

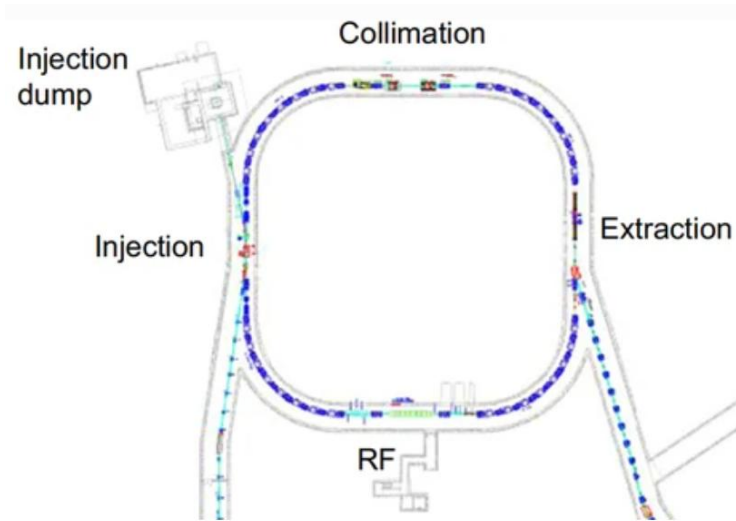
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# Proton Driver Design

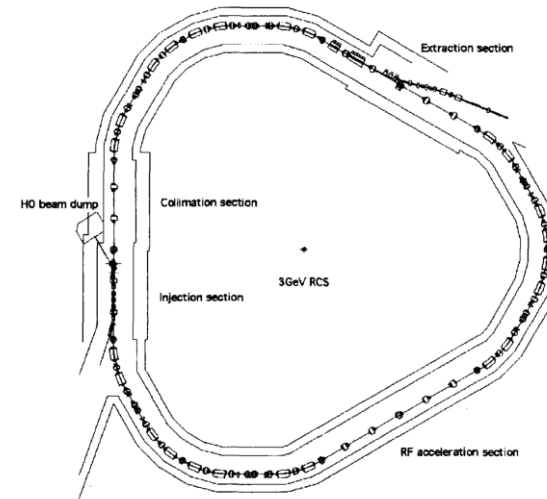
Jeffrey Eldred

for **US Muon Collider Accelerator School**  
**University of Chicago, August 3-6th 2025**

# Existing Proton Machines (~2007)



- 1) Oak Ridge SNS  
Neutron Spallation Source  
1.0 GeV AR, ~1.4MW  
1.3-GeV AR, -> 2.8MW



- 2) J-PARC RCS  
Meson/Muon Production  
0.2-3 GeV RCS, ~0.5MW  
0.4-3 GeV RCS, -> 1.4MW

**Both MW-class facilities that serve as a model for proton driver...**  
Except we need 2-4MW, 5-12 GeV, 1-3ns pulses (from ~us).

# Proton Driver as Muon Collider R&D Topic

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Candidly, proton driver design is not biggest factor for technical risk, muon collider performance, or overall cost.

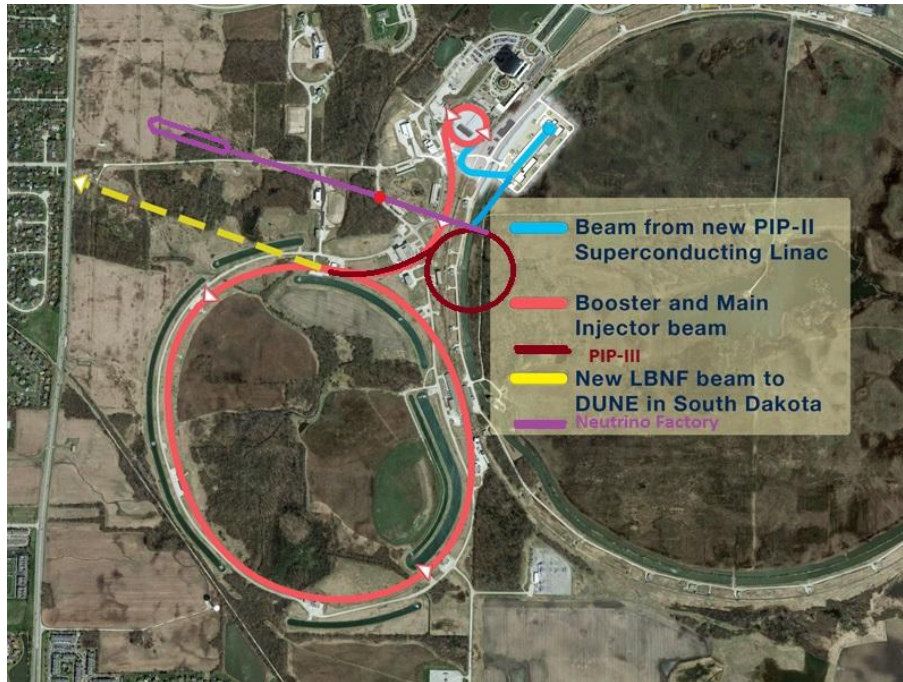
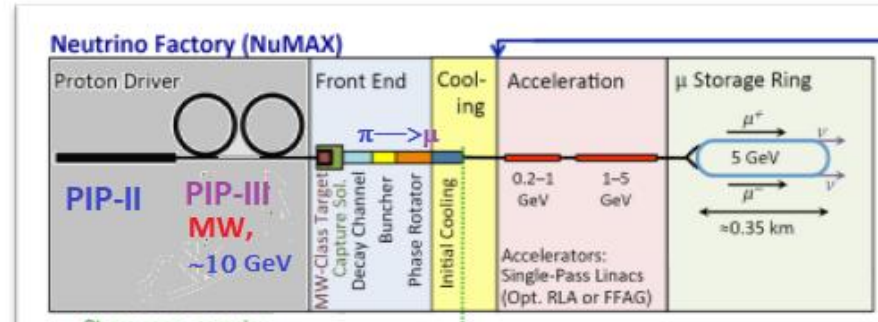
- Closest to existing technology and operational experiences.

So why study proton driver as a muon collider R&D topic?

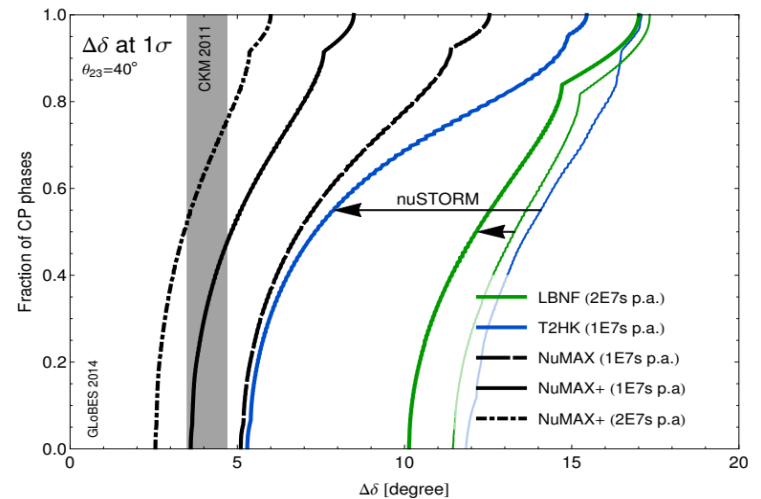
- We are still working towards a self-consistent baseline scenario.
- Any technical risk that we can offload onto the proton driver we will be highly cost-effective.
- Proton driver is chronologically first by virtue of being upstream. For the hypothetical Gantt chart of building a Muon Collider, it needs the highest level of technical readiness.
- Proton drivers present opportunities for HEP staging.

# Possible Proton Driver HEP Staging

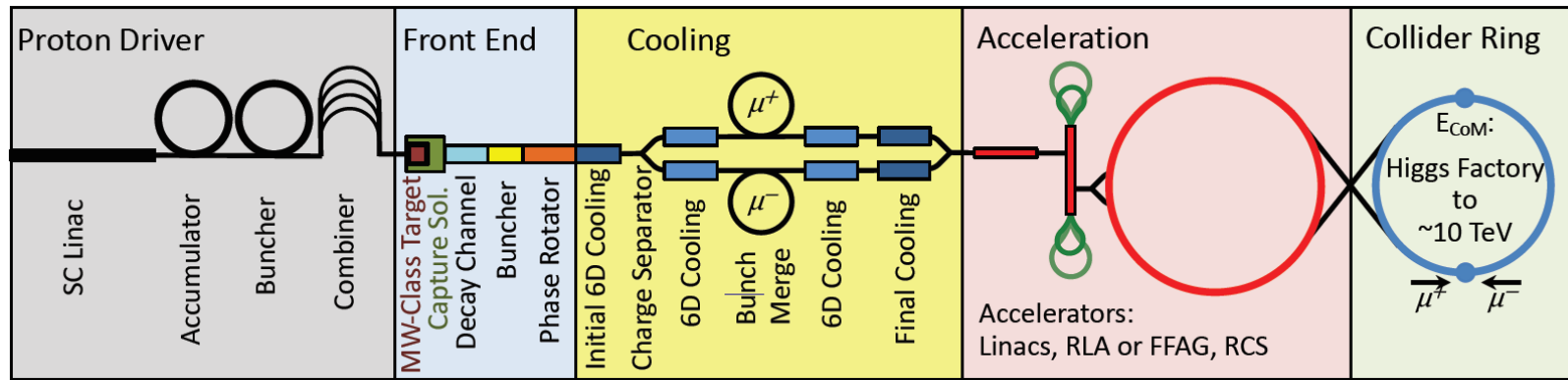
- 1) Built Proton Driver
- 2) Built separate muon ring for neutrino program.
- 3) Continue MC construction.



