



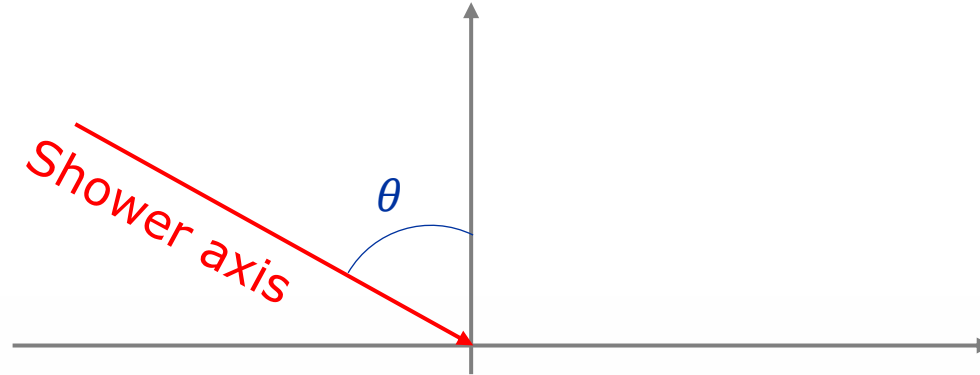
# GEOMETRIC BOOSTING IN VERY INCLINED AIR SHOWERS ARENA MEETING 11/06/2024

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## WHAT ARE VERY INCLINED AIR SHOWERS?

Airshowers that propagate nearly horizontally (zenith angle  $\theta > 80^\circ$ )



Of particular interest for next generation radio based detectors

Emission from these airshowers could require more careful treatment  
=> Do simulations have to take this into account?

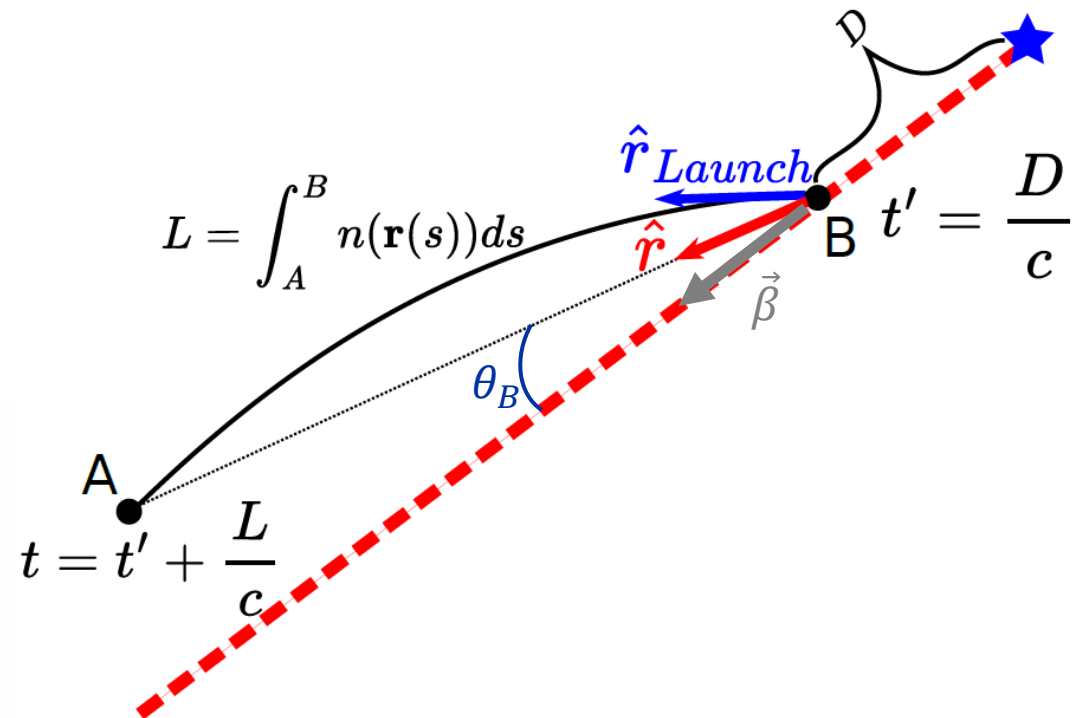
# HOW TO DESCRIBE THE EFFECTS OF A NON-UNIFORM ATMOSPHERE?

