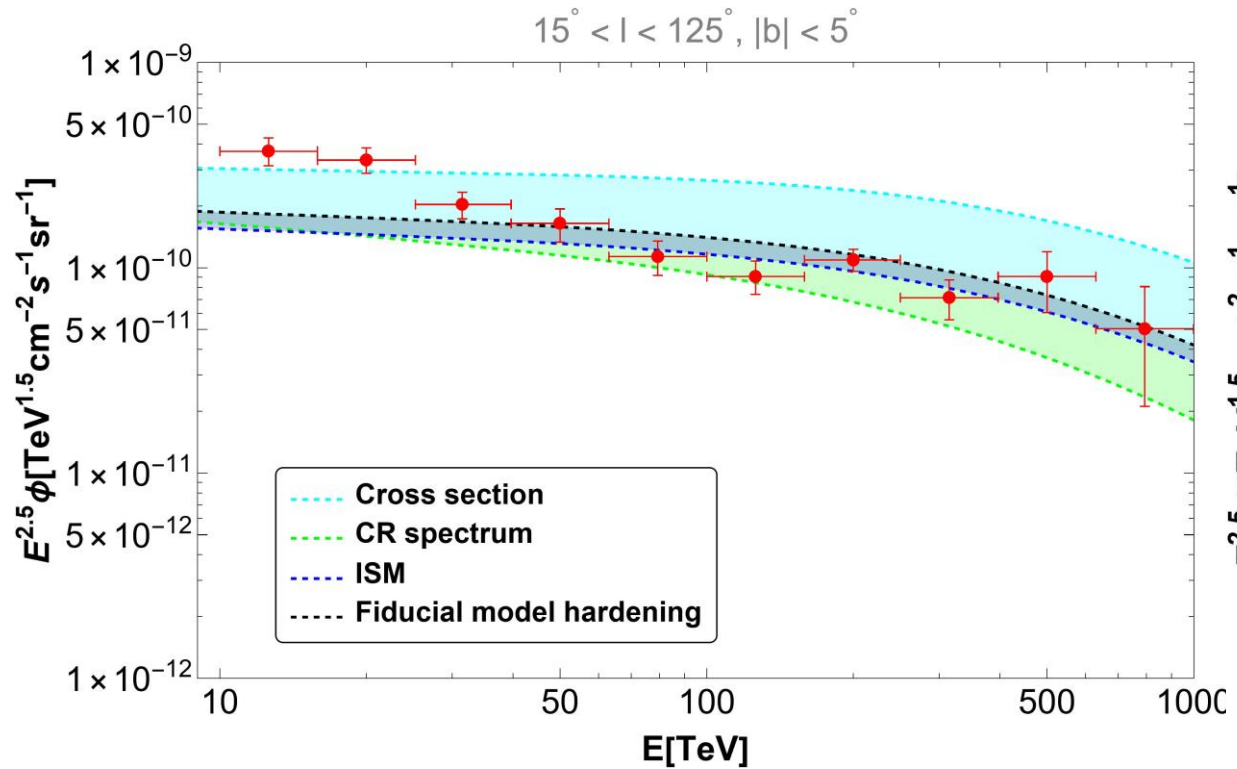
- CRs: fit protons (KASKADE) + All elements (*Dembinki et al 2018*);
- Gas: Dust;
- Cross section: Sybill;
- CR spatial distribution of CR:  $g(\vec{r}, \infty)$

Assumptions unresolved sources:

- Spectrum: power-law with exponential cut-off ( $E_{\text{cut}} = 100$  TeV), spectral index fixed to: 2.4.
- Thickness of the disk  $H=0.05$  kpc.

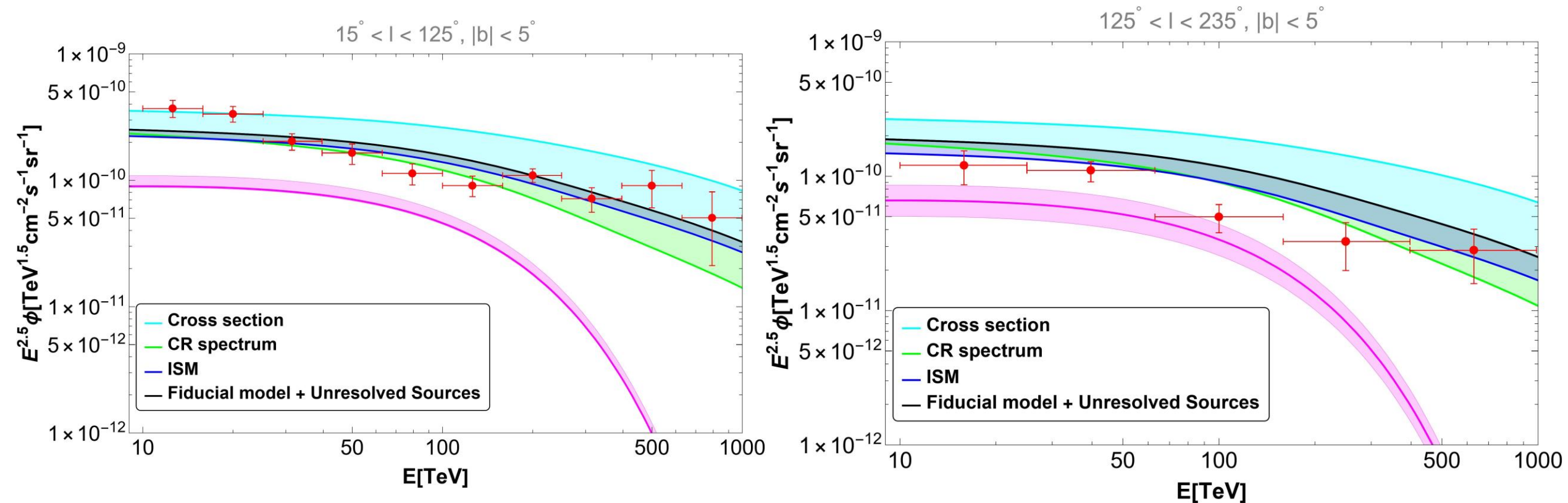


# Comparison with LHAASO (hardening effect):



- LHAASO cannot be used to distinguish the hardening hypothesis from the standard one;

# Comparison with LHAASO (standard diffusion and size 40 pc):



- Unresolved sources contribute  $\sim 54\%$  of the fiducial model at 50 TeV in both regions;
- The 40 pc size case corresponds to an upper limit for the unresolved source contribution.