A Multi-Probe Analysis of the 3D Shapes and Non-Thermal Pressure of Galaxy Clusters

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CHEX-MATE

- CHEX-MATE: The Cluster HEritage project with XMM-Newton Mass Assembly and Thermodynamics at the Endpoint of structure formation (CHEX-MATE Collaboration 2021)
 - 3 Msec XMM-Heritage program
 - Planck SZ selected 118 clusters
 - Tier-1: volume-limited sample in the local universe z < 0.2 and dec > 0 $2 \times 10^{14} M_{sun} < M_{500} < 9 \times 10^{14} M_{sun}$
 - Tier -2: sample of the most massive objects to have formed

z < 0.6 and $M_{500} > 7.25 \times 10^{14} M_{sur}$



Science Goals

Shape Measurement

 Unbiased estimates of cluster properties require understanding of shape & orientation of halo

Spherical \rightarrow Triaxial modeling

• Example: WL-derived masses are extremely sensitive to line of sight elongation



Non-thermal Pressure Support

- Assessing equilibrium status of cluster outskirts, where new material is being acccreted
- No large sample studies have probed the cluster outskirts → CHEX-MATE dataset will enable measurement out to R₂₀₀



Weak lensing shear

- Shear maps constructed from archival observations from the Subaru Suprime-Cam instrument
 - We use two component reduced shear maps
- Enables reconstruction of the total mass profile, providing constraints on the mass and concentration



Thermal SZ-y

 SZ Compton-y map from the combination of ACT and *Planck* measurements

$$y \equiv \frac{\sigma_{\rm T}}{m_e c^2} \int_{\parallel} P_e dl = \frac{\sigma_{\rm T} k_{\rm B}}{m_e c^2} \int_{\parallel} n_e T_e dl,$$



X-ray Surface Brightness and Temperature

- X-ray SB observations from XMM-*Newton* in [0.7-1.2] keV range
 - 2D data used in radial region that encloses 80% of the emission
 - 1D data used in exterior region to mitigate biases from gas clumping
- Spectroscopic temperature measurements constructed via spectroscopic fits to SB data
 - Assume ICM is ideal gas to estimate electron temperature



$$SB = \frac{1}{4\pi(1+z)^3} \int_{\parallel} n_e^2 \Lambda_{\text{eff}}(T_e, Z) dl$$

$$T_{\rm sp} = \frac{\int W T_e dV}{\int W dV} \text{ keV}; W = \frac{n_e^2}{T_e^{3/4}},$$



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Triaxial Modeling



Assumed thermodynamic profiles of ICM, the electron density (n_e) and electron pressure (P_e) are represented as functions of the ellipsoidal radius:

$$\zeta^2 = \frac{x_1^2}{q_1^2} + \frac{x_2^2}{q_2^2} + x_3^2$$

- SZ and X-ray SB redundantly probe the LOS extent of the ICM (ICM temp measured from X-ray)

$$S_X \propto \int n_e^2 \Lambda(T_e, Z) dl + B_{SZ} \propto \int n_e T_e dl \longrightarrow \Delta l \sim \frac{B_{SZ}^2 \Lambda(T_e, Z)}{S_X T_X^2} \sim \frac{B_{SZ}^2}{S_X T_X^2}$$

Mass Reconstruction



Projection

npton-y 3

-2

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x [arcmin]

y [arcn

6

4

2

x [arcmin]

0 -2 -4 -6

v [arcmin]

• To make models of our 2D observables, we must project the assumed 3D profiles describing the signal along the LOS:

$$F_{2D}(x_{\xi}; l_p, p_i) = 2l_p e_{\parallel} \int_{x_{\xi}}^{\infty} F_{3D}(x_{\xi}; l_s, p_i) \frac{x_{\xi}}{\sqrt{x_{\xi}^2 + x_{\xi}^2}} dx_{\xi}$$

$$\Psi = \frac{D_{LS}}{D_L D_S} \frac{2}{c^2} \int \Phi dl$$

$$P_e \quad n_e \quad \Phi$$

$$\gamma_1 = \frac{1}{2} (\Psi_{11} - \Psi_{22}) \text{ and } \gamma_2 = \Psi_{12} = \Psi_{21}$$

$$Model (SZ, convolved) \quad e^{-5} \int_{a}^{b} \int_{a}^{b} \frac{W_{2}}{2} \int_{a}$$