Developed a raytracer based on Fermat's principle, combined with a line model of a cascade

Geometric boosting described by the boostfactor  
For **uniform** media:

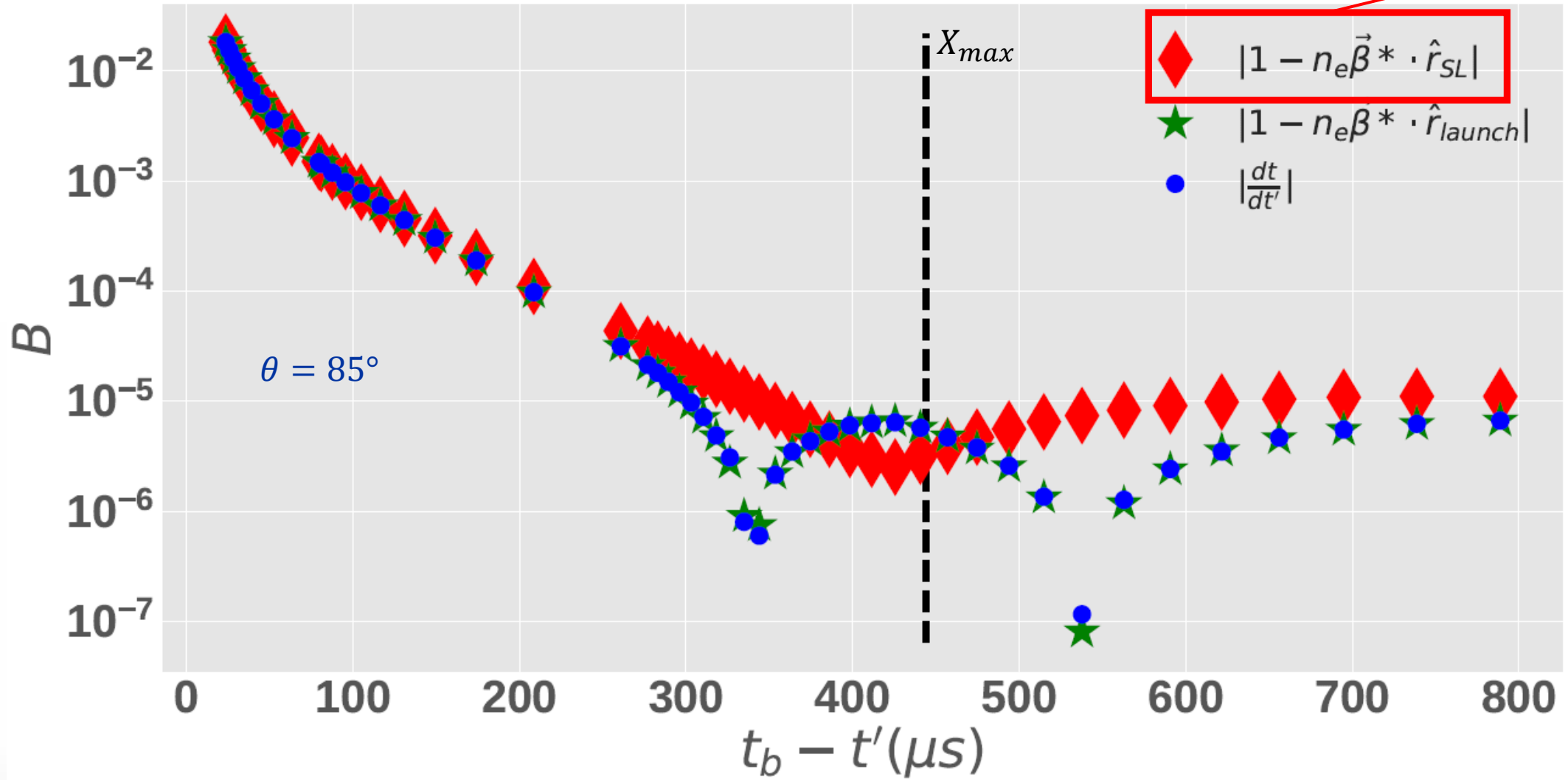
$$\frac{dt}{dt'} = 1 - n \cdot \beta \cdot \cos(\theta_B) = \text{Boostfactor}^{-1}$$

We wish to generalise this to non uniform media



# HOW TO DESCRIBE A NON-UNIFORM ATMOSPHERE?

Currently used in CoREAS



## WHERE DOES THE BOOSTFACTOR MATTER?

Appears in, for example, the end point formalism which is used in CoREAS

$$\vec{E}_{\pm}(\vec{x}, t) = \pm \frac{1}{\Delta t} \frac{q}{c} \left( \frac{\hat{r} \times [\hat{r} \times \vec{\beta}^*]}{\underbrace{(1 - n\vec{\beta}^* \cdot \hat{r})}_{\text{Boostfactor}^{-1}}} R \right)$$