Vastly improves  $\delta_{CP}$  resolution.  
Unitarity tests of PMNS matrix.



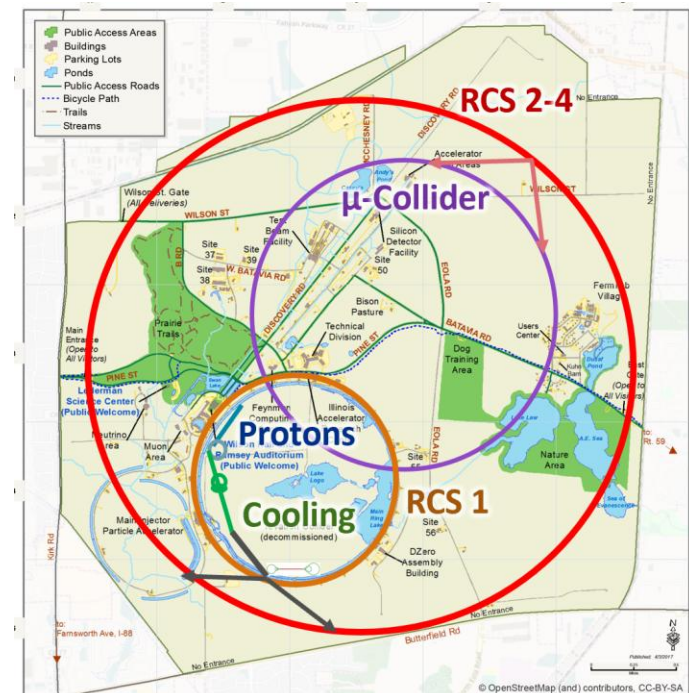
# Proton Driver on the MAP



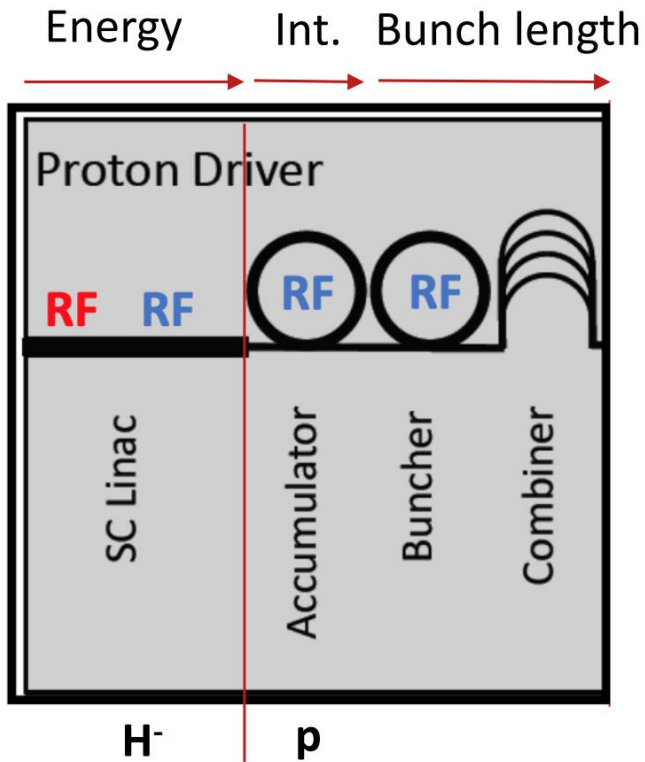
The **Proton Driver** determines the proton energy, pulse intensity, beam power, spot size on target, bunch length on target for muon production stage.

There in turn feed in the x, y, and t phase-space of the initial cooling stage, muon rate, and pulse repetition rate.

By the end, just luminosity and rep. rate.



# Proton Driver Overview



**H- Linac**, accelerates up to 5-12 GeV

**Accumulation Ring (AR)**, collects  $H^- \sim 3e14$  particles and paint into stable bunches.

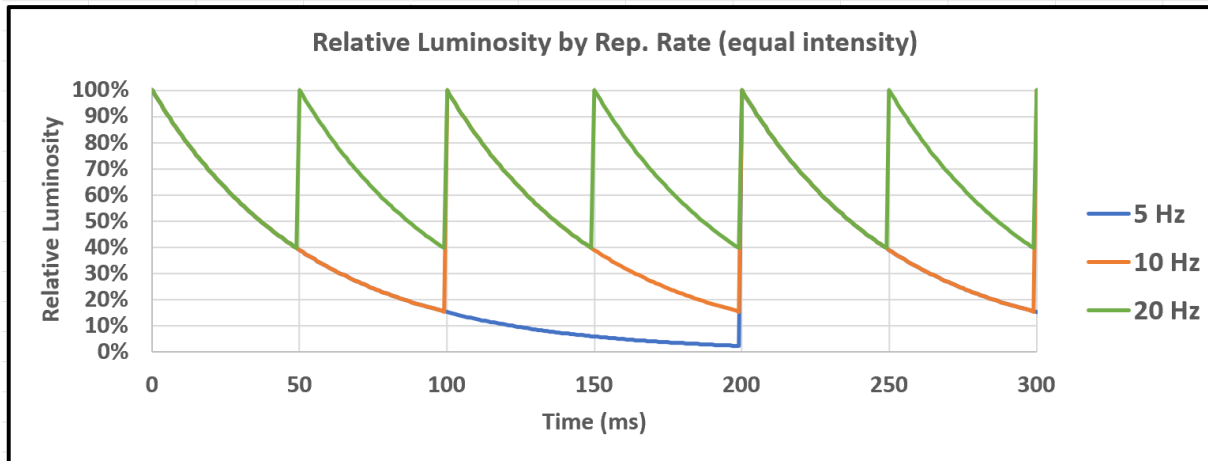
**Compressor Ring (CR)**, performs bunch rotation to extract 1-3ns bunches.

**Combiner**, ensures that all (typically 1-6) bunches converge on the target simultaneously.



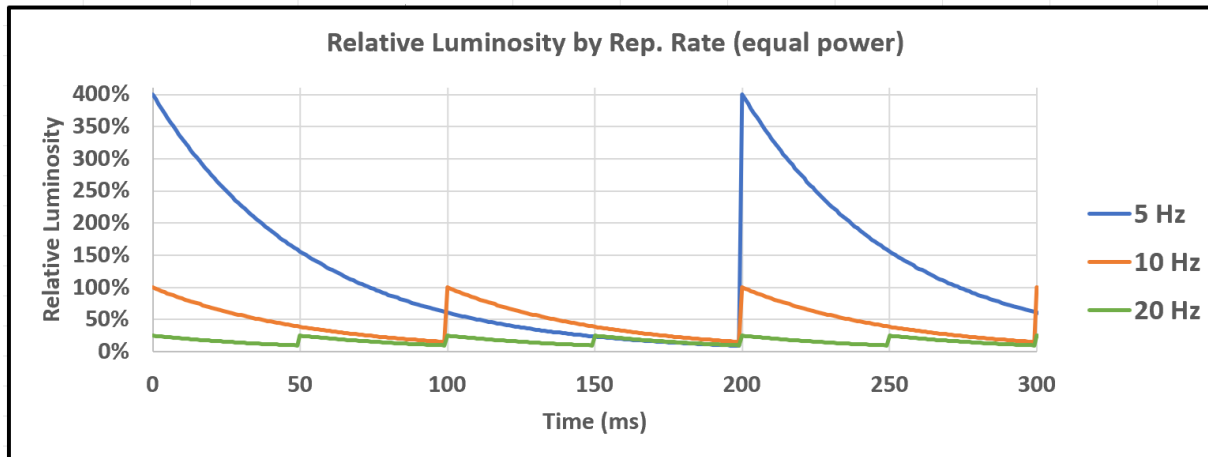
# Muon Collider Proton Driver Parameters

**Rep Rate 5-10 Hz:** At 5 TeV beams, muon decay time is 0.1s in the lab frame. Luminosity depends on bunch intensity squared (neglecting beam quality losses).



Higher rep rates are better, given the same number of particles per pulse.

Above 20 Hz, diminishing returns for luminosity.



However, lower rep rates are better for a fixed beam power, because each pulse is more intense.

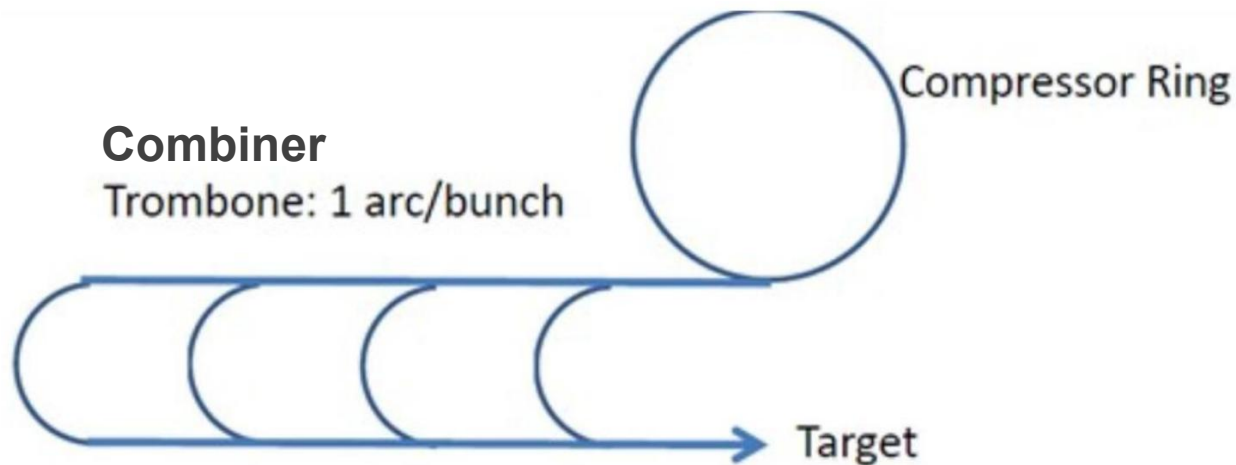
**10 Hz** may be a good balance of considerations.