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ture [keV]

-10

-5 0

x [arcmin]

10

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x [arcmin]

2

-2 -4 -6

s) [cou

-3.50

3.75

-0.1

-0.2

-0.3

5 10

Non-Thermal Pressure Calculation

From gas analysis:

Can calculate elliptical gas density profile:

 $\rho_{gas} = \mu_e m_p n_e$

Can calculate elliptical gas pressure profile:

 $P_{th} = P_e \frac{\mu_e}{\mu}$

From addition of WL analysis:

Can calculate gravitational potential on triaxial ellipsoid.

Assume generalized HSE

$$\nabla P_{\rm tot} = -\rho_{\rm gas} \nabla \Phi,$$

Undo this numerically to get the total pressure needed to offset gravity

Non-thermal pressure:

Spherically average the total pressure and thermal pressure, then:

$$P_{nt} = P_{tot} - P_{th}$$

This calculation is computationally expensive, so it is done outside the fit

Demonstration on A1689

Multiprobe Fit

- Geometrical constraints are consistent with what is found in the literature
 - $\cos\theta = 0.99 \Rightarrow A1689$ is almost perfectly aligned with the line of sight
 - $R_{LP} = 1.27 \Rightarrow A1689$ is elliptical and elongated along the line of sight



Radial Profiles



We find good agreement between the best fit model and the input data

- Limited d.o.f. in model → two independent thermodynamic profiles shapes (density and pressure) must simultaneously describe three observables
 Higher S/N SZ data primarily constrain P_e
 Elongation ensures
 - normalization of fitted temp profile is in agreement with obs data

WLResults

Prior spherical fits to A1689 suggested a high c₂₀₀ for this system

 High c₂₀₀ retained in triaxial fit → may be intrinsic rather than a result of projection

Higher mass obtained in spherical fit agrees with expectations from fitted geometry given LOS elongation = 1.27



Non-thermal Pressure Fraction



We obtain a non-thermal pressure fraction with ≤10% uncertainty

- Chappuis et al 2025 is a 1D analysis that uses the same CHEX-MATE data but different modeling formalism
- Green and yellow lines are independent analyses of simulated clusters showing ensemble average profiles

Conclusions

We are capable of measuring spherical and triaxial masses

- Will investigate the spherical-triaxial mass bias in the full CHEX-MATE sample 0
- We obtain a radial profile of the non-thermal pressure fraction with $\leq 10\%$ uncertainty
 - Will apply pipeline to full CHEX-MATE sample to obtain an ensemble average radial profile 0

Next steps:

- Apply pipeline to sample of \sim 50 simulated clusters from The300 to investigate any bias 0 introduced by the fit
- Apply pipeline to sub-sample of CHEX-MATE clusters 0



II. Combined Gas and Dark Matter Analysis from X-ray, SZE, and WL

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Extra Slides

Mock Data Fits

Fitting to toy model data generated using model equations tests how well the pipeline returns known input parameters

| ${ m M}_{ m _{200}}[10^{14}{ m M}_{ m _{sun}}]$ | C ₂₀₀ | |
|---|------------------|---|
| 8 | 3.8 | - |
| 10 | 3.8 | |
| 15 | 3.8 | |

| ${ m M}_{ m 200}~[10^{14}~{ m M}_{ m sun}]$ | C ₂₀₀ |
|---|------------------|
| 8 | 2 |
| 10 | 2 |
| 15 | 2 |

| | ${ m M}_{ m 200}~[10^{14}~{ m M}_{ m sun}]$ | C ₂₀₀ |
|---|---|------------------|
| | 8 | 10 |
| * | 10 | 10 |
| | 15 | 10 |

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| q ₁ | 0.60 |
|-----------------------|----------------------|
| q ₂ | 0.75 |
| cosθ | 0.6 |
| φ | -25 |
| Ψ | 60 |
| n _o | 0.01 cm ³ |
| ζ_{c} | 200 kpc |
| ζ _t | 2.5 Mpc |
| β _e | 0.65 |
| Ŋ _e | 0.60 |
| γ _e | 4.2 |
| P ₀ * * | 30 |
| α _p | 0.8 |
| | |