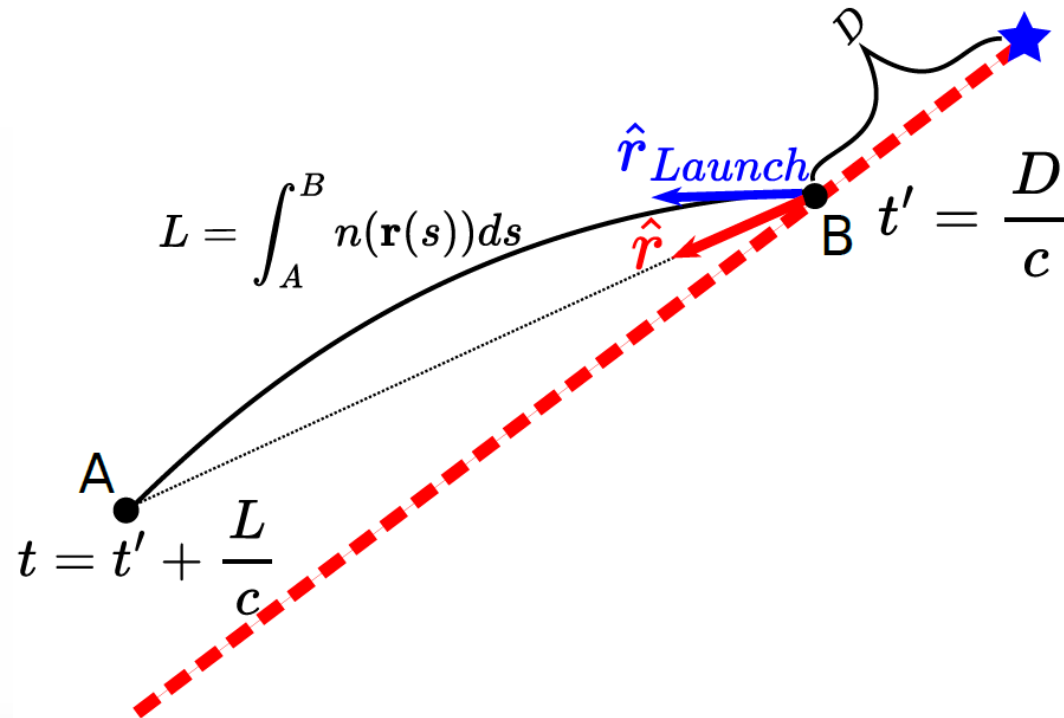
Where for  $\hat{r}$  the straight line vector is used

What is the effect of choosing the straight line vector for this boostfactor?

# HOW CAN WE ADAPT THE BOOSTFACTOR?

Live raytracing for each emitter is too computationally expensive => Tabulation

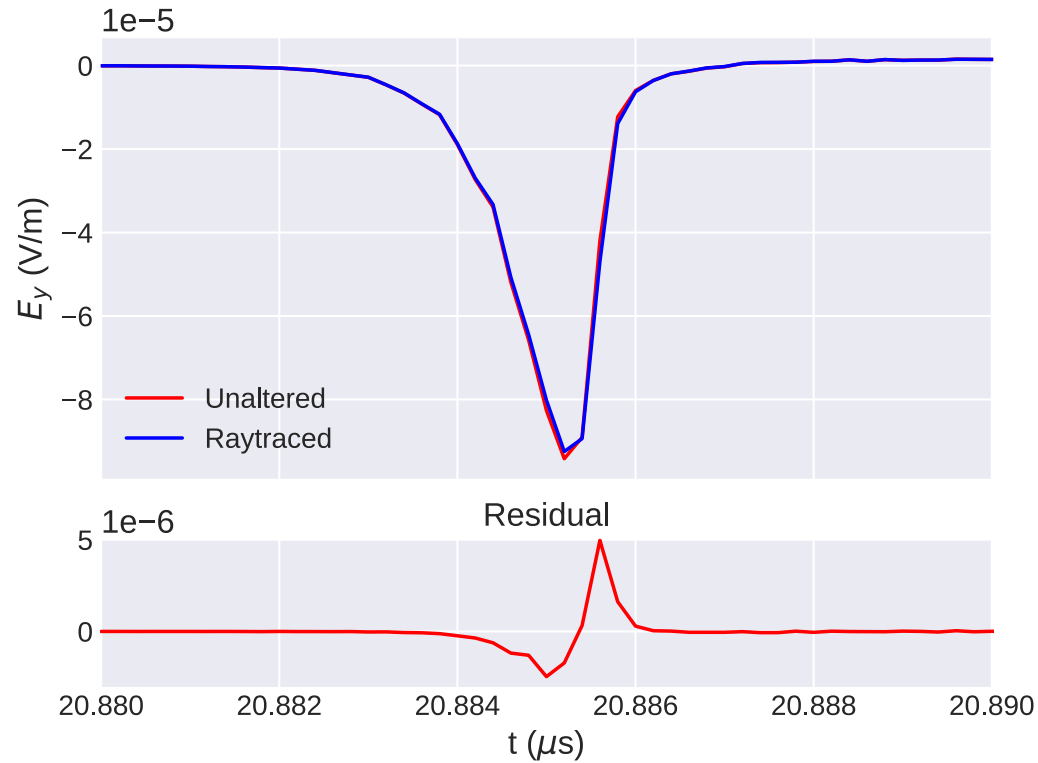
Per observer position: 1 table generated with raytracer that maps  $\hat{r} \rightarrow \hat{r}_{Launch}$