# Muon Collider Proton Driver Parameters

**Beam Energy 5-12 GeV:** Muon yield per kW on target is relatively flat from 5-10 GeV, smaller beyond that. Using **Fermilab 8 GeV** is a reasonable choice.

**Bunch Length 1-3 ns:** The time structure of the proton beam gets largely passed on to the pion/muon beam produced for the target, so **1-3ns** bunch length alleviates the 6D cooling requirements as much as possible.

**Proton Beam Power 1-4 MW:** To achieve overall luminosity goals ( $20 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ ). At 8-GeV and 10 Hz, the corresponding proton pulse is **80-320e12 protons**. However, split among four bunches sent to combiner, **20-80e12 protons/bunch**.





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# Linac & H- Injection

# Linac Macropulse Structure

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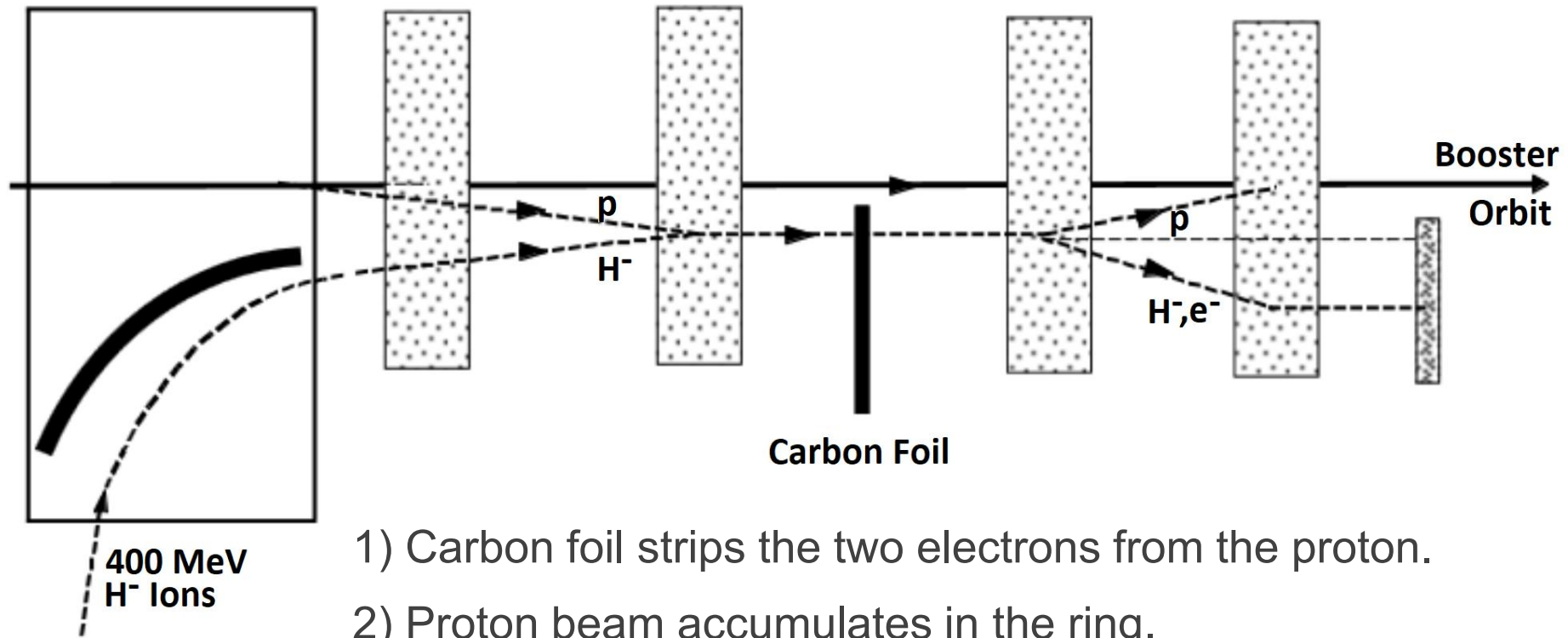
For a pulsed linac klystrons are used for RF sources and for a continuously operating linac solid state amplifiers are used.

Pulsed power is much more affordable and much closer to what the linac needs to deliver  $160\text{-}320\text{e}12$  every 5-20 Hz.

Recent FNAL ACE design exercise for an 8-GeV Linac

- Injector chain begins with PIP-II linac with a beam current upgrade.
- **$10\text{ Hz} \times 5\text{ mA} \times 2\text{ ms} \times 8\text{ GeV} = 0.8\text{ MW}$**
- 1.3 GHz ILC style cavity, LCLS-II style cryomodule, E-XFEL style klystron RF power sources 1.5-2ms.
- Use higher linac current of  **$6\text{-}25\text{mA}$** , that becomes  **$1\text{-}4\text{ MW}$** .
- Might be better to let PIP-II serve 2mA CW users, and build new injector.

# Charge Exchange Injection

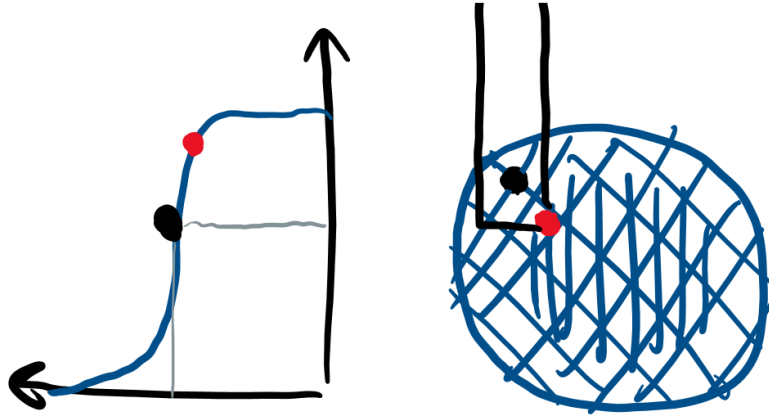


- 1) Carbon foil strips the two electrons from the proton.
- 2) Proton beam accumulates in the ring.
- 3) Unstripped  $\text{H}^-$  ions are sent to an absorber.

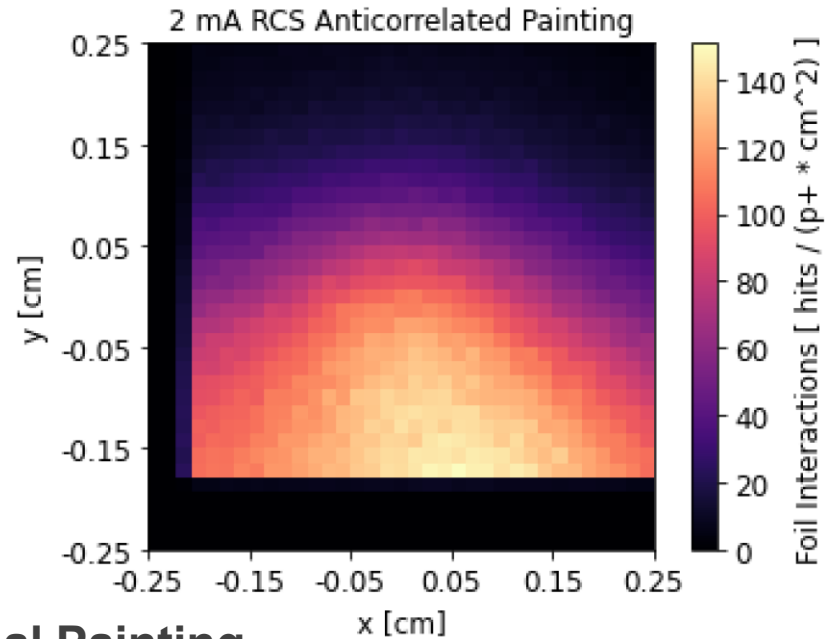
**Linac current is concentrated 100s to 1000s times over!**