Mock Data Fits



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Mock Data Fits



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Gas Mass Fraction



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Science Goal: Shape Measurement

- What is the distribution of three-dimensional shapes of galaxy clusters?
 - Knowledge of the mass and concentration of cluster crucial for understanding formation & evolution
 - Unbiased estimates require understanding of shape & orientation of halo

Spherical \rightarrow Triaxial modeling

- Cosmological models make strong prediction for the shape of DM halos (e.g., Yoshida et al. 2000)
 - CLASH results (Sereno et al. 2018) potentially suggest more extreme axial ratios compared to simulations
 - CLASH measurement has lower $q_{mat1} \rightarrow$ more recent formation time than in simulation?
 - Larger axial ratios could potentially point to a non-zero self-interaction cross-section for dark matter



Triaxial Gas Analysis

- Assume the following model profiles:
 - Electron density profile (Vikhlinin et al. 2006; Ettori et al. 2009)

 $n_e(\zeta) = n_0 \left(\frac{\zeta}{\zeta_c}\right)^{-\eta_e} \left[1 + \left(\frac{\zeta}{\zeta_c}\right)^2\right]^{-3\beta_e/2 + \eta_e/2} \left[1 + \left(\frac{\zeta}{\zeta_t}\right)^3\right]^{-\gamma_e/3} \quad \text{cm}^{-3}$

• Gas pressure profile: gNFW (Nagai et al. 2007; Arnaud et al. 2010)

 $\frac{P_e(x)}{P_{500}} = \frac{P_0}{(c_{500}x)^{\gamma_p} [1 + (c_{500}x)^{\alpha_p}]^{(\beta_p - \gamma_p)/\alpha_p}} \qquad x = \zeta/R_{500}$

SZ and X-ray SB redundantly probe the LOS extent of the ICM

$$J_X \propto \int n_e^2 \Lambda(T_e, Z) dl + B_{SZ} \propto \int n_e T_e dl \longrightarrow \Delta l \sim \frac{B_{SZ}^2 \Lambda(T_e, Z)}{S_X T_X^2} \sim \frac{B_{SZ}^2}{S_X T_X^2}$$

Triaxial Total Mass Reconstruction

- We assume an NFW density profile to model the total matter distribution and that the gravitational potential resulting from it is elliptically symmetric
 - Allows us to assume $q_{pot} = q_{ICM}$ (motivated by simulation)

$$\Phi(\zeta) = 4\pi G \Delta^{-1} \rho(\zeta') = 4\pi G \Delta^{-1} \left[\frac{\rho_s}{(\zeta'/\zeta_s)(1+\zeta'/\zeta_s)^2} \right]$$

- Define μ_{200} and γ_{200} as elliptical analogs to M_{200c} and c_{200c} . These parameters are defined exactly w.r.t the axial ratios of the matter distribution
 - For computational efficiency, make approximation in connecting μ_{200} and γ_{200} to ζ_s and ρ_s that $q_{mat} \sim q_{pot}$

$$\zeta_s = \frac{\zeta_{200}}{\gamma_{200}} \cong \frac{1}{\gamma_{200}} \left[\frac{3}{800\pi\rho_{cr}} \frac{1}{q_{\text{pot},1}q_{\text{pot},2}} \mu_{200} \right]^{1/3} \qquad \rho_s = \frac{200}{3} \rho_{cr} \frac{\gamma_{200}^3}{\ln(1+\gamma_{200}) - \gamma_{200}/(1+\gamma_{200})}$$

Convert M_{200c} and $c_{200c} \rightarrow \mu_{200}$ and γ_{200} using a look-up table

Science Goal: Shape Measurement

- Unbiased estimates of cluster properties require understanding of shape & orientation of halo
 - Spherical \rightarrow Triaxial modeling
 - Cluster abundance cosmology

Ο

- Slight " S_8 tension" between cluster measurements and other probes such as the CMB, with cluster measurements producing lower values of S_8
 - Most likely explanation is mass calibration
 - Simulations used to quantify the mass bias due to the assumption of spherical symmetry to calibrate WL-derived masses
 - An observationally derived benchmark to compare with simulations does not yet exist