Note: for boostfactor calculations, assume the emitter to be on the shower axis when mapping  $\hat{r} \rightarrow \hat{r}_{Launch}$

# WHAT IS THE EFFECT ON THE COREAS OUTPUT?

CoREAS output are electric field time traces



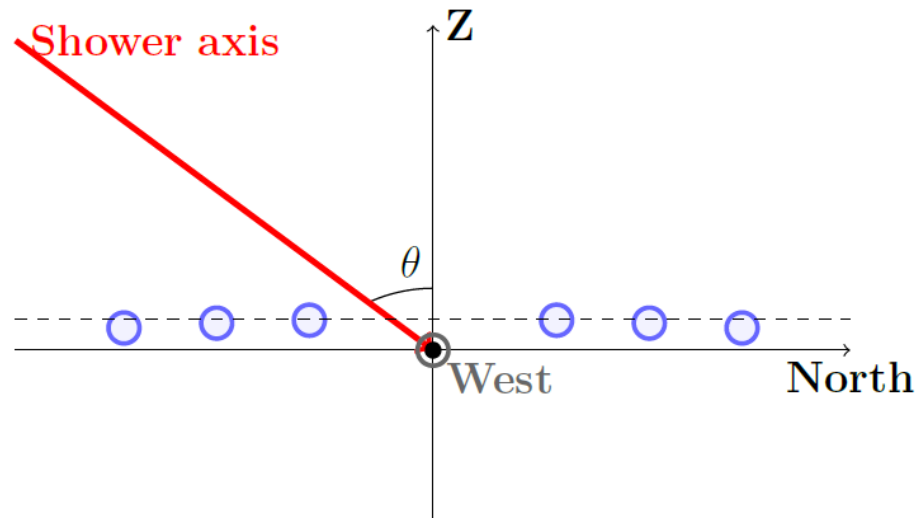
These are limited to single observer positions and divided into different polarisations  
-> hard to get a general picture

# HOW CAN WE GET A MORE GENERAL IDEA OF THE EFFECT?

Look at the fluence:

$$f = \epsilon_0 c \Delta t \sum E_i^2$$

Proton primary  
 $10^{17} eV$   
 $\phi = 0^\circ$   
 $|\vec{B}| = 50 \mu T$   
 $\vec{B} = |\vec{B}| \cdot \vec{1}_z$