# Longitudinal & Transverse Painting

## Beam Distribution during Injection

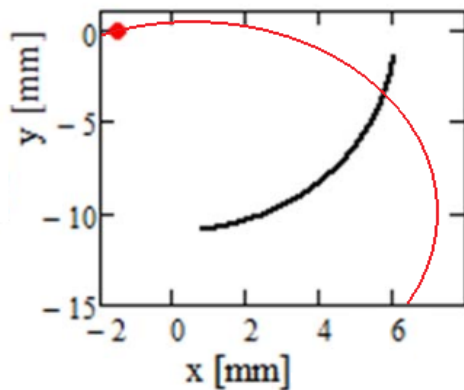


## Foil Interaction Distribution

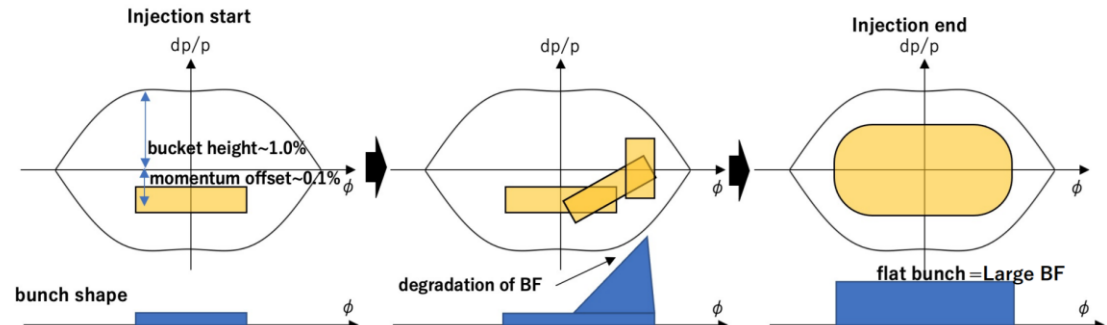


## Anti-correlated Painting Injection

Circulating orbit relative to injection orbit



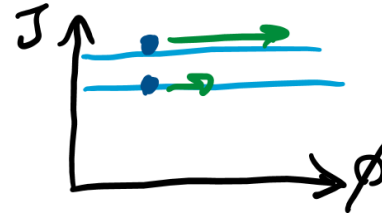
## Longitudinal Painting



# Equilibrium Charge Distributions

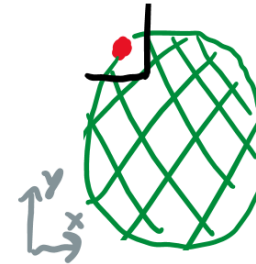
Due to nonlinear decoherence effects, particles tend to become uniformly distributed on surfaces of constant action  $J_x$  and  $J_y$ .

$$\rho_{\perp}(x, p_x, y, p_y) \rightarrow \rho_{\perp}(J_x, J_y)$$



Due to nonlinear coupling effects (inseparable transverse degrees of freedom), particles may also become uniformly distributed on equipotential surfaces of the Hamiltonian  $H$ .

$$\rho_{\perp}(J_x, J_y) \rightarrow \rho_{\perp}(H)$$



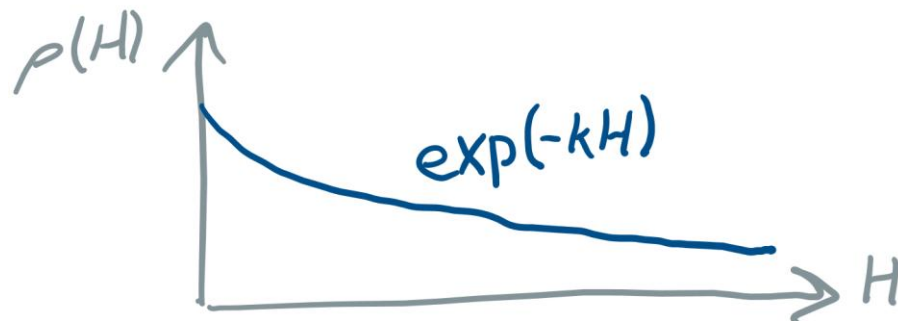
In the limit of weak space-charge, we can write down the single-particle Hamiltonian, consider space-charge as a perturbation, find an equilibrium.

In general however, writing a self-consistent space-charge distribution, not already at equilibrium is challenging.

One special case is the KV distribution, which is an important case study.

# Some Particle Distributions

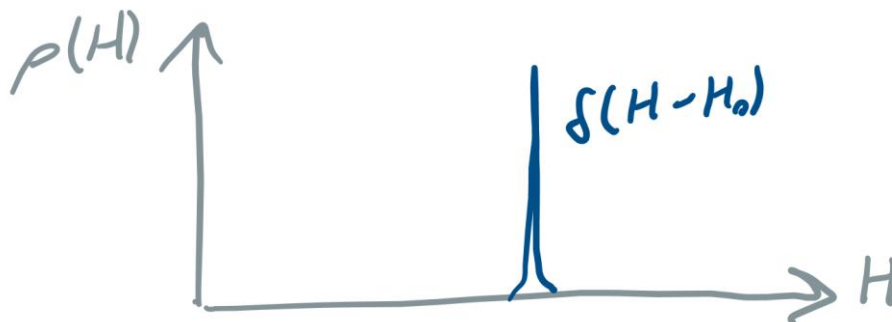
Gaussian:



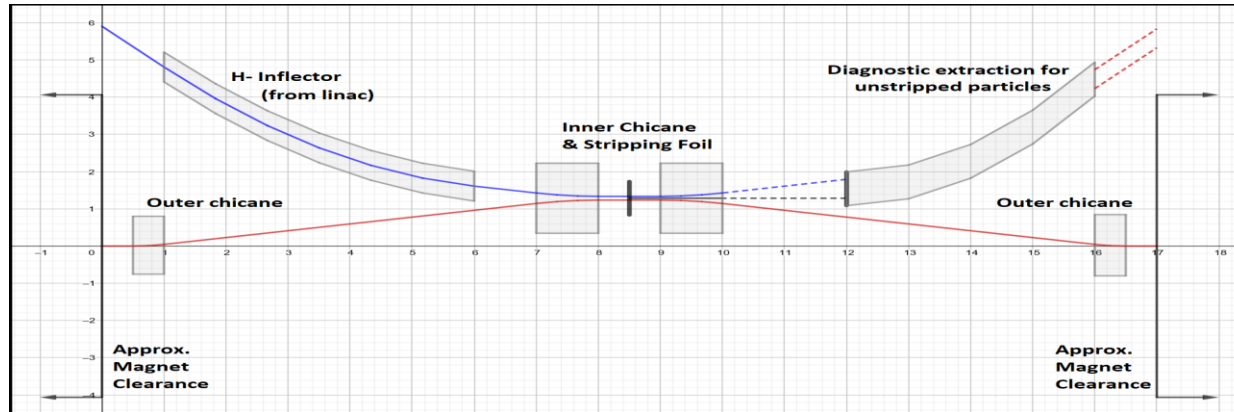
Waterbag:



KV:



# H- Injection Foil Challenges



**Long Injection Chicane** to avoid early H- stripping, worse with energy.

**Extraction to absorber**, to remove unstripped H-/H<sup>0</sup> particles

**Foil Failure due to beam heating**, Foil failure at 2000-2300 K

**Circulating Proton Beam Scatters off Injection Foil**, hard to absorb

**At GeV-scale energy helps, but at higher energy it hurts.**

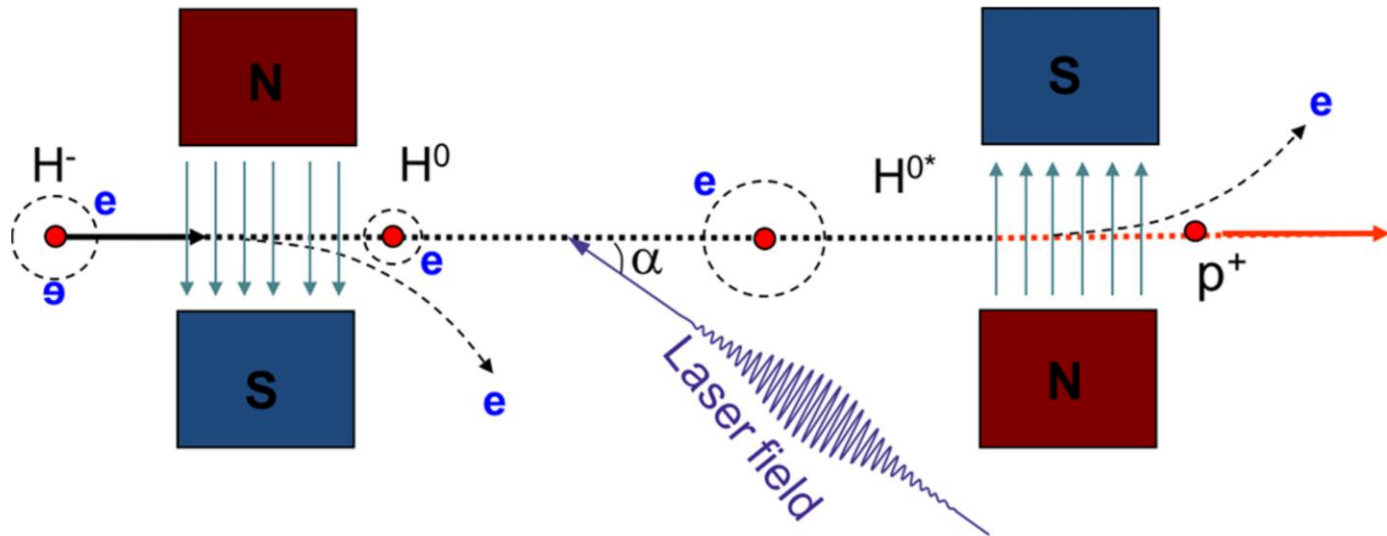
- Proton scattering persists, absorbing particles is much harder.



# Laser Stripping injection for H<sup>-</sup> Injection

Is foil stripping even feasible for a 4MW proton driver AR?

In laser-assisted charge-exchange (LACE), H<sup>-</sup> is Lorentz-stripped to H<sup>0</sup>, H<sup>0</sup> excited with a laser, and then H<sup>0\*</sup> stripped to proton with a second magnet.



If laser stripping efficiency is 95%, 200 kW is being directed to a beam dump!