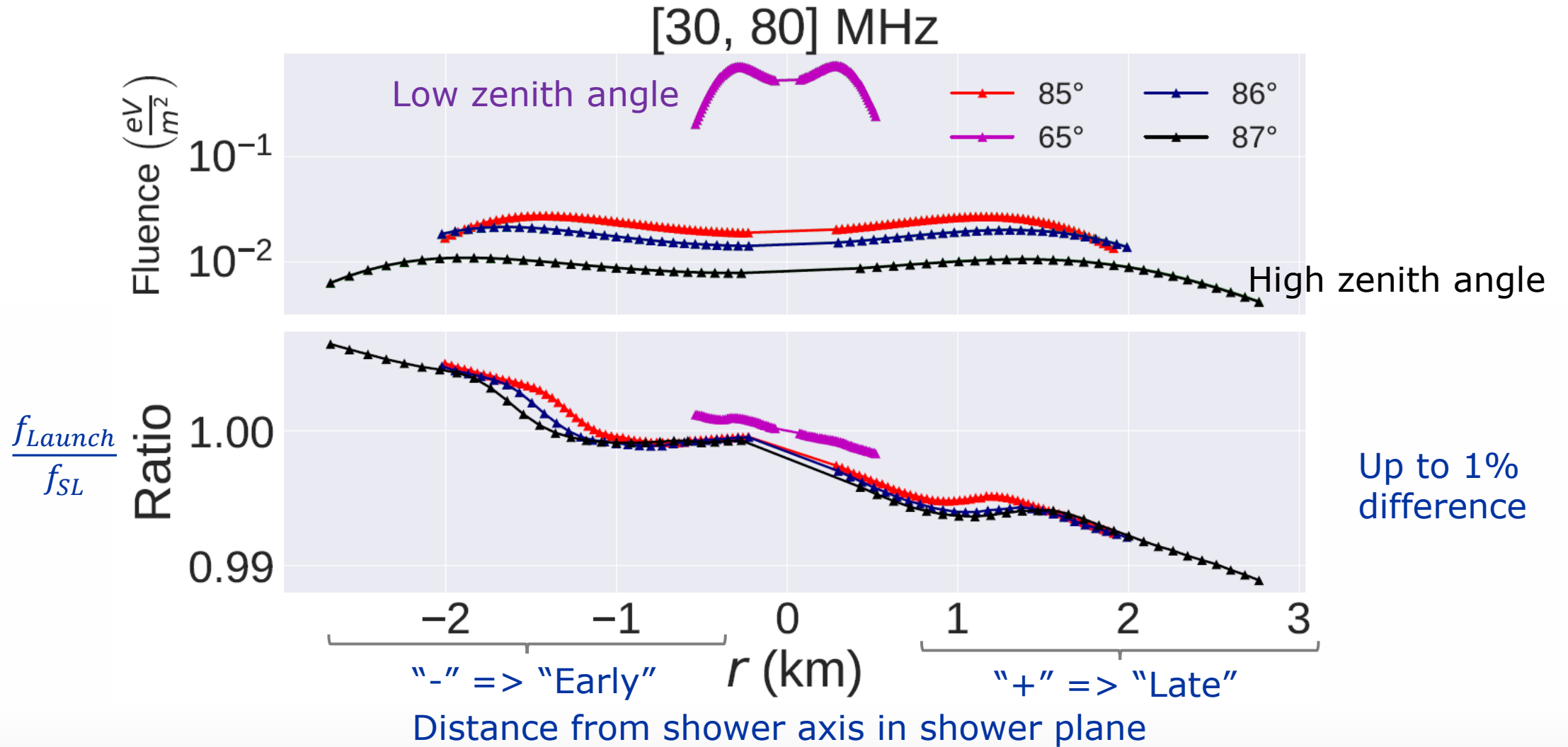


Traces for observers in plane with shower axis -> filter traces for desired frequency range -> calculate fluence -> apply early late correction

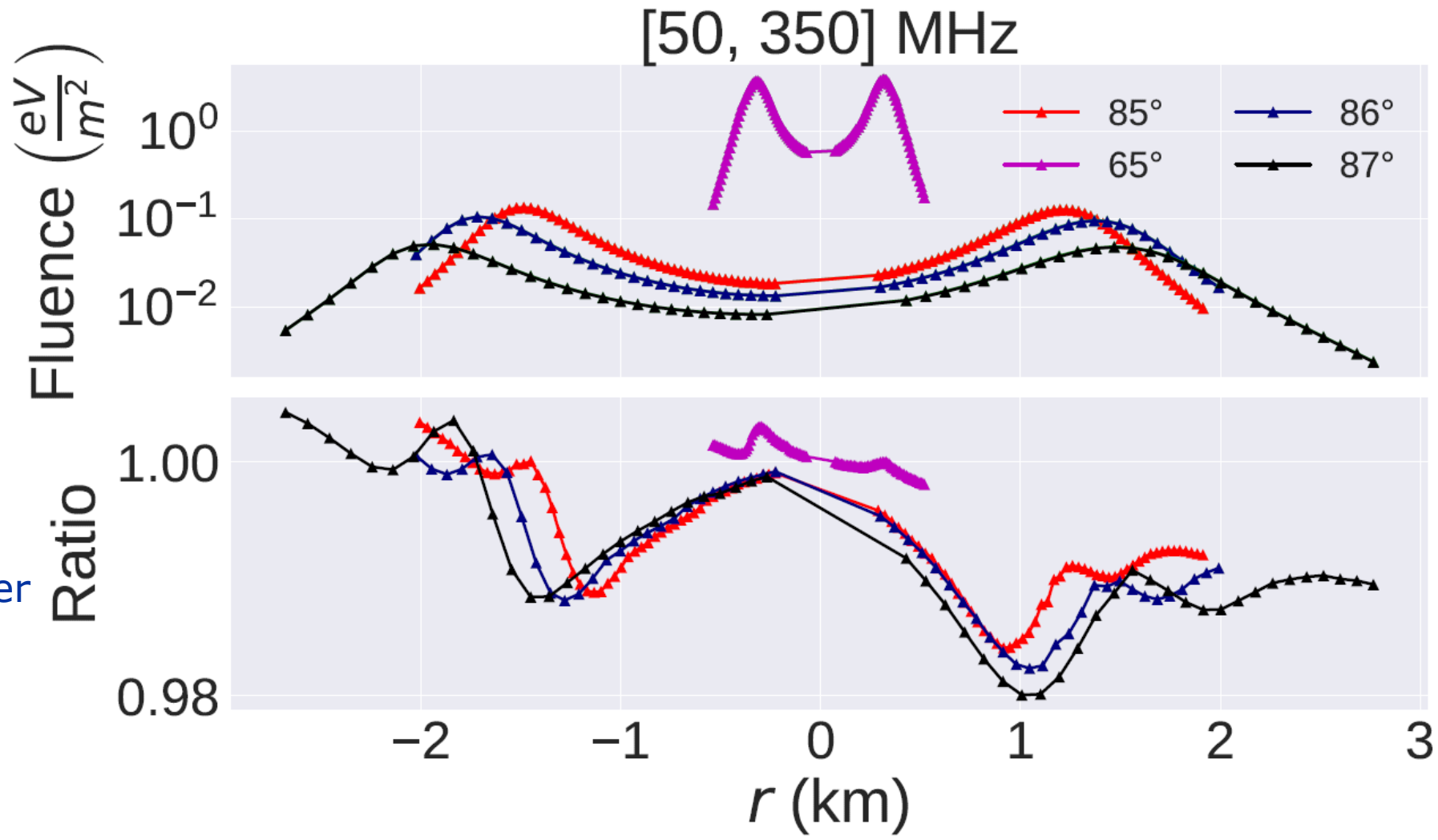
Compare fluence with adapted boostfactor ( $f_{Launch}$ ) to fluence with standard CoREAS ( $f_{SL}$ )



# WHAT DO WE LEARN FROM LOOKING AT THE FLUENCE?



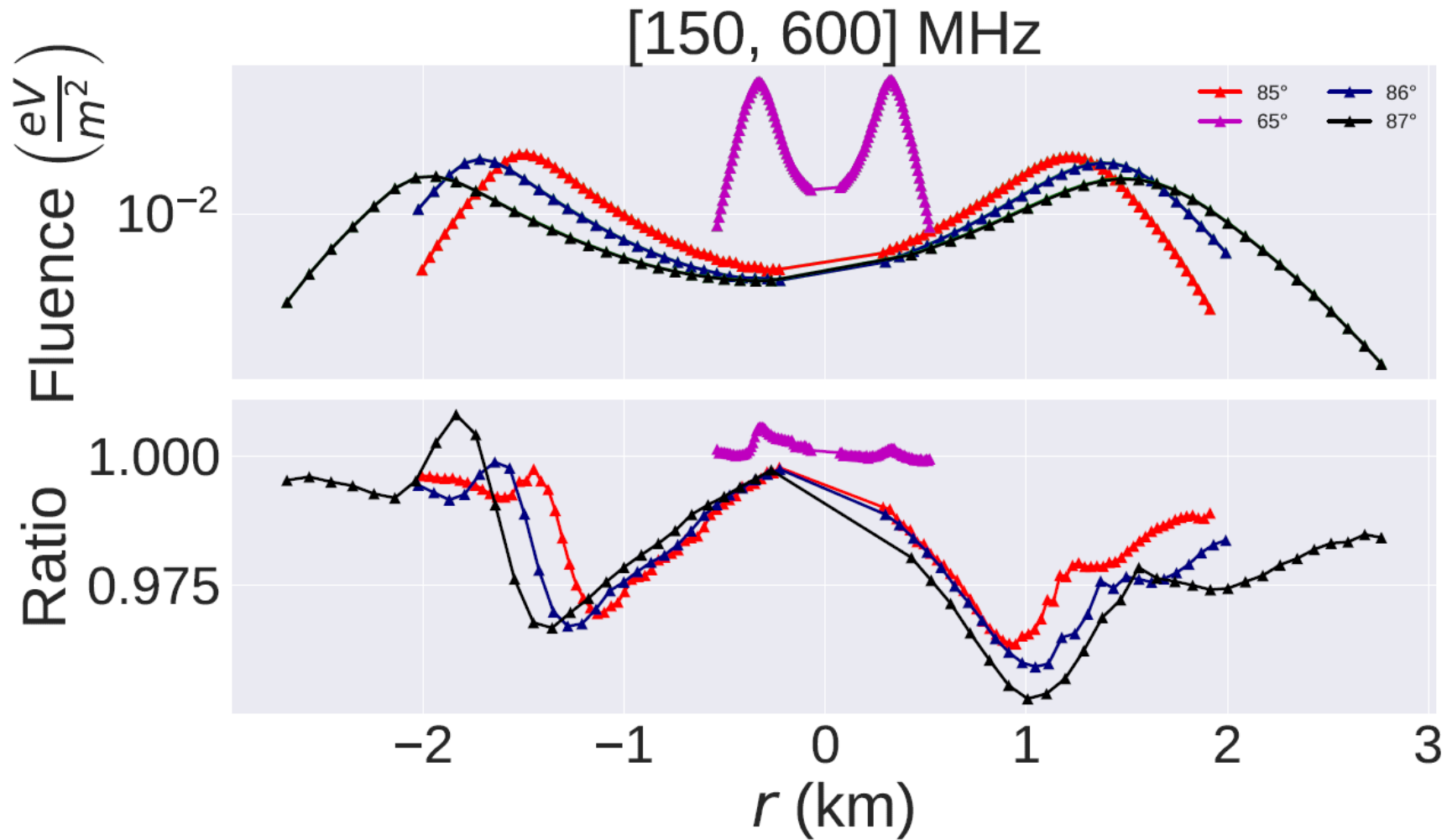
# WHAT DO WE LEARN FROM LOOKING AT THE FLUENCE?