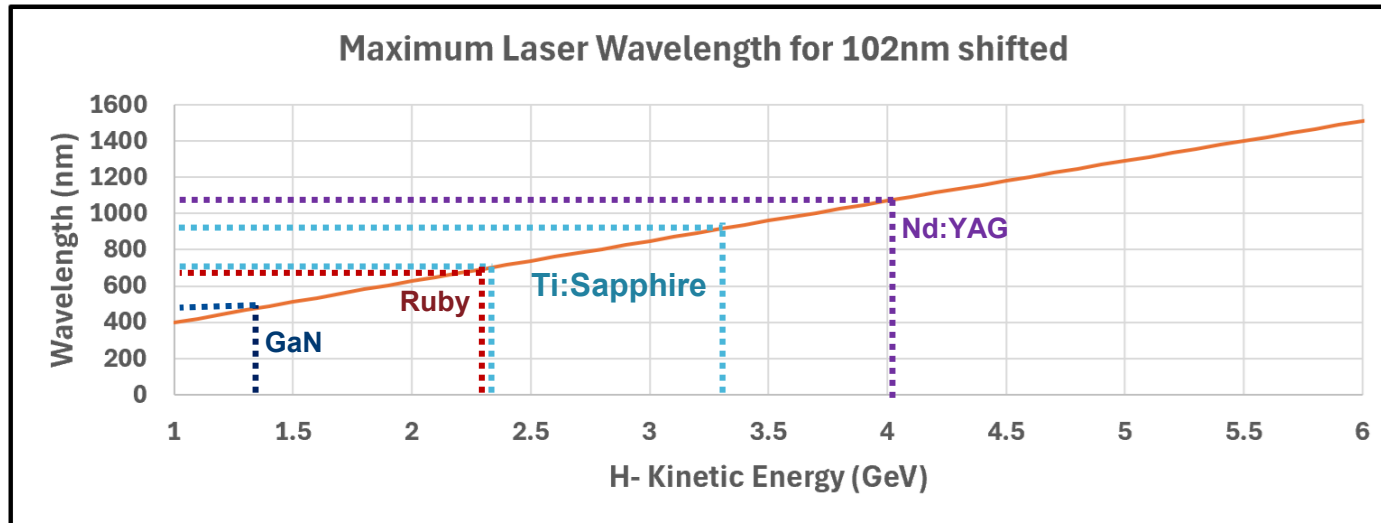
This technology is being developed at SNS and J-PARC

# Energy Dependence on H- Laser Stripping injection

Laser wavelength is Doppler-boosted to target Lyman series

$n=1 \rightarrow n=3$ ,  $\lambda_0 = 102 \text{ nm}$  line.

$$\lambda_{laser} / \gamma_0 (1 + \beta_0 \cos \theta) = \lambda_0$$



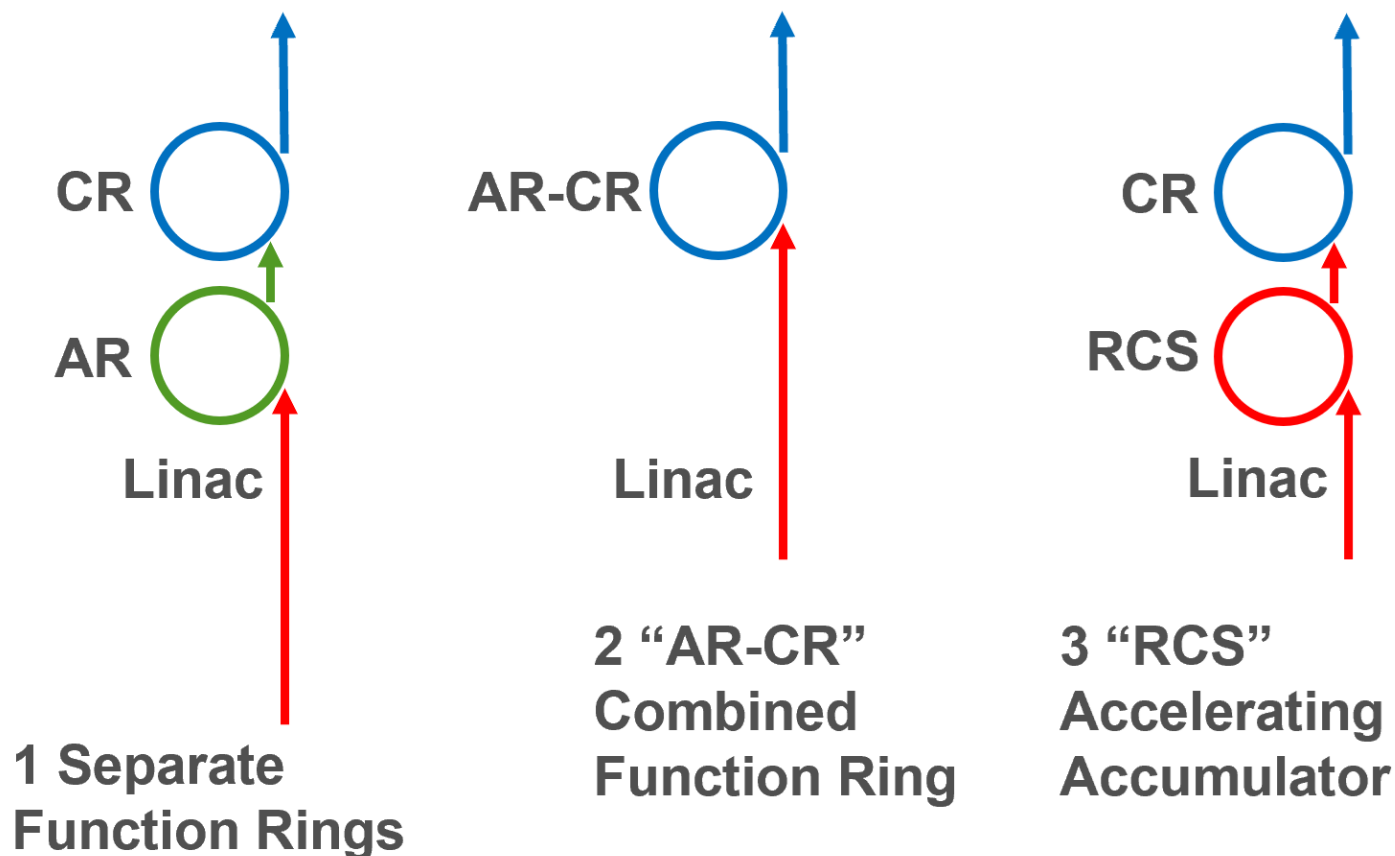
Changing the angle, any shorter wavelength laser can still be used at higher energies.

**Above ~4 GeV**, essentially the full range of laser technology and optics.

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# Accumulator Ring & Compressor Ring

# Architecture



Can the CR accumulate without losing compression performance?

Can the AR accelerate without losing accumulation performance?

For now keep the rings separate when in reality they may be the same...

# Space-charge Limit

Laslett Spacecharge  
Tuneshift parameter:

$$\Delta Q_{SC} = \frac{Nr_p}{2\pi\epsilon_n\beta_0\gamma^2} B_f F < \sim 0.2$$

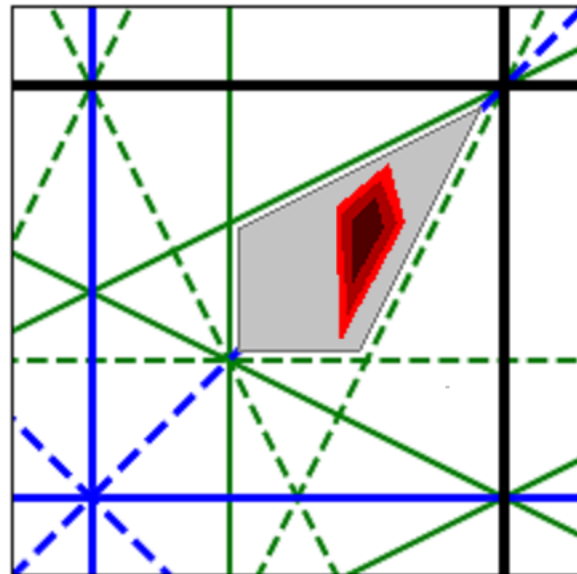
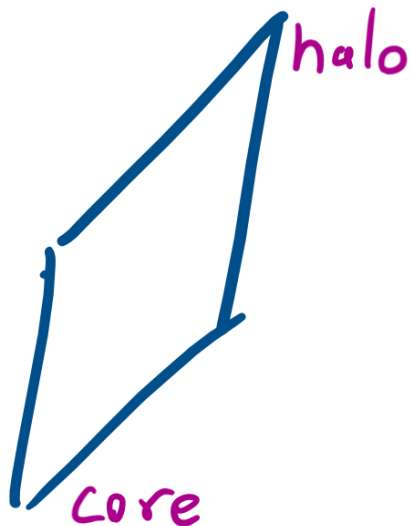
Number of particles  $\rightarrow N$

Transverse emittance (prop to beam-size squared)  $\rightarrow r_p$

Relativistic  $\beta, \gamma$  factors  $\rightarrow \beta_0, \gamma^2$

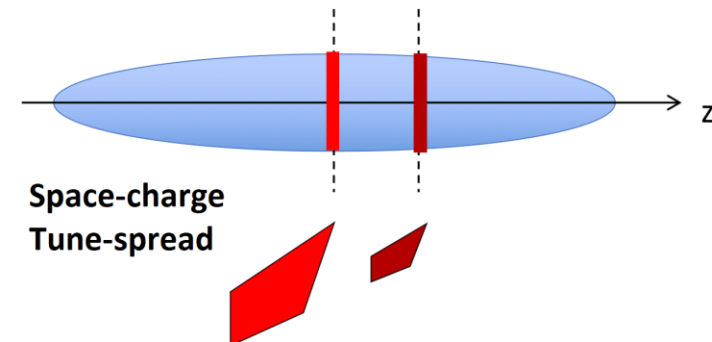
Bunching factor (inverse to bunch length)  $\rightarrow B_f$