Effect more pronounced towards higher  $\theta$

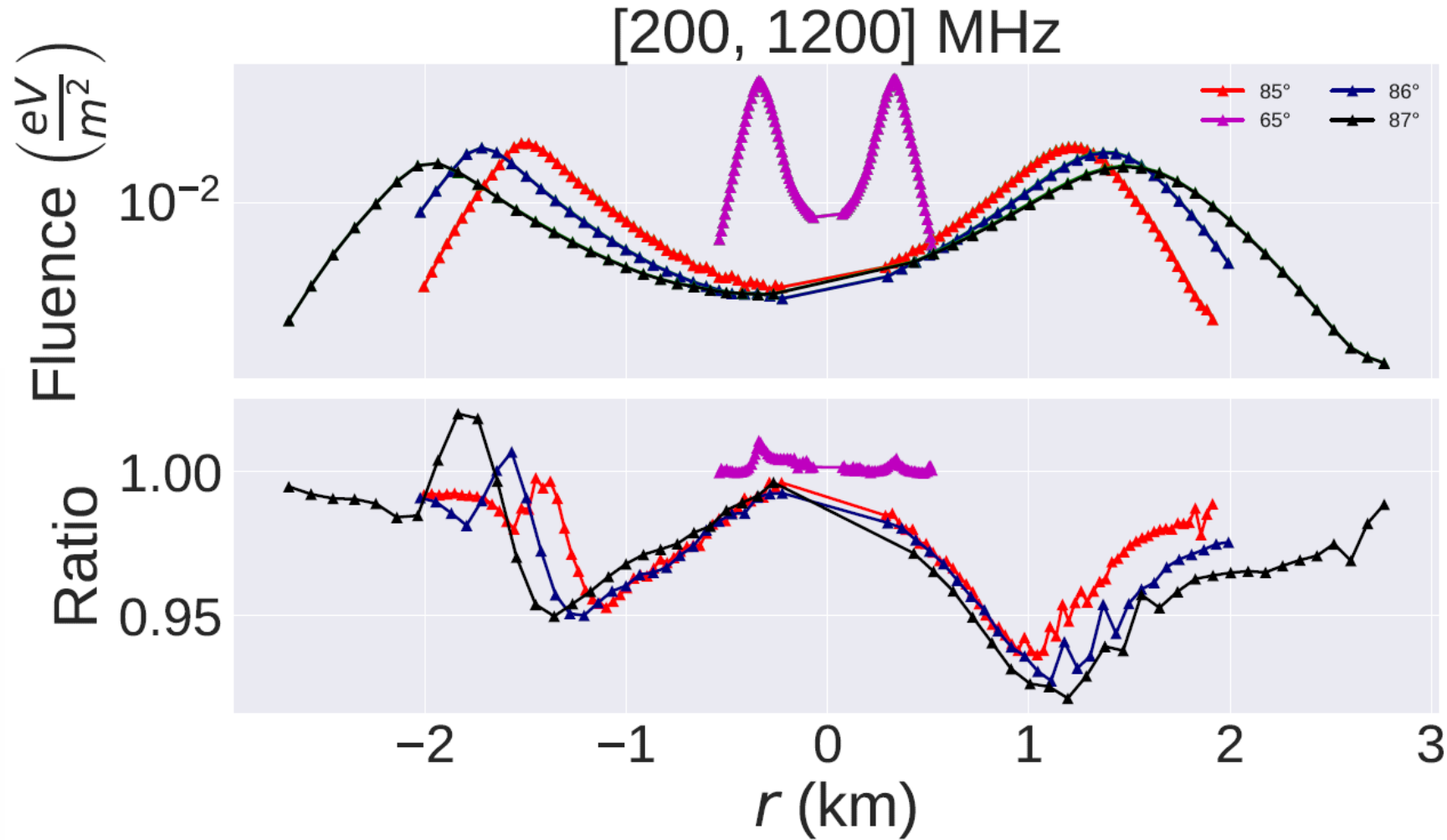
Up to 2% difference

# WHAT DO WE LEARN FROM LOOKING AT THE FLUENCE?



Around 2.5%  
difference

# WHAT DO WE LEARN FROM LOOKING AT THE FLUENCE?

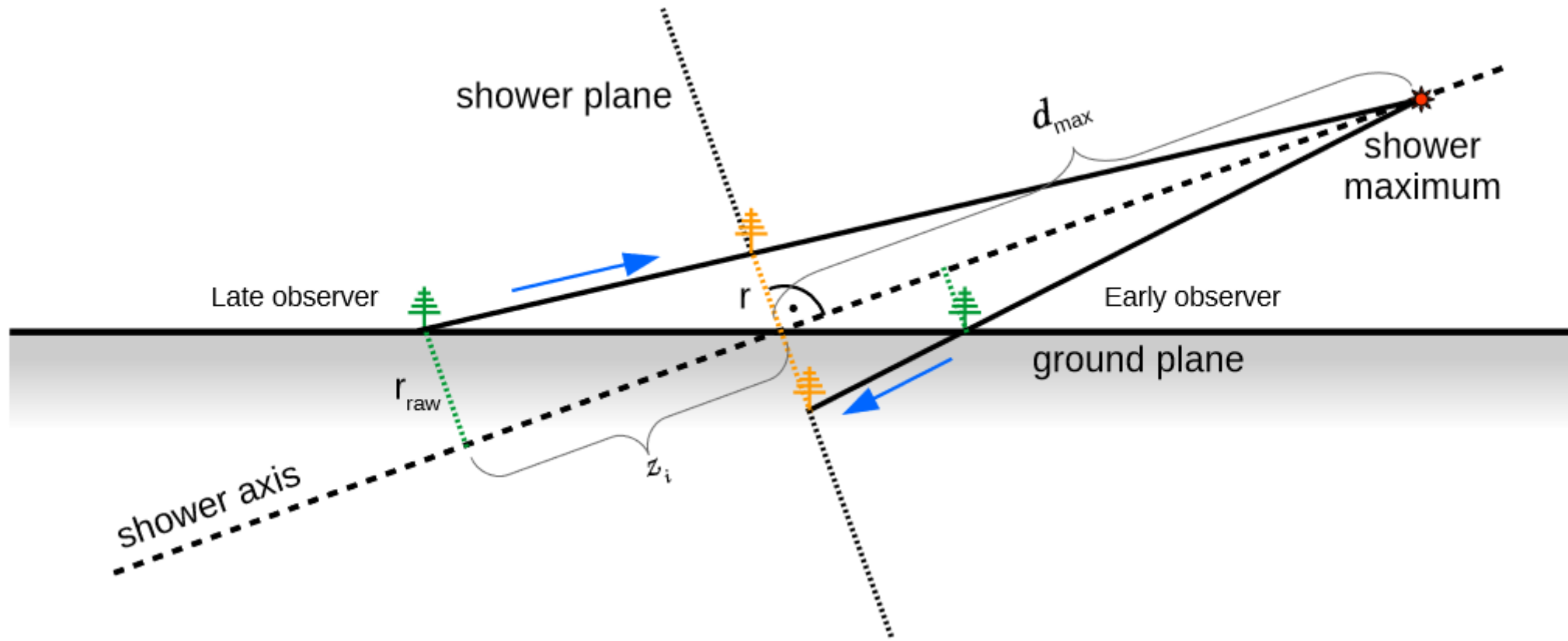


Around 4%  
to 5%  
difference

## CONCLUSION

- The boostfactor can be correctly calculated by using the index of refraction at the emission point and the original launch direction of the ray
- The effect of changing the boostfactor leads to a few % difference in fluence for zenith angles up to  $87^\circ$ , with the effect more pronounced towards higher frequencies.
- Coherence between emitters is most important

BACKUP



$$c_{\text{el}} \equiv \frac{d_{\max} + \vec{x}_i \cdot \vec{e}_v}{d_{\max}} = 1 + \frac{z_i}{d_{\max}}.$$

$$f = f_{\text{raw}} \cdot c_{\text{el}}^2, \quad r = \frac{r_{\text{raw}}}{c_{\text{el}}},$$