Form factor (deviation from round gaussian)  $\rightarrow F$



Available Tunespace

Tune Footprint  
(density of particles)



# Space-charge Tune-shift for other Round Beams

## “Laslett” Space-charge Tune-shift:

$$\Delta\nu_{sc} = FB \frac{Nr_0}{2\pi\epsilon_N\beta\gamma^2}$$

**Form Factor F:**

$$F = \frac{\text{Max}[\rho(x, p_x)]}{\text{Max}[\rho_{KV}(x, p_x)]}$$

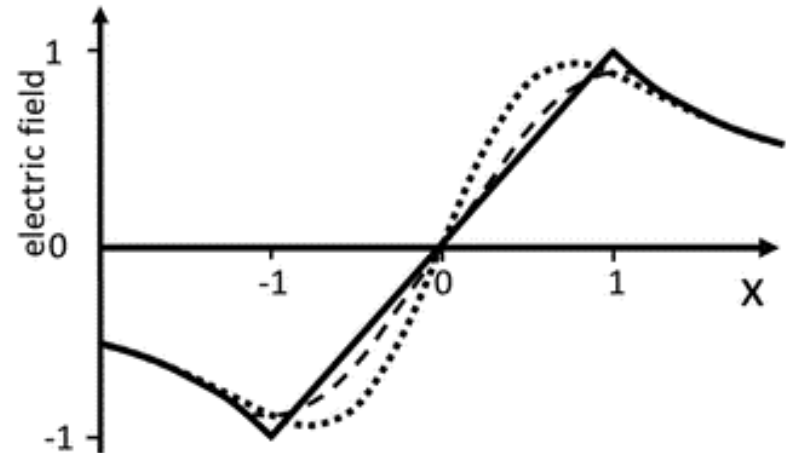
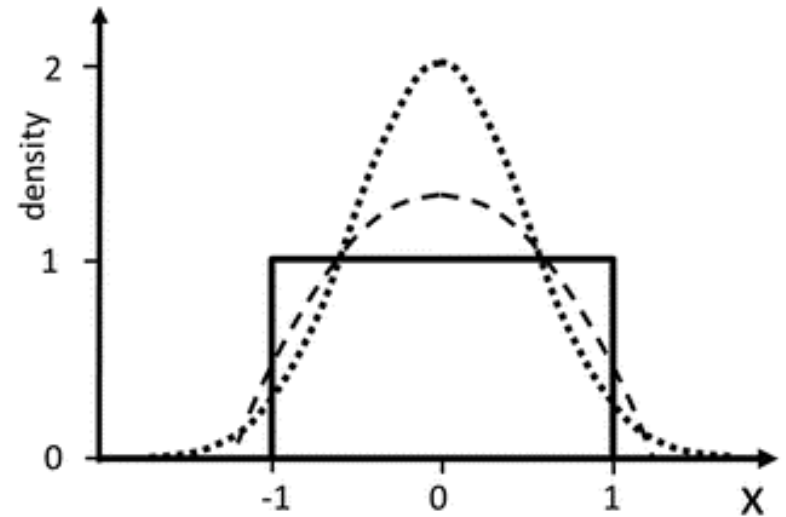
For a KV beam,  $F = 1$ , gauss  $F=1/2$ .

**Bunching Factor B:**

$$B_f = \frac{I(z)}{I_{ave}}$$

For a DC beam,  $B = 1$ .

For a partially filled ring, divide by fill factor.



# Space-charge Limit

Laslett Spacecharge Tuneshift parameter:

$$Q_{SC} = \frac{Nr_p}{4\pi\epsilon_n\beta_0\gamma^2} B_f F < \sim 0.2$$

Number of particles  $\rightarrow N$

Bunching factor (inverse to bunch length)  $\rightarrow B_f$

Form factor (deviation from round gaussian)  $\rightarrow F$

Transverse emittance (prop to beam-size squared)  $\rightarrow \epsilon_n$

Relativistic  $\beta, \gamma$  factors  $\rightarrow \beta_0, \gamma^2$

Bunching factor  $B_f$  means that the smaller the bunch length the more intense the space-charge force! Going from 2 $\mu$ s to 2ns a major challenge

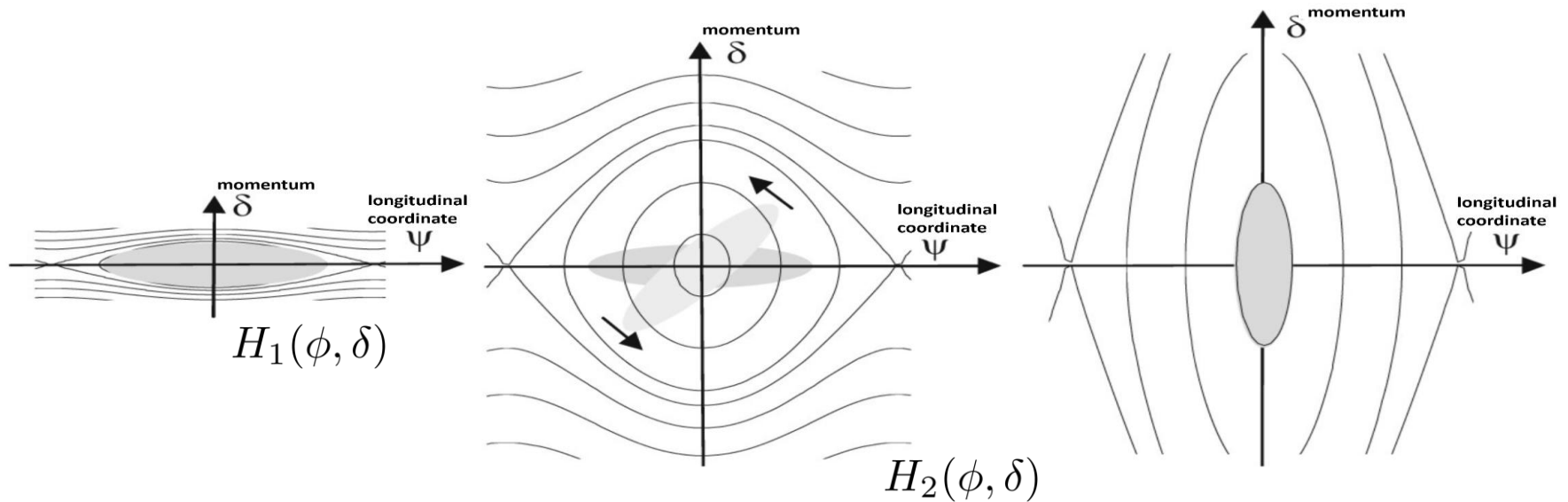
However also scales with  $\gamma^2$ , so accumulating at 5-10 GeV instead of at 0.4-1 GeV helps counterbalance.

Increasing emittance  $\epsilon_N$  is effective, but not cost-free for AR/CR design. Furthermore, we can't support an arbitrarily large spot size or divergence on muon production target.

More compact ARs and CRs are more effective, but can we use superconducting magnets in a messy charge-dominated ring?



# Snap Bunch Rotation in Phase-Space



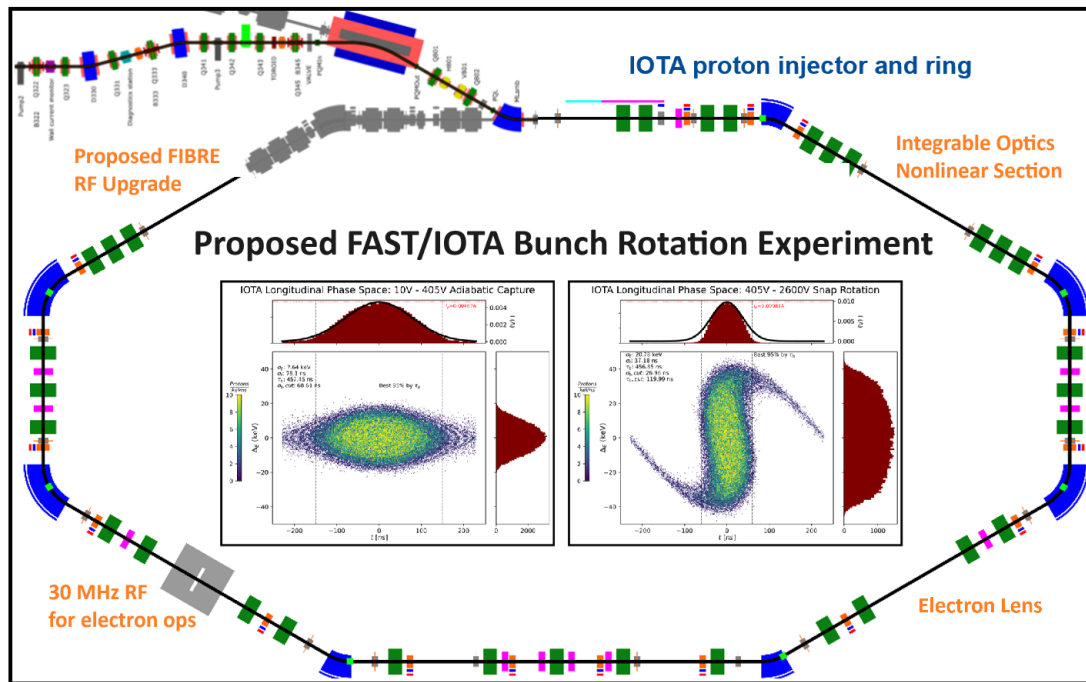
- 1) Start with as low as momentum spread beam as possible.
- 2) Jump up the voltage as high as possible as quickly as possible.
- 3) Particles follow new phase-space contours to rotate 1/4 synch period.

This achieves  $\sigma_{\delta,f} = \left(\frac{V_2}{V_1}\right)^{1/2} \sigma_{\delta,i}$ , versus adiabatic changes  $\sigma_{\delta,f} = \left(\frac{V_2}{V_1}\right)^{1/4} \sigma_{\delta,i}$ ,

Only the particles near the center full rotate, RF flattening may help.

There are also resonant schemes (modulate voltage at 1/2 synch period).

# Bunch Rotation & Strong Space-charge (R&D)



## Proposed FAST/IOTA Bunch Rotation Experiment (FIBRE)

- Space-charge R&D is a primary purpose of IOTA.
- Optimize bunch rotation in an extreme space-charge environment.
- Just needs a  $\sim 10$  kV  $\sim 2.2$  MHz RF cavity.

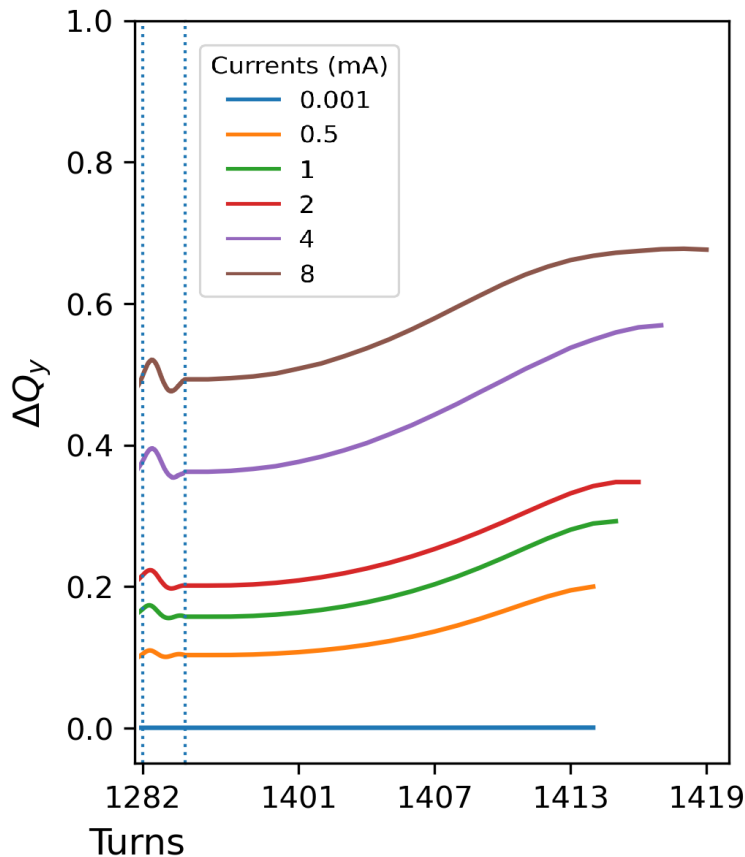
See poster from Ben Simons (NIU) for FAST/IOTA.

Other facility options ... ORNL SNS, RAL ISIS, FNAL Booster, GSI SIS...

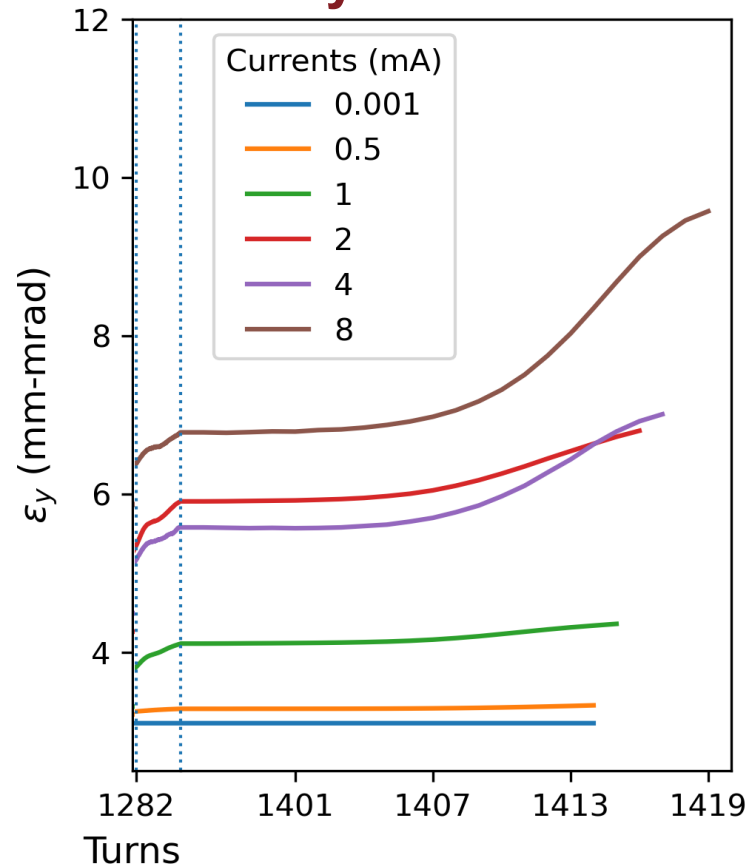
Talk from **Austin Hoover (ORNL)** in Muon Collider meeting.

- They have the RF, MW beams, new HEP application for BES facility.

# Bunch Rotation & Strong Space-charge (Simulation)



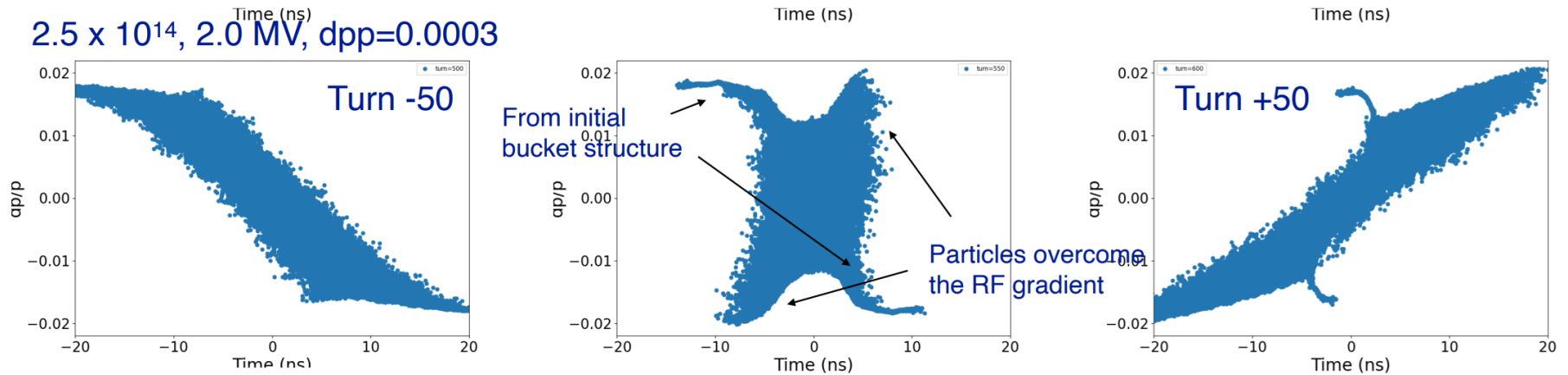
## Preliminary Results! Simons 2025



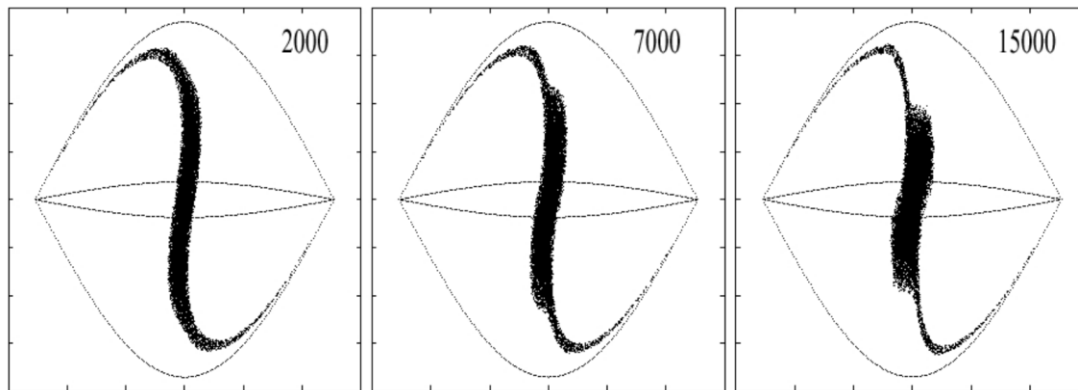
Can push the Laslett parameter into extreme values with manageable impact on transverse emittances. Rotation occurs in tens of us!

# Bunch Rotation & Longitudinal SC

Shinji Machida, Proton Driver simulation 2025



Ng Simulation of IUCF in 2001

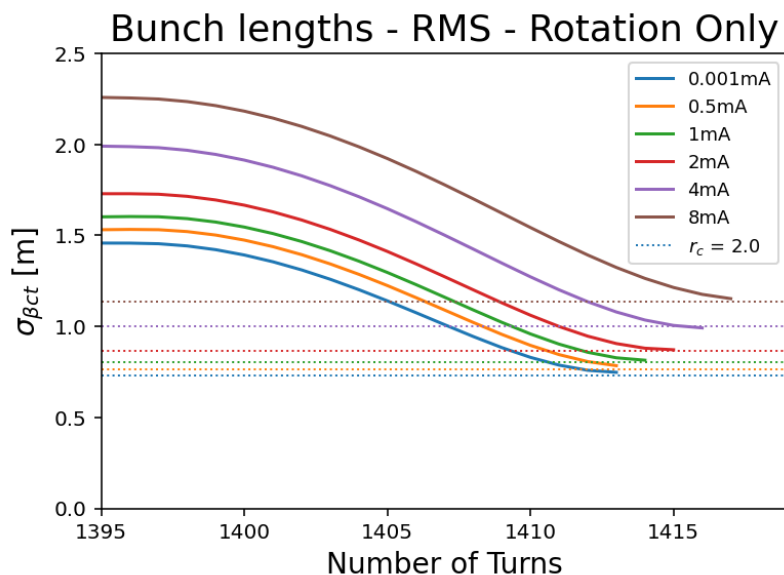


# Bunch Rotation & Longitudinal SC

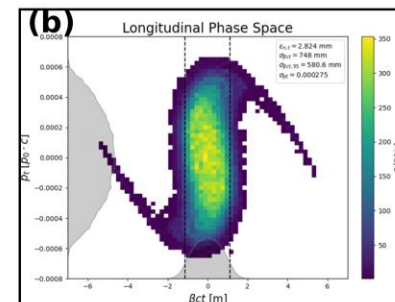
## Effective Longitudinal Focusing/defocusing

$$\left. \frac{\text{Sp-ch force}}{\text{Rf force}} \right|_{\text{critical}} = \frac{eN_b |Z/n|_{\text{spch}}}{\sqrt{2\pi} h \omega_0^2 \sigma_\tau^3 V_{\text{rf}}}$$

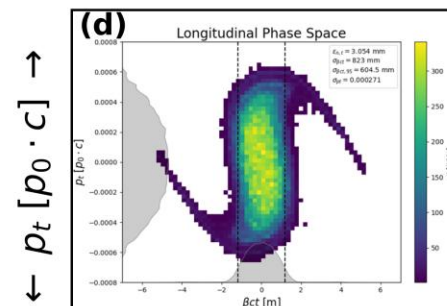
$$V_{\text{eff}} \approx -h \left( \frac{L_{\text{RF}}}{\sqrt{2\pi} \sigma} \right)^3 \frac{q k_e}{c} I_0 \frac{g}{\beta_0 \gamma_0^2}$$



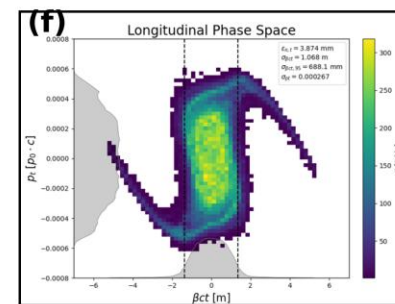
Simons 2025



1  $\mu\text{A}$



1 mA



4 mA

$\leftarrow \beta ct \text{ [m]} \rightarrow$

# Bunch Rotation & Phase-slip factor

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Should the phase-slip factor of the CR be large or small?

## Large $\eta$

- Most extreme SC occurs briefly.
- Mode-coupling instabilities are suppressed
- Higher voltage requirements allow more extreme bunch compression, needs more RF.
- Less bunch compression requires more of Linac/AR.

## Small $\eta$

- Beam should be SC stable.
- Wake effects accumulate from head to tail of bunch.
- Lower voltage requirements allow more extreme bunch compression given RF power
- More bunch compression relaxes requirements on Linac/AR

My conclusion? We will need to keep studying both.

*If the gains of larger  $\eta$  can be realized, it is worth it.*

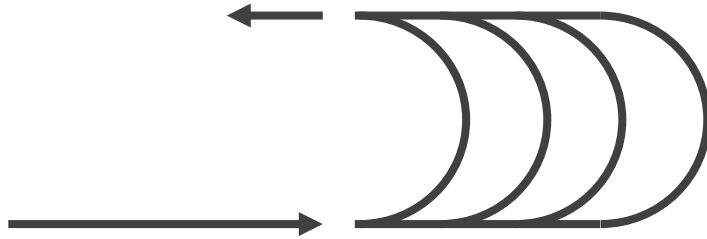
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# Combiner & Target



# Combiner Geometry

The combiner puts each bunch on a different path, so that the time of flights have all bunches coincide

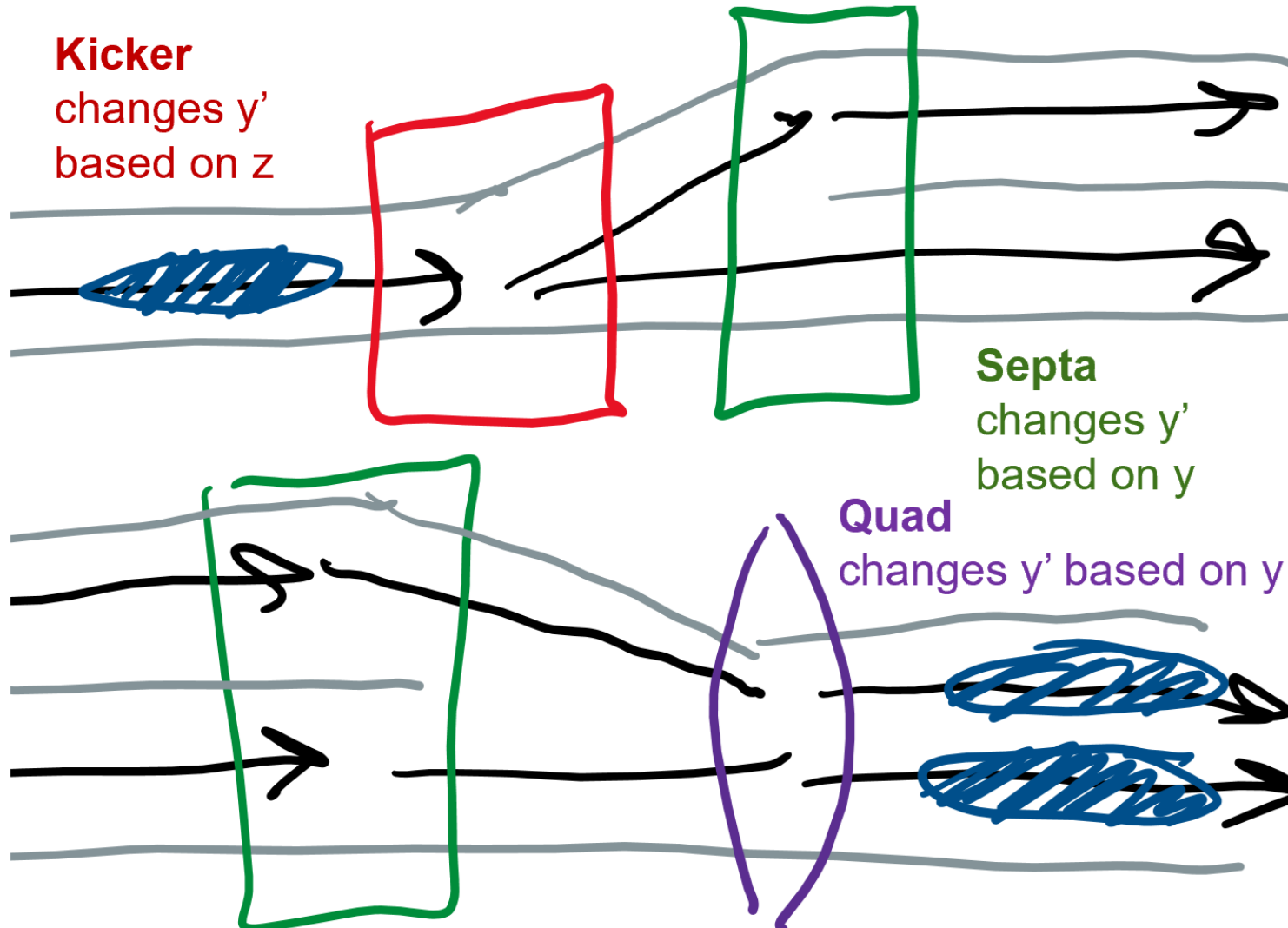


Each arc is roughly half the circumference  $C/2$  of the AR/CR.  
Then if  $N$  bunches fill up roughly  $f$  of the circumference, then the straight sections of the combined must each by  $fC/2$ .  
So the total length of all transfer lines is  $NC/2 + fC$ .

For modest numbers  $f = 0.25$ ,  $C = 300\text{m}$ ,  $N = 4$ , that's 675m of accelerator.  
Not much RF needed, but more than double the magnet/tunnel costs.

It will be worth if it can lower the linac or the CR/AR costs, or improve overall performance.

# Combiner Transverse Optics



# Impact of Multiple Bunches

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You cannot use a combiner to cheat Louiville's theorem, so the implications of using more bunches is that each bunch is narrower.

**What are the implication of using many bunches in a combiner?**

- Technical complexity and cost of combiner line.
- Emittance growth from combining bunches.
- Transverse aperture in CR and AR smaller, lower costs.
- Less SC during bunch compression process.
- Linac beam injects into a large azimuth of AR, which means that a smaller bunch charge for the same macropulse average current.

**My Conclusion?** More bunches may help, but don't use any more than you have to do in order to achieve performance goals.

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# Summary & Questions